

Complement H<sub>II</sub>

## EXERCISES

## Dirac notation. Commutators. Eigenvectors and eigenvalues

1.  $|\varphi_n\rangle$  are the eigenstates of a Hermitian operator  $H$  ( $H$  is, for example, the Hamiltonian of an arbitrary physical system). Assume that the states  $|\varphi_n\rangle$  form a discrete orthonormal basis. The operator  $U(m, n)$  is defined by:

$$U(m, n) = |\varphi_m\rangle \langle \varphi_n|$$

a. Calculate the adjoint  $U^\dagger(m, n)$  of  $U(m, n)$ .

b. Calculate the commutator  $[H, U(m, n)]$ .

c. Prove the relation:

$$U(m, n)U^\dagger(p, q) = \delta_{nq}U(m, p)$$

d. Calculate  $\text{Tr} \{ U(m, n) \}$ , the trace of the operator  $U(m, n)$ .

e. Let  $A$  be an operator, with matrix elements  $A_{mn} = \langle \varphi_m | A | \varphi_n \rangle$ . Prove the relation:

$$A = \sum_{m,n} A_{mn} U(m, n)$$

f. Show that  $A_{pq} = \text{Tr} \{ AU^\dagger(p, q) \}$ .

2. In a two-dimensional vector space, consider the operator whose matrix, in an orthonormal basis  $\{ |1\rangle, |2\rangle \}$ , is written:

$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

a. Is  $\sigma_y$  Hermitian? Calculate its eigenvalues and eigenvectors (giving their normalized expansion in terms of the  $\{ |1\rangle, |2\rangle \}$  basis).

b. Calculate the matrices which represent the projectors onto these eigenvectors. Then verify that they satisfy the orthogonality and closure relations.

c. Same questions for the matrices:

$$M = \begin{pmatrix} 2 & i\sqrt{2} \\ -i\sqrt{2} & 3 \end{pmatrix}$$

and, in a three-dimensional space

$$L_y = \frac{\hbar}{2i} \begin{pmatrix} 0 & \sqrt{2} & 0 \\ -\sqrt{2} & 0 & \sqrt{2} \\ 0 & -\sqrt{2} & 0 \end{pmatrix}$$

3. The state space of a certain physical system is three-dimensional. Let  $\{|u_1\rangle, |u_2\rangle, |u_3\rangle\}$  be an orthonormal basis of this space. The kets  $|\psi_0\rangle$  and  $|\psi_1\rangle$  are defined by:

$$|\psi_0\rangle = \frac{1}{\sqrt{2}}|u_1\rangle + \frac{i}{2}|u_2\rangle + \frac{1}{2}|u_3\rangle$$

$$|\psi_1\rangle = \frac{1}{\sqrt{3}}|u_1\rangle + \frac{i}{\sqrt{3}}|u_3\rangle$$

a. Are these kets normalized?

b. Calculate the matrices  $\rho_0$  and  $\rho_1$  representing, in the  $\{|u_1\rangle, |u_2\rangle, |u_3\rangle\}$  basis, the projection operators onto the state  $|\psi_0\rangle$  and onto the state  $|\psi_1\rangle$ . Verify that these matrices are Hermitian.

4. Let  $K$  be the operator defined by  $K = |\varphi\rangle\langle\psi|$ , where  $|\varphi\rangle$  and  $|\psi\rangle$  are two vectors of the state space.

a. Under what condition is  $K$  Hermitian?

b. Calculate  $K^2$ . Under what condition is  $K$  a projector?

c. Show that  $K$  can always be written in the form  $K = \lambda P_1 P_2$  where  $\lambda$  is a constant to be calculated and  $P_1$  and  $P_2$  are projectors.

5. Let  $P_1$  be the orthogonal projector onto the subspace  $\mathcal{E}_1$ ,  $P_2$  the orthogonal projector onto the subspace  $\mathcal{E}_2$ . Show that, for the product  $P_1 P_2$  to be an orthogonal projector as well, it is necessary and sufficient that  $P_1$  and  $P_2$  commute. In this case, what is the subspace onto which  $P_1 P_2$  projects?

6. The  $\sigma_x$  matrix is defined by:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Prove the relation:

$$e^{i\alpha\sigma_x} = I \cos \alpha + i\sigma_x \sin \alpha$$

where  $I$  is the  $2 \times 2$  unit matrix.

7. Establish, for the  $\sigma_y$  matrix given in exercise 2, a relation analogous to the one proved for  $\sigma_x$  in the preceding exercise. Generalize for all matrices of the form:

$$\sigma_u = \lambda\sigma_x + \mu\sigma_y$$

with:

$$\lambda^2 + \mu^2 = 1$$

Calculate the matrices representing  $e^{2i\sigma_x}$ ,  $(e^{i\sigma_x})^2$  and  $e^{i(\sigma_x + \sigma_y)}$ . Is  $e^{2i\sigma_x}$  equal to  $(e^{i\sigma_x})^2$ ?  $e^{i(\sigma_x + \sigma_y)}$  to  $e^{i\sigma_x} e^{i\sigma_y}$ ?

8. Consider the Hamiltonian  $H$  of a particle in a one-dimensional problem defined by:

$$H = \frac{1}{2m} P^2 + V(X)$$

where  $X$  and  $P$  are the operators defined in § E of chapter II and which satisfy the relation:  $[X, P] = i\hbar$ . The eigenvectors of  $H$  are denoted by  $|\varphi_n\rangle$ :  $H|\varphi_n\rangle = E_n|\varphi_n\rangle$ , where  $n$  is a discrete index.

a. Show that:

$$\langle \varphi_n | P | \varphi_{n'} \rangle = \alpha \langle \varphi_n | X | \varphi_{n'} \rangle$$

where  $\alpha$  is a coefficient which depends on the difference between  $E_n$  and  $E_{n'}$ . Calculate  $\alpha$  (hint: consider the commutator  $[X, H]$ ).

b. From this, deduce, using the closure relation, the equation:

$$\sum_{n'} (E_n - E_{n'})^2 |\langle \varphi_n | X | \varphi_{n'} \rangle|^2 = \frac{\hbar^2}{m^2} \langle \varphi_n | P^2 | \varphi_n \rangle$$

9. Let  $H$  be the Hamiltonian operator of a physical system. Denote by  $|\varphi_n\rangle$  the eigenvectors of  $H$ , with eigenvalues  $E_n$ :

$$H|\varphi_n\rangle = E_n|\varphi_n\rangle$$

a. For an arbitrary operator  $A$ , prove the relation:

$$\langle \varphi_n | [A, H] | \varphi_n \rangle = 0.$$

b. Consider a one-dimensional problem, where the physical system is a particle of mass  $m$  and of potential energy  $V(X)$ . In this case,  $H$  is written:

$$H = \frac{1}{2m} P^2 + V(X)$$

$\alpha$ . In terms of  $P$ ,  $X$  and  $V(X)$ , find the commutators:  $[H, P]$ ,  $[H, X]$  and  $[H, XP]$ .

$\beta$ . Show that the matrix element  $\langle \varphi_n | P | \varphi_n \rangle$  (which we shall interpret in chapter III as the mean value of the momentum in the state  $|\varphi_n\rangle$ ) is zero.

$\gamma$ . Establish a relation between  $E_k = \langle \varphi_n | \frac{P^2}{2m} | \varphi_n \rangle$  (the mean value of the kinetic energy in the state  $|\varphi_n\rangle$ ) and  $\langle \varphi_n | X \frac{dV}{dX} | \varphi_n \rangle$ . Since the mean value of the potential energy in the state  $|\varphi_n\rangle$  is  $\langle \varphi_n | V(x) | \varphi_n \rangle$ , how is it related to the

mean value of the kinetic energy when:

$$V(X) = V_0 X^\lambda$$

( $\lambda = 2, 4, 6 \dots$ ;  $V_0 > 0$ )?

**10** Using the relation  $\langle x | p \rangle = (2\pi\hbar)^{-1/2} e^{ipx/\hbar}$ , find the expressions  $\langle x | XP | \psi \rangle$  and  $\langle x | PX | \psi \rangle$  in terms of  $\psi(x)$ . Can these results be found directly by using the fact that in the  $\{ | x \rangle \}$  representation,  $P$  acts like  $\frac{\hbar}{i} \frac{d}{dx}$  ?

**Sets of commuting observables and C.S.C.O.'S**

**11.** Consider a physical system whose three-dimensional state space is spanned by the orthonormal basis formed by the three kets  $| u_1 \rangle, | u_2 \rangle, | u_3 \rangle$ . In the basis of these three vectors, taken in this order, the two operators  $H$  and  $B$  are defined by:

$$H = \hbar\omega_0 \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \quad B = b \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

where  $\omega_0$  and  $b$  are real constants.

- a. Are  $H$  and  $B$  Hermitian?
- b. Show that  $H$  and  $B$  commute. Give a basis of eigenvectors common to  $H$  and  $B$ .
- c. Of the sets of operators:  $\{ H \}, \{ B \}, \{ H, B \}, \{ H^2, B \}$ , which form a C.S.C.O.?

**12.** In the same state space as that of the preceding exercise, consider two operators  $L_z$  and  $S$  defined by:

$$\begin{aligned} L_z | u_1 \rangle &= | u_1 \rangle & L_z | u_2 \rangle &= 0 & L_z | u_3 \rangle &= - | u_3 \rangle \\ S | u_1 \rangle &= | u_3 \rangle & S | u_2 \rangle &= | u_2 \rangle & S | u_3 \rangle &= | u_1 \rangle \end{aligned}$$

- a. Write the matrices which represent, in the  $\{ | u_1 \rangle, | u_2 \rangle, | u_3 \rangle \}$  basis, the operators  $L_z, L_z^2, S, S^2$ . Are these operators observables?
- b. Give the form of the most general matrix which represents an operator which commutes with  $L_z$ . Same question for  $L_z^2$ , then for  $S^2$ .
- c. Do  $L_z^2$  and  $S$  form a C.S.C.O. ? Give a basis of common eigenvectors.