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NO REFERENCIAL CU, DO PROCESSO AB → CD TEMOS,

$$dQ = (2\pi)^4 \delta^{(4)}(p_c + p_D - p_A - p_B) \frac{d^3 p_c}{(2\pi)^3 2E_c} \frac{d^3 p_D}{(2\pi)^3 2E_D}$$

USAMOS A DADA  $\delta^{(4)}(\vec{p}_c + p_D - \vec{p}_A - p_B)$ , ENTÃO  
 $\vec{p}_c + \vec{p}_D = \vec{p}_A + \vec{p}_B$

$$dQ = (2\pi)^4 \frac{d^3 p_c}{(2\pi)^3 2E_c} \frac{1}{2E_D} \delta^{(4)}(p_c + p_D - p_A - p_B)$$

$$\delta^{(4)}(\vec{E}_c + E_D - E_A - E_B)$$

INTEGRAL DE FUNÇÃO DETA

$$\int \delta(g(x)) dx = \int \frac{\delta(x - x_i)}{|g'(x_i)|}$$

OMÉ KE SIA A  
 RAÍZ DE  $\delta(K_e) = 0$

~~MESMO CASO~~  $d^3 p_c = |\vec{p}_c|^2 d|\vec{p}_c| d\Omega$

$$\int \frac{d|\vec{p}_c| |\vec{p}_c|^2 d\Omega}{(2\pi)^2 2E_c 2E_D} \delta^{(4)}(E_c + E_D - E_A - E_B)$$

$$g(p_c) = \sqrt{p_c^2 + m_c^2} + \sqrt{p_D^2 + m_B^2} + \cancel{E_A} + \cancel{E_B}$$

↑  
↳  $p_A + p_B = p_c$

$$g(p_c) = 0 \Rightarrow \cancel{E_A} + \cancel{E_B}$$

$$\int \frac{p_c^2 dp_c}{2\tilde{E}_c 2\tilde{E}_D} \delta^0(\tilde{E}_c + \tilde{E}_D - E_A - E_B)$$

$$E_c^2 = p_c^2 + m_c^2$$

$$2\tilde{E}_c d\tilde{E}_c = p_c dp_c$$

$$\int \frac{E_c dE_c \sqrt{E_c^2 - m_c^2}}{2\tilde{E}_c 2\tilde{E}_D} \delta^0(\tilde{E}_c + \tilde{E}_D - E_A - E_B)$$

NÃO PODEMOS  
TRAZÍ-LOS.

$$E_D^2 = (p_A + p_B - p_c)^2 + m_B^2$$

$$g'(p_c) = \frac{p_c}{\sqrt{p_c^2 + m_c^2}} + \frac{(-1) p_D}{\sqrt{p_D^2 + m_B^2}} = \frac{p_c}{\tilde{E}_c} - \frac{p_D}{\tilde{E}_D} = p_D \frac{(E_D + E_c)}{\tilde{E}_c \tilde{E}_D}$$

$$\int \frac{p_c^2 dp_c}{2\tilde{E}_c 2\tilde{E}_D} \frac{1}{|p_c - p_D|} \frac{d}{dp_c} \sqrt{(p_A + p_B - p_c)^2 + m_B^2} =$$

$$\frac{(1/2) 2(p_D - p_c)}{\sqrt{\dots}}$$

EMÃO

$$\Gamma(A \rightarrow 1+2) = \frac{1}{32\pi^2 m_A} \int |\mathcal{M}|^2 dR$$



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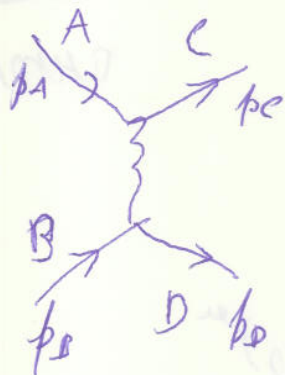
(41)

SEÇÃO 4.5 DO URZEW

ESPALHAMENTO ELETRON-ELÉTRON

PRETENDIA IDENTIFICAR  
NO ESTADO INICIAL (P<sub>1</sub>)

MESTE COMO ALÉM DO DUCHNA



PODEROS TNR

OUTRO DIAGRAMA TROCANDO (A ↔ D) (B ↔ C)



QUEM DIAGRAMA OUTRO? AMBOS PODEM NÃO PODEREM

DISCRIMINAR COMO FOI FEITO O ESPALHAMENTO.



$$-iM = \left[ (tie)(p_A + p_C)^\mu \left( \frac{-i\gamma_\nu}{(p_B - p_C)^2} \right) (tie)(p_B + p_D)^\nu \right.$$

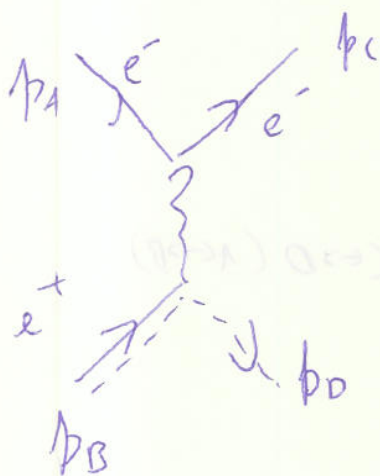
$$\left. + (tie)(p_A + p_D)^\mu \left( \frac{-i\gamma_\nu}{(p_C - p_B)^2} \right) (tie)(p_B + p_C)^\nu \right]$$

O PROCESSO DE ESPALHAMENTO ELETRON-POSITRON

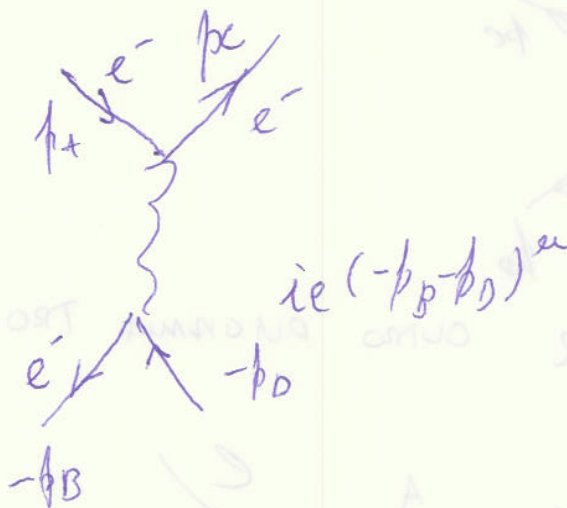
OU ESPALHAMENTO ~~REACTIVO~~

$$e^+ e^- \rightarrow e^+ e^-$$

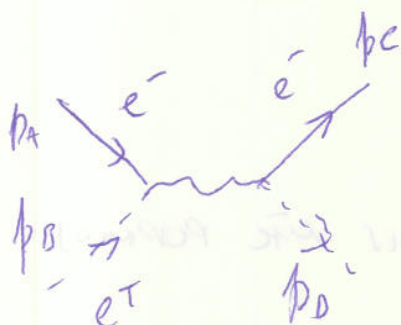
DIAGRAMA



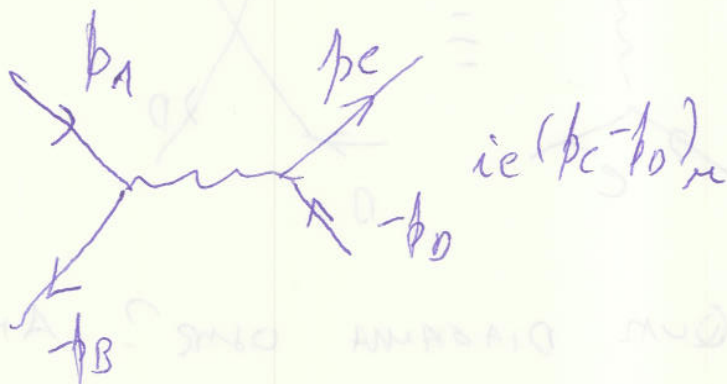
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E O OUTRO DIAGRAMA,



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43

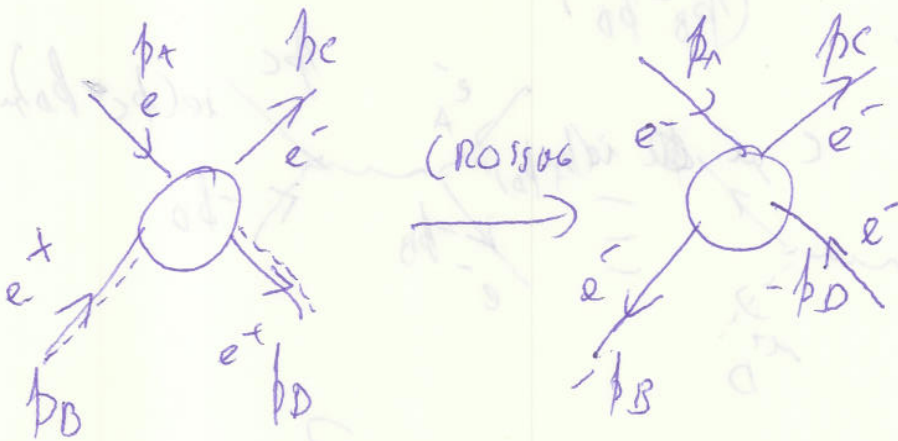
ERFAC PORTAS ESCRITA

IDENTIC

$$-id_{e^-} = \left\{ ie (-p_D - p_B)^{\mu} \left( \frac{-i g_{\mu\nu}}{(-p_D + p_B)^2} \right) \right\} ie (p_A + p_C)^{\nu}$$

$$+ ie (p_A - p_B)^{\mu} \left( \frac{-i g_{\mu\nu}}{(p_C + p_D)^2} \right) (-p_D + p_C)^{\nu}$$

OU PORTAS USAR O CROSSING



$$M_{e^- e^+ \rightarrow e^- e^+} (p_A, p_B, p_C, p_D) = M_{e^- e^+ \rightarrow e^- e^+} (p_A, p_D, p_C, p_B)$$

$$AD \rightarrow CB$$

$$AB \rightarrow CD$$

ESTA EXPRESSÃO É SIMÉTRICA

POIS TRATA DE

DOS ELÉTRONS  $p_C \leftrightarrow -p_B$

NAZEM PROCESSO

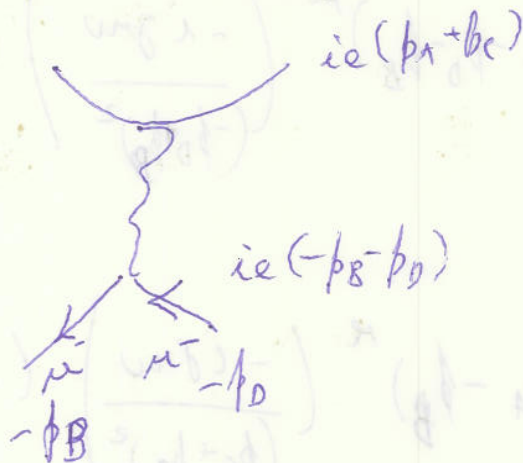
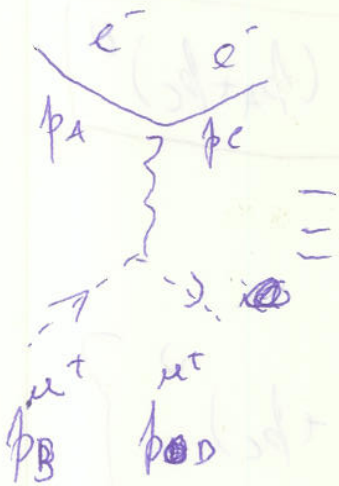
$$\begin{aligned} e^- \mu^+ &\rightarrow e^- \mu^+ \\ e^- e^+ &\rightarrow \mu^- \mu^+ \end{aligned}$$

$$e^- \mu^- \rightarrow e^- \mu^-$$

$$A B \rightarrow C D$$

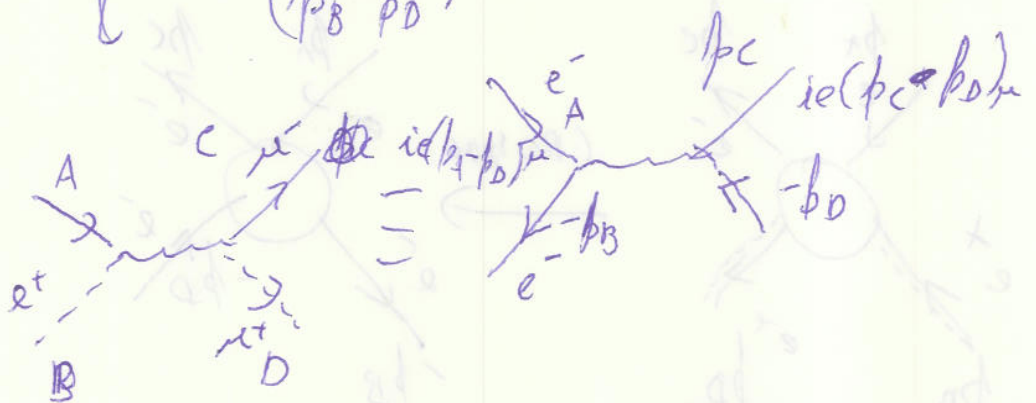
$$1) A \bar{D} \rightarrow C \bar{B}$$

$$2) A \bar{C} \rightarrow \bar{B} D$$



$$-i \mathcal{M}_{e^- \mu^+ \rightarrow e^- \mu^+} = -i \left[ \frac{-e^2 (p_A + p_C)_\mu (-p_B - p_D)^\mu}{(p_B - p_D)^2} \right]$$

E  $e^- e^+ \rightarrow \mu^- \mu^+$



$$-i \mathcal{M}_{e^- e^+ \rightarrow \mu^- \mu^+} = -i \left[ \frac{-e^2 (p_A - p_B)^\mu (p_C - p_D)_\mu}{(p_C + p_D)^2} \right]$$



AS VARIÁVEIS DE HAMILTON ~~REPRESENTAM~~ SÃO USADAS PARA ESCREVER A AMPLITUDE DE FORMA COVARIANTE. ESTAS VARIÁVEIS TÊM DIFERENTES SIGNIFICADOS FÍSICOS.

EM GERAL DESCRIVEMOS AS REGIÕES QUANTÍCAS PERMITIDAS NA SEGUNTE FORMA, COMO AS TRÊS VARIÁVEIS TÊM O VÍNCULO

$$A + t + u = J \text{ m}^2$$

ENTÃO PODEMOS DIZER UMA DESCRIÇÃO COM GRÁFICO BIDIMENSIONAL.

NO CASO DE  $AB \rightarrow CD$  COM TOMBAS AS PARTÍCULAS IGUAIS

TEMOS

$$A + t + u = 4 \text{ m}^2$$

NO CASO EM QUE  $t = u = 0$ , ~~temos~~ ~~o~~ ~~valor~~ ~~máximo~~ ~~de~~ ~~A~~ OBTÉM O VALOR

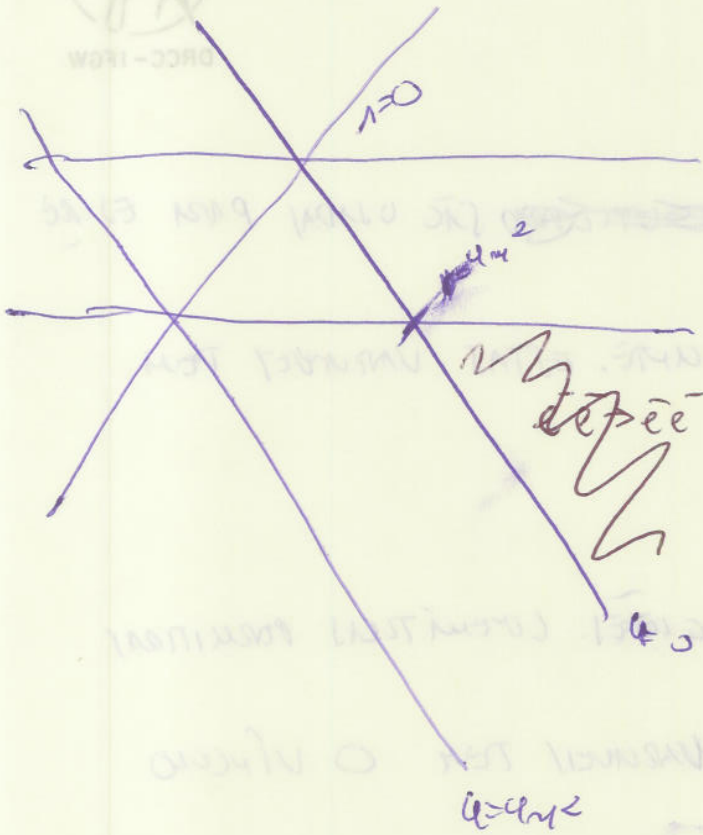
~~de~~ ~~A~~  $A = 4 \text{ m}^2$ , PORÉMOS MAIOR QUE NO CASO  $t = u = 0$

$t + u < 0$ , PORTANTO  $A = 4 \text{ m}^2$  É O VALOR MÁXIMO.



PURPOSOS FAZER O DIAGRAMA

$$(t = t_{24} = 4m^2)$$



$$t = 4m^2$$

$$t = 0$$

NO CASO DE

$$e^- e^- \rightarrow e^- e^-$$

A REGIÃO SOMBREADA É.

EM GERAL OS PROCESSOS TEM MAIS DE UMA DIAGRAMA,

COMO NO ESPANHAMENTO  $e^- e^- \rightarrow e^- e^-$ , ESTE PROCESSO É

<sup>PRINCIPAL</sup>  
(MUITO COMUM)

CONSIDERANDO UM PROCESSO  $QQ$ ,  $AB \rightarrow CD$  EXISTE

UMA DIAGRAMA QUE SE CORRESPONDE A  $\sigma = (E_{CM}^{TOT})^2$

NESTE CASO COMUNS DE CASO 1,  
PURPOSOS GERE

$AB \rightarrow CD$   $\left\{ \begin{array}{l} e^- e^- \rightarrow e^- e^- \\ e^- e^+ \rightarrow e^- e^+ \end{array} \right.$

E O CASO SEJA O CASO 1. !!



como

$$s = (\vec{p}_1 + \vec{p}_2)^2$$

ENTÃO NO REFERENCIAL CM, PARA PARTÍCULAS  
 COM MASSAS IGUAIS (M),

$$s = (\vec{p}_1 + \vec{p}_2)^2 = (E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2 = (E_{CM})^2$$

Como partículas com MOMA MISA  $\vec{p}_1 = -\vec{p}_2 \Rightarrow E_1 = E_2$

$$E_{CM} = E_1 + E_2 = 2E_1$$

CONSIDERANDO  $AB \rightarrow CD$  A REAÇÃO  $e^- e^+ \rightarrow e^- e^+$

ENTÃO ESTE QUANTUM É O CM S.



PARA ESTE PROCESSO,  $e^- e^+ \rightarrow e^- e^-$

$$s = (\vec{p}_1 + \vec{p}_2)^2 = (E_{CM})^2$$

NO SISTEMA CM  $\vec{p}_1 = -\vec{p}_2$ ,  $\vec{p}_C = -\vec{p}_D$

~~OU~~ OU  $|\vec{v}_C| = v$

$$s = (2(\gamma^2 v^2 + m^2))^2 = 4(\gamma^2 m^2)$$

$$t = (p_A - p_C)^2 = (p_i - p_o)^2$$

Logo

$$E_A + E_p = E_C + E_D$$

$$\Rightarrow E_A = E_C \Rightarrow E_A = E_C \Rightarrow$$

$$|\vec{v}_p| = |\vec{v}_i|$$

$$E_A = E_D$$

ASSUMINDO

$$\vec{v}_p \cdot \vec{v}_i = |\vec{v}|^2 \cos \theta$$

ENTÃO

$$t = (p_A - p_C)^2 = (E_A - E_C)^2 - (\vec{p}_A - \vec{p}_C)^2$$

$$= -(\vec{p}_A^2 + \vec{p}_C^2 - 2\vec{p}_A \cdot \vec{p}_C)$$

$$t = -(u^2 + u^2 - 2u^2 \cos \theta) = -2u^2(1 - \cos \theta)$$

PARA O TEMPO

$$\vec{p}_D = -\vec{p}_C \quad \vec{p}_A \cdot \vec{p}_D = -\vec{p}_A \cdot \vec{p}_C = -u^2 \cos \theta$$

$$u^2 (p_A - p_D)^2 = (E_A - E_D)^2 - (\vec{p}_A - \vec{p}_D)^2 = -[u^2 + u^2 + 2u^2 \cos \theta]$$

$$= -2u^2(1 + \cos \theta)$$

PORTANTO

$$s > 4m^2, t < 0$$

PMU  $\theta = 0 \Rightarrow t = 0 \Rightarrow$  ESPALHAMENTO FRONTAL

$\theta = \pi \Rightarrow u = 0 \Rightarrow$  ESPALHAMENTO TRASEIRO

$$s = 4(u^2 + m^2)$$

$$t = -2u^2(1 - \cos \theta)$$

$$u = -2u(1 + \cos \theta)$$

OS PROCESSOS  $A\bar{D} \rightarrow C\bar{B}$   $\bar{D}B \rightarrow C\bar{A}$  SÃO O CASO  $\varphi$   $\pi$ .

SE  $AB \rightarrow CD$   $e^-e^+ \rightarrow e^-e^+$   $e^-e^+ \rightarrow e^-e^+$   $e^-e^+ \rightarrow e^-e^+$

DA REAÇÃO ANTERIOR

$$M_{e^-e^+ \rightarrow e^-e^+}(p_A, p_B, p_C, p_D) = M_{e^-e^+ \rightarrow e^-e^+}(p_A, -p_B, p_C, p_D)$$

A REAÇÃO REVERSA É

$$M_{e^-e^+ \rightarrow e^-e^+}(p_A, p_B, p_C, p_D) = M_{e^-e^+ \rightarrow e^-e^+}(p_A, -p_D, p_C, -p_B)$$

$$s = (p_A + p_B)^2 \rightarrow u = (p_A - p_D)^2$$

AS VARIÁVEIS FICAM

$$t = (p_A - p_C)^2 \rightarrow u' = (p_A - p_C)^2$$

$$u = (p_A - p_D)^2 \rightarrow s' = (p_A + p_B)^2$$

ENTÃO A SIMETRIA DE CROSSING PODE SER INTERPRETADA

COMO A TROCA DAS VARIÁVEIS DE MANDSTAM.



NESTE CASO A REAÇÃO OCORRE EM ~~AB~~ ~~CD~~  
 $e^-e^- \rightarrow e^-e^-$   
 COM VELOCIDADES  $v, x, u$  QUE SÃO CROSSES DE

TRANSFORMAÇÃO NA REAÇÃO ~~AB~~ ~~CD~~  
~~AB~~ ~~CD~~  
 $A\bar{D} \rightarrow C\bar{B}$

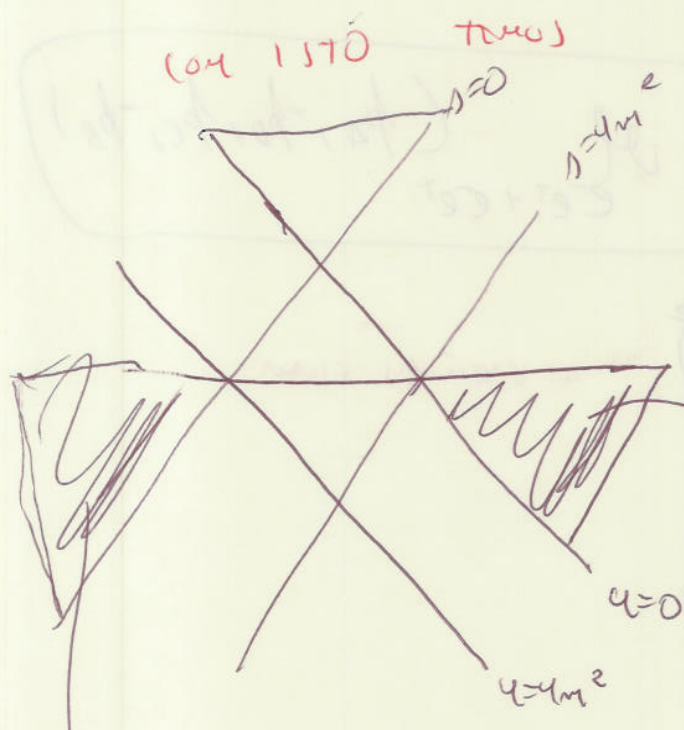
NESTA LOCALIDADE

NOTE A ORDEM DAS VELOCIDADES.

$$\mathcal{M}(s, t, u) = \mathcal{M}(y, x, 0)$$

$$e^-e^- \rightarrow e^-e^-$$

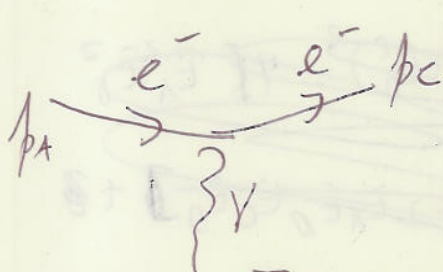
$$e^-e^- \rightarrow e^-e^-$$



PROCESSO NO CM  $u$   
 $AB \rightarrow CD$   
 $(e^-e^- \rightarrow e^-e^-)$   
 Tempo de  $\mathcal{M}(E_{CM}^2)$

PROCESSO NO CM  $u$   
 $e^-e^- \rightarrow e^-e^-$   
 Tempo de  $u = (E_{CM}^2)^2$

SE O PROCESSO DO CM  $\gamma \rightarrow e^- \mu^- \rightarrow e^- \mu^-$

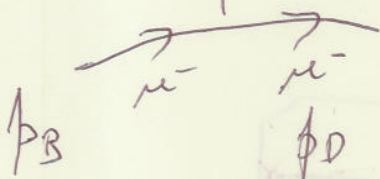


$$s = (p_A + p_B)^2$$

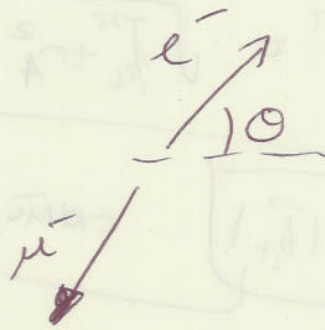
$$t = (p_A - p_C)^2$$

$$u = (p_A - p_D)^2$$

$$u + t = 2(m_\mu^2 - m_e^2)$$



NO CM,



COMO  $\vec{p}_B = -\vec{p}_A$

$$E_A + E_B = E_C + E_D$$

$$p_A = (E_A, \vec{p}_A)$$

$$p_C = (E_C, \vec{p}_C)$$

$$p_B = (E_B, \vec{p}_B)$$

$$p_D = (E_D, \vec{p}_D)$$

$$\sqrt{p_A^2 + m_e^2} + \sqrt{p_B^2 + m_\mu^2} = \sqrt{p_C^2 + m_e^2} + \sqrt{p_D^2 + m_\mu^2}$$

$$p_A^2 + m_e^2 + p_B^2 + m_\mu^2 = p_C^2 + m_e^2 + p_D^2 + m_\mu^2$$

$$p_A^2 + p_B^2 = p_C^2 + p_D^2$$

$$E_A^2 - p_A^2 + E_B^2 - p_B^2 = E_C^2 - p_C^2 + E_D^2 - p_D^2$$

$$E_A^2 + E_B^2 - p_A^2 - p_B^2 = E_C^2 + E_D^2 - p_C^2 - p_D^2$$

$$s = (p_A + p_B)^2 = (E_A + E_B)^2 - (\vec{p}_A + \vec{p}_B)^2$$

$$t = (p_A - p_C)^2 = (E_A - E_C)^2 - (\vec{p}_A - \vec{p}_C)^2$$

$$4(E_A E_B - E_C E_D) = E_A^2 + E_B^2 - E_C^2 - E_D^2 + 2E_A E_B - 2E_C E_D$$

$$4(E_A E_B - E_C E_D) = (E_A^2 + E_B^2 - E_C^2 - E_D^2) + 2(E_A E_B - E_C E_D)$$

$$2(E_A E_B - E_C E_D) = (E_A^2 + E_B^2 - E_C^2 - E_D^2)$$



$$\left( \cancel{E_c + E_D} - \cancel{E_A - E_B} \right)^2 = \left( \cancel{E_A - E_B} - \cancel{E_c + E_D} \right)^2$$

$$\cancel{E_c + E_D} - \cancel{E_A - E_B}$$

$$\cancel{(E_c + E_D)^2 - 2(E_A - E_B)(E_c + E_D) + (E_A - E_B)^2} = \cancel{4E_A E_B}$$

$$\cancel{(E_c + E_D)^2 - 2E_A E_B - (E_c + E_D)^2} = \cancel{4E_A E_B}$$

$$\cancel{(E_c + E_D)^2 + (E_A - E_B)^2 - 2}$$

$$\sqrt{p_A^2 + m_A^2} + \sqrt{p_B^2 + m_B^2} = \sqrt{p_C^2 + m_A^2} + \sqrt{p_C^2 + m_B^2}$$

solução  $|\vec{p}_C| = |\vec{p}_A| \Rightarrow m_A c$

$$\begin{cases} E_A = E_C \\ E_B = E_D \end{cases}$$

$$x = (E_A - E_C)^2 - (\vec{p}_A - \vec{p}_C)^2 = -2\beta^2(1 - \cos\theta)$$

$$E_A = \sqrt{p_A^2 + m_A^2}$$

$$E_B = \sqrt{p_A^2 + m_B^2}$$

$$y = (E_A - E_B)^2 = \cancel{E_A^2 - 2E_A E_B + E_B^2}$$

$$\cancel{E_A^2 + E_B^2 - 2E_A E_B + p_A^2 + m_A^2 + p_A^2 + m_B^2 - 2p_A^2 - 2m_A m_B \cos\theta}$$

$$y = (E_A + E_B)^2 - (m_A + m_B)^2$$

$$x = (p_A - p_B)^2 = (E_A - E_B)^2 - (\vec{p}_A - \vec{p}_B)^2$$



$$u+t = (E_A - E_B)^2$$

53

$$u = (E_A - E_B)^2 - (\vec{p}_A - \vec{p}_B)^2 = (E_A - E_B)^2 - (p_A^2 + p_B^2 - 2\vec{p}_A \cdot \vec{p}_B)$$

$$u = (E_A - E_B)^2 - p_A^2 (2 + 2\cos\theta) = (E_A - E_B)^2 - 2p_A^2 (1 + \cos\theta)$$

$$(E_A - E_B)^2 = E_A^2 + E_B^2 - 2E_A E_B \cos\theta$$

$$\left( \sqrt{p_A^2 + m_A^2} - \sqrt{p_B^2 + m_B^2} \right)^2 \leq (E_A - E_B)^2 - 4p_A^2 \leq (E_A - E_B)^2$$

$$E_A^2 + E_B^2 - 2E_A E_B \cos\theta - 4p_A^2 = 2p_A^2 + m_A^2 - m_B^2 - 2E_A E_B - 4p_A^2$$

$$= -2p_A^2 + m_A^2 + m_B^2$$

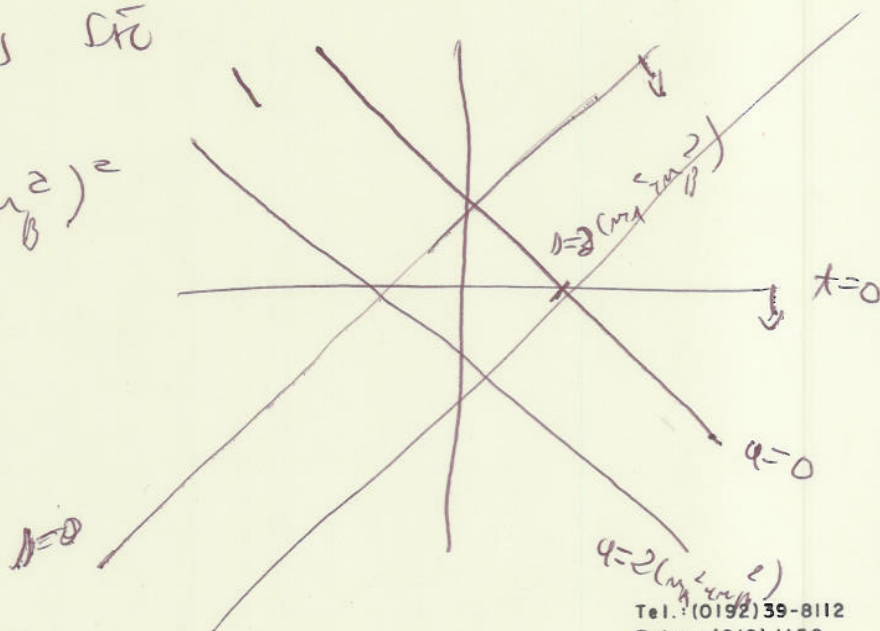
$$u_{\text{max}} = (E_A - E_B)^2$$

$$u > (E_A + E_B)^2$$

$$u_{\text{max}} = (E_A - E_B)^2 = (p_A^2 + m_A^2 - p_B^2 - m_B^2)^2 = (m_A^2 - m_B^2)^2$$

EM TÃO OS CASOS DE

$$\left. \begin{array}{l} u < (m_A^2 - m_B^2)^2 \\ t < 0 \end{array} \right\}$$





OBTIVAMOS TAMBEÉM QUE

$$\mathcal{M}_{e^-e^- \rightarrow e^-e^-}(s, t, u) = \mathcal{M}_{e^-e^- \rightarrow e^-e^-}(u, t, s)$$

PODEMOS

$$\mathcal{M}_{e^-e^- \rightarrow e^-e^-}(s, t, u) = \mathcal{M}_{e^-e^- \rightarrow e^-e^-}(u, t, s)$$

~~COM ESTO AGORA PODEMOS ESTABELECE~~

VAMOS ESTABELECE AS RELAÇÕES AUTOMÁTICAS EM TERMO DE

~~s, t, u.~~

$$p_1 + p_2 = k_1 + k_2$$

$$-i\mathcal{M}_{e^-e^+} = -i \left[ \frac{-e^2 (p_1 + p_2)_\mu (p_3 + p_4)^\mu}{(p_0 - p_2)^2} - \frac{-e^2 (p_1 + p_2)_\mu (p_3 + p_4)^\mu}{(p_2 - p_3)^2} \right]$$

$$-i\mathcal{M}_{e^-e^+} = -i \left[ \frac{-e^2 (p_1 + k_1)_\mu (-p_0 - p_3)^\mu}{(p_0 - p_3)^2} - \frac{-e^2 (p_1 - p_3)_\mu (-p_0 + k_1)^\mu}{(p_2 + p_3)^2} \right]$$



