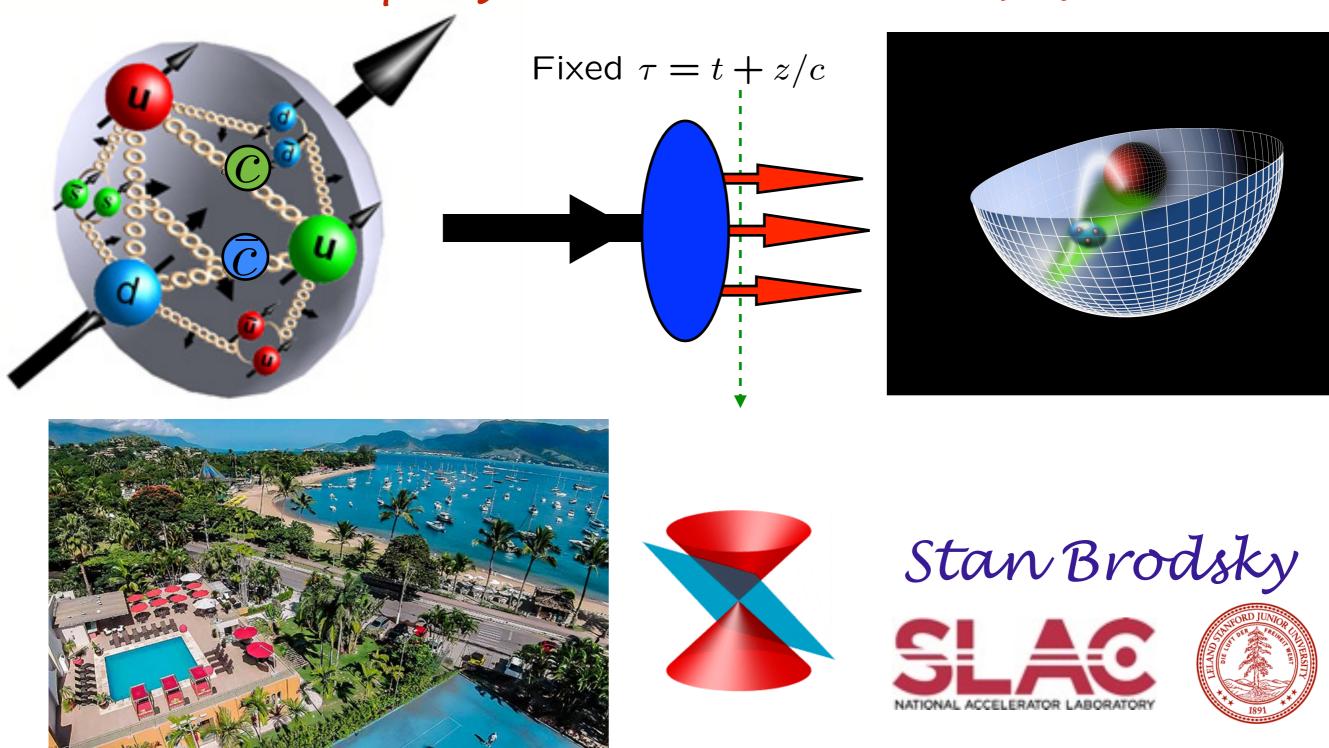
Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD

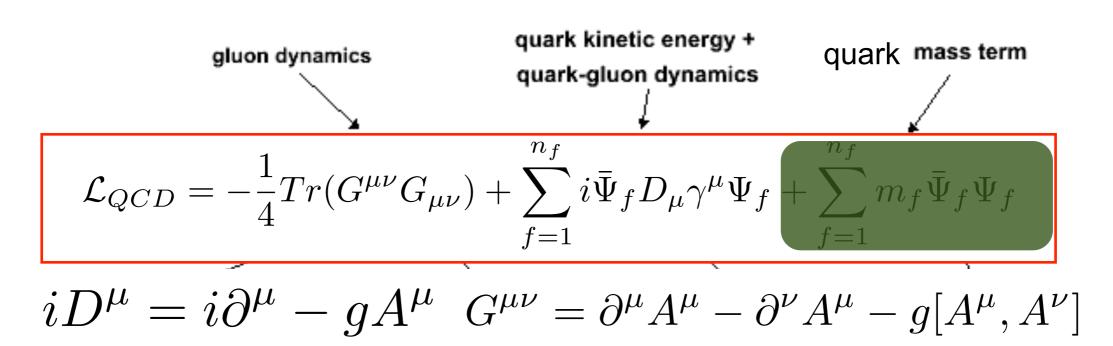


QCD-TNT4 Unraveling the Organization of the Tapestry of QCD

IhaBela, São Paulo, Brazil - August 31 to September 04, 2015

QCD Lagrangian

Fundamental Theory of Hadron and Nuclear Physics



Classically Conformal if $m_q=0$

Yang Mills Gauge Principle: Color Rotation and Phase Invariance at Every Point of Space and Time Scale-Invariant Coupling Renormalizable Asymptotic Freedom Color Confinement

QCD Mass Scale from Confinement not Explicit

Tony Zee

"Quantum Field Theory in a Nutshell"

Dreams of Exact Solvability

"In other words, if you manage to calculate m_P it better come out proportional to Λ_{QCD} since Λ_{QCD} is the only quantity with dimension of mass around.

Light-Front Holography:

Similarly for m_{ρ} .

$$m_p \simeq 3.21 \ \Lambda_{\overline{MS}}$$

$$m_{
ho} \simeq 2.2 \ \Lambda_{\overline{MS}}$$

Put in precise terms, if you publish a paper with a formula giving m_{ρ}/m_{P} in terms of pure numbers such as 2 and π , the field theory community will hail you as a conquering hero who has solved QCD exactly."

$$(m_q = 0)$$

$$m_{\pi} = 0$$

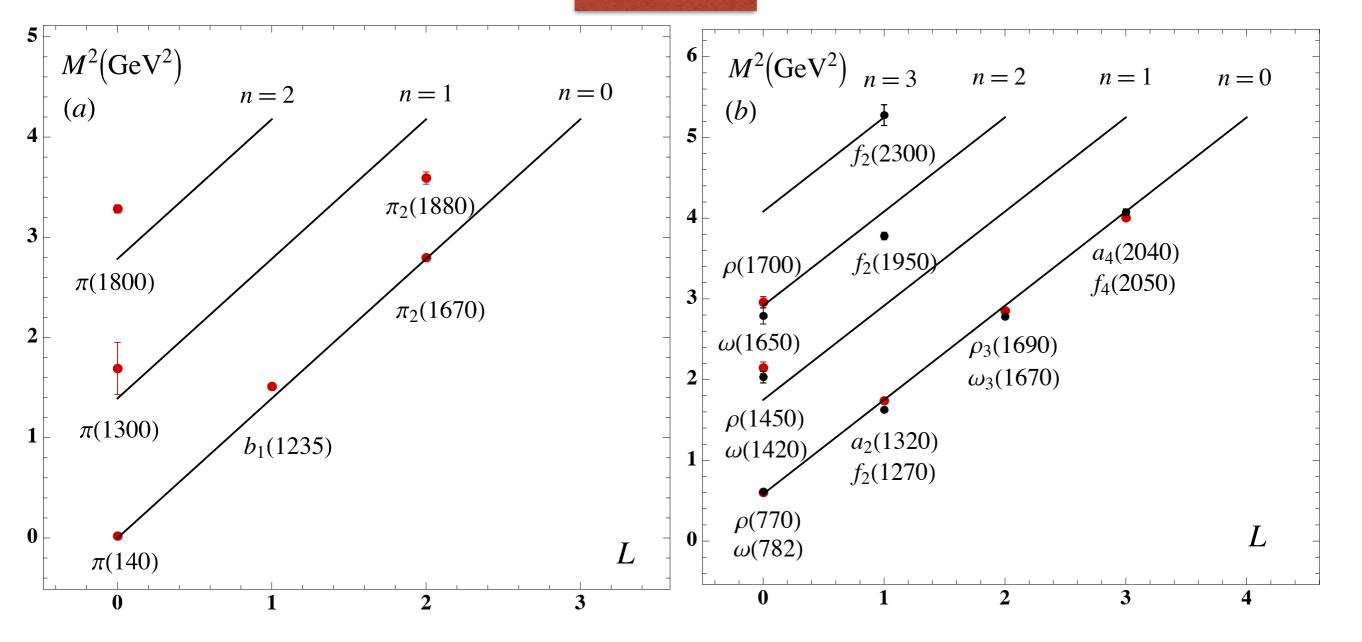
$$\frac{m_{\rho}}{m_{P}} = \frac{1}{\sqrt{2}}$$

$$\frac{\Lambda_{\overline{MS}}}{m_{\rho}} = 0.455 \pm 0.031$$

$$m_u = m_d = 0$$

Preview

de Tèramond, Dosch, sjb



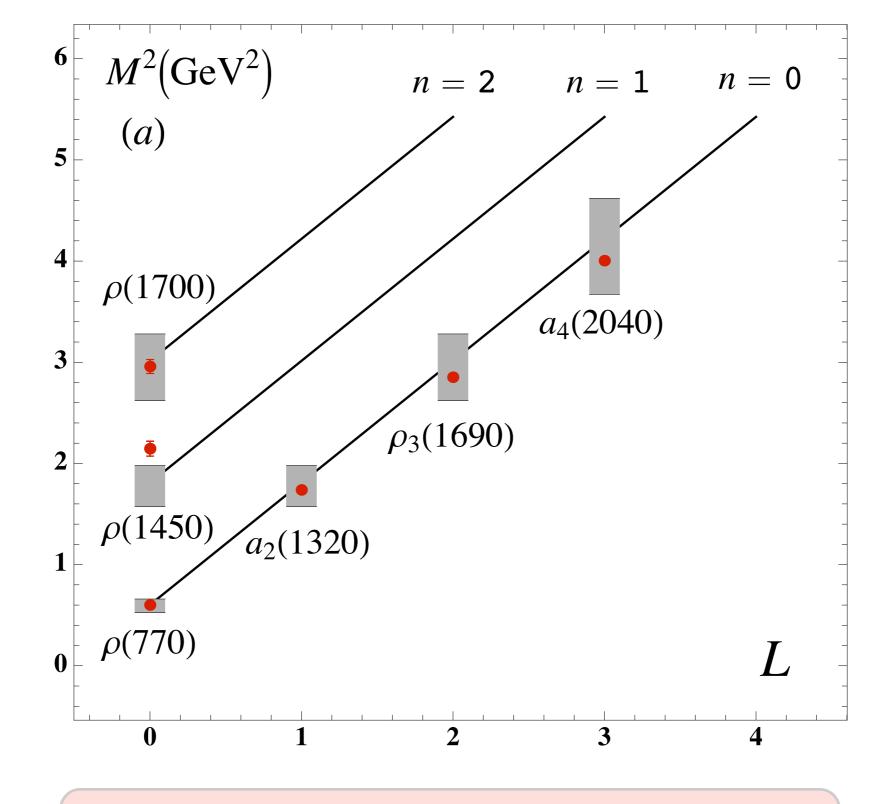
$$M^{2}(n, L, S) = 4\kappa^{2}(n + L + S/2)$$

QCD-TNT4
Ilhabela, Brazil
August 31, 2015

Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD

Stan Brodsky





$$M^{2}(n, L, S) = 4\kappa^{2}(n + L + S/2)$$

QCD-TNT4
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Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD

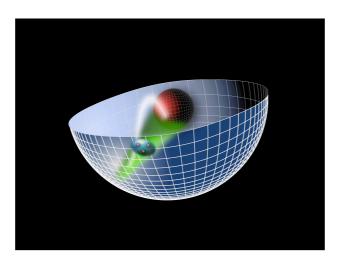
Stan Brodsky



de Tèramond, Dosch, sjb

AdS/QCD Soft-Wall Model

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$



$$\zeta^2 = x(1-x)\mathbf{b}_{\perp}^2.$$

Light-Front Holography

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

 $\kappa \simeq 0.6 \; GeV$

Confinement scale:

$$1/\kappa \simeq 1/3 \ fm$$

- de Alfaro, Fubini, Furlan:
 - Fubini, Rabinovici:

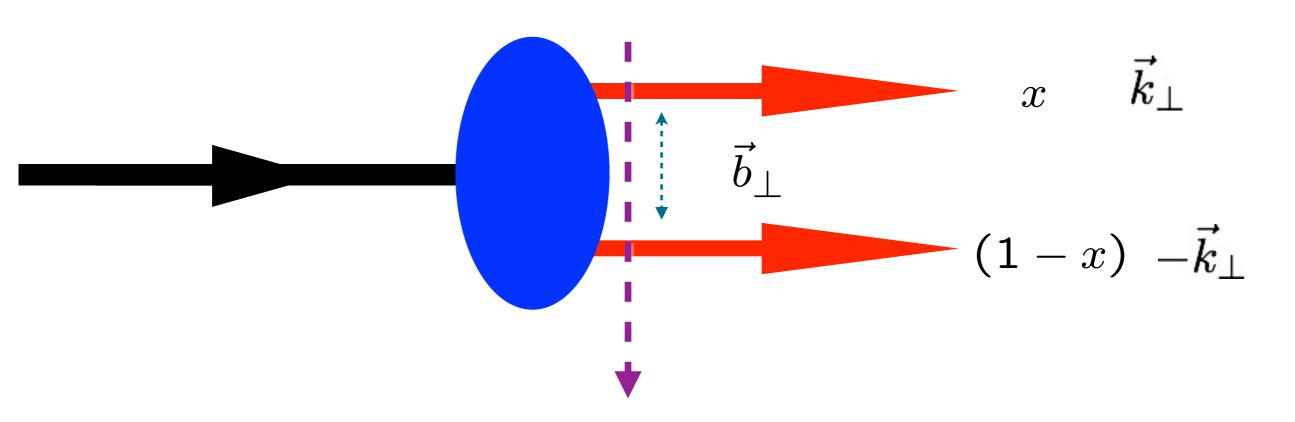
Unique Confinement Potential!

Scale can appear in Hamiltonian and EQM

without affecting conformal invariance of action!

Preserves Conformal Symmetry of the action

Fixed $\tau = t + z/c$



$$\zeta^2 \equiv b_\perp^2 x (1-x)$$

Invariant transverse separation

$$\zeta^2$$
 conjugate to $\frac{k_\perp^2}{x(1-x)}=(p_q+p_{\bar q})^2=\mathcal{M}_{q+\bar q}^2$

$$\int dk^- \Psi_{BS}(P,k) \to \psi_{LF}(x,\vec{k}_\perp)$$

LF Holography

Baryon Equation

Superconformal Algebra

$$\left(-\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}(L_{B} + 1) + \frac{4L_{B}^{2} - 1}{4\zeta^{2}}\right)\psi_{J}^{+} = M^{2}\psi_{J}^{+}$$

$$\left(-\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}L_{B} + \frac{4(L_{B} + 1)^{2} - 1}{4\zeta^{2}}\right)\psi_{J}^{-} = M^{2}\psi_{J}^{-}$$

$$M^2(n, L_B) = 4\kappa^2(n + L_B + 1)$$

S=1/2, P=+

Meson Equation

both chiralities

$$\left(-\partial_{\zeta}^{2} + \kappa^{4}\zeta^{2} + 2\kappa^{2}(J-1) + \frac{4L_{M}^{2}-1}{4\zeta^{2}}\right)\phi_{J} = M^{2}\phi_{J}$$

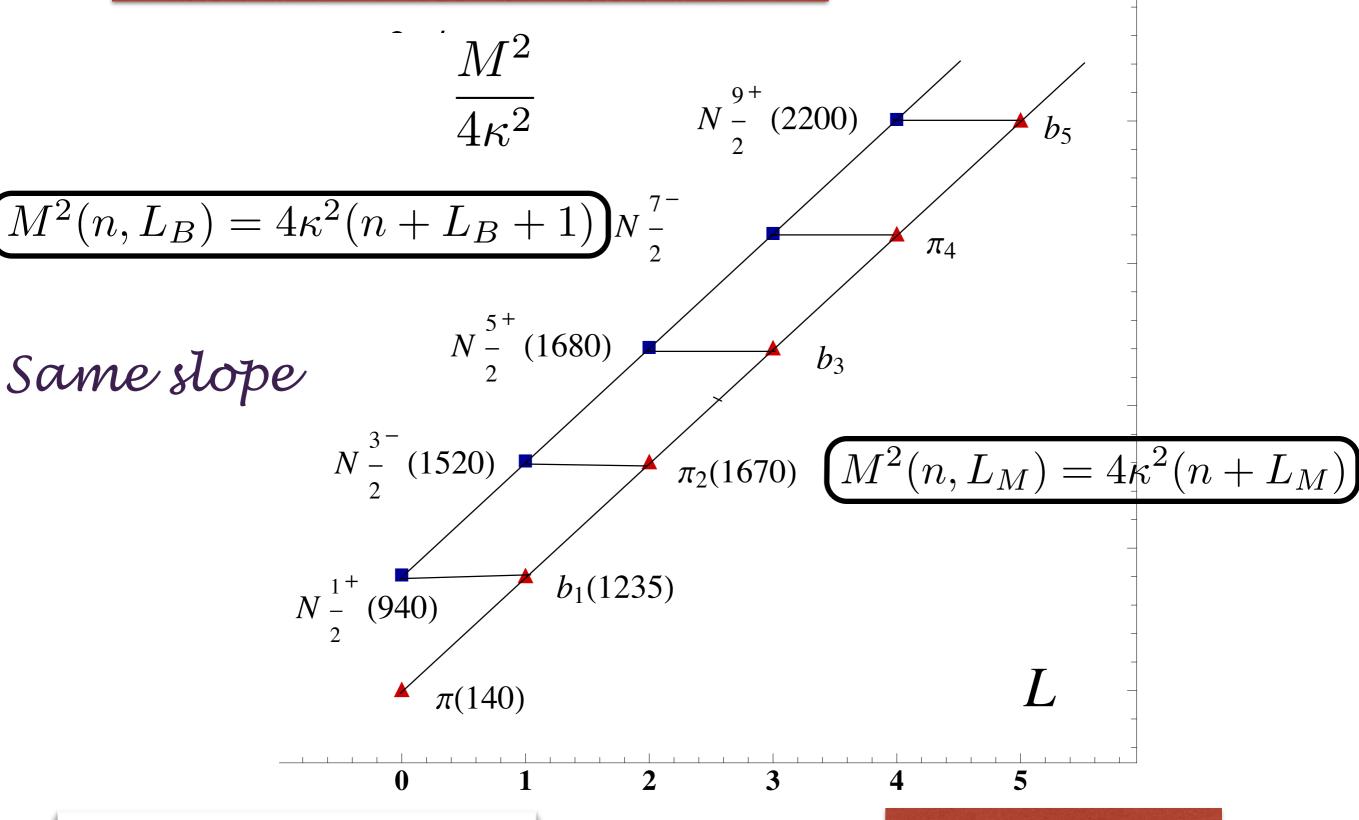
$$M^2(n, L_M) = 4\kappa^2(n + L_M)$$

Same k!

S=0, I=I Meson is superpartner of S=I/2, I=I Baryon Meson-Baryon Degeneracy for $L_M=L_B+1$

Superconformal Algebra

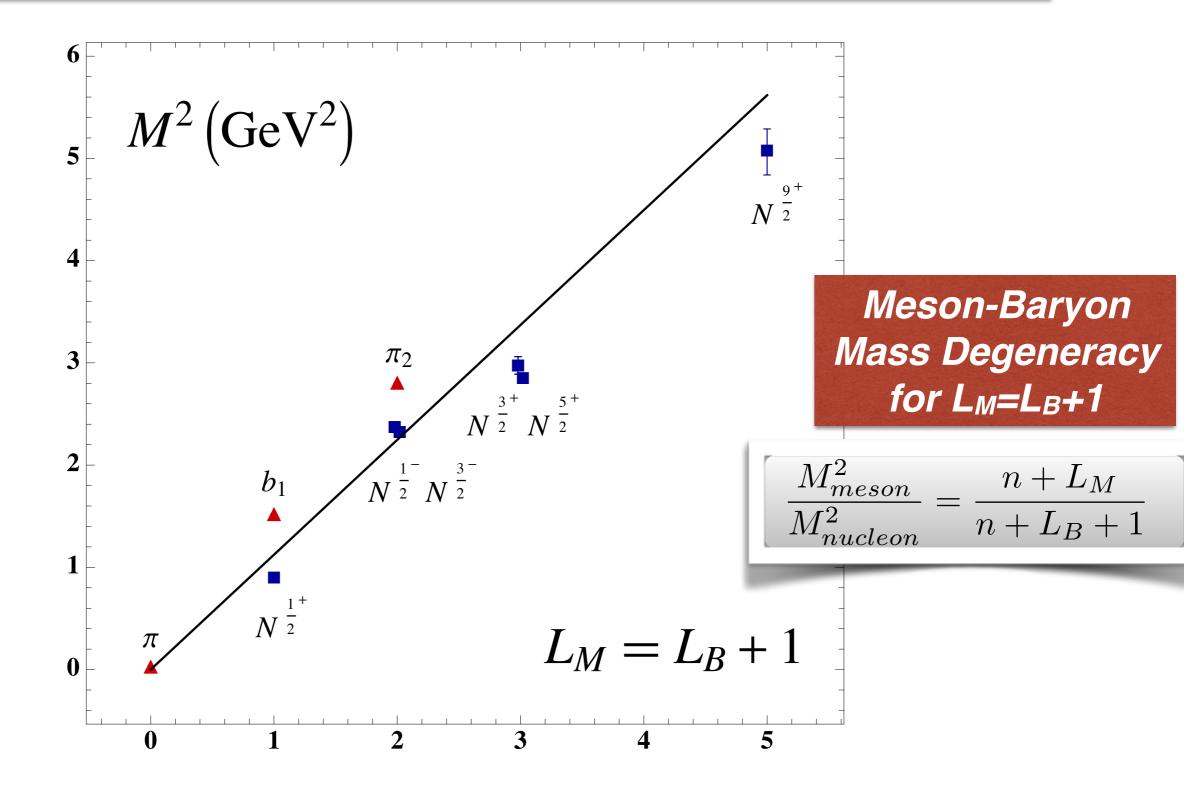
de Tèramond, Dosch, sjb



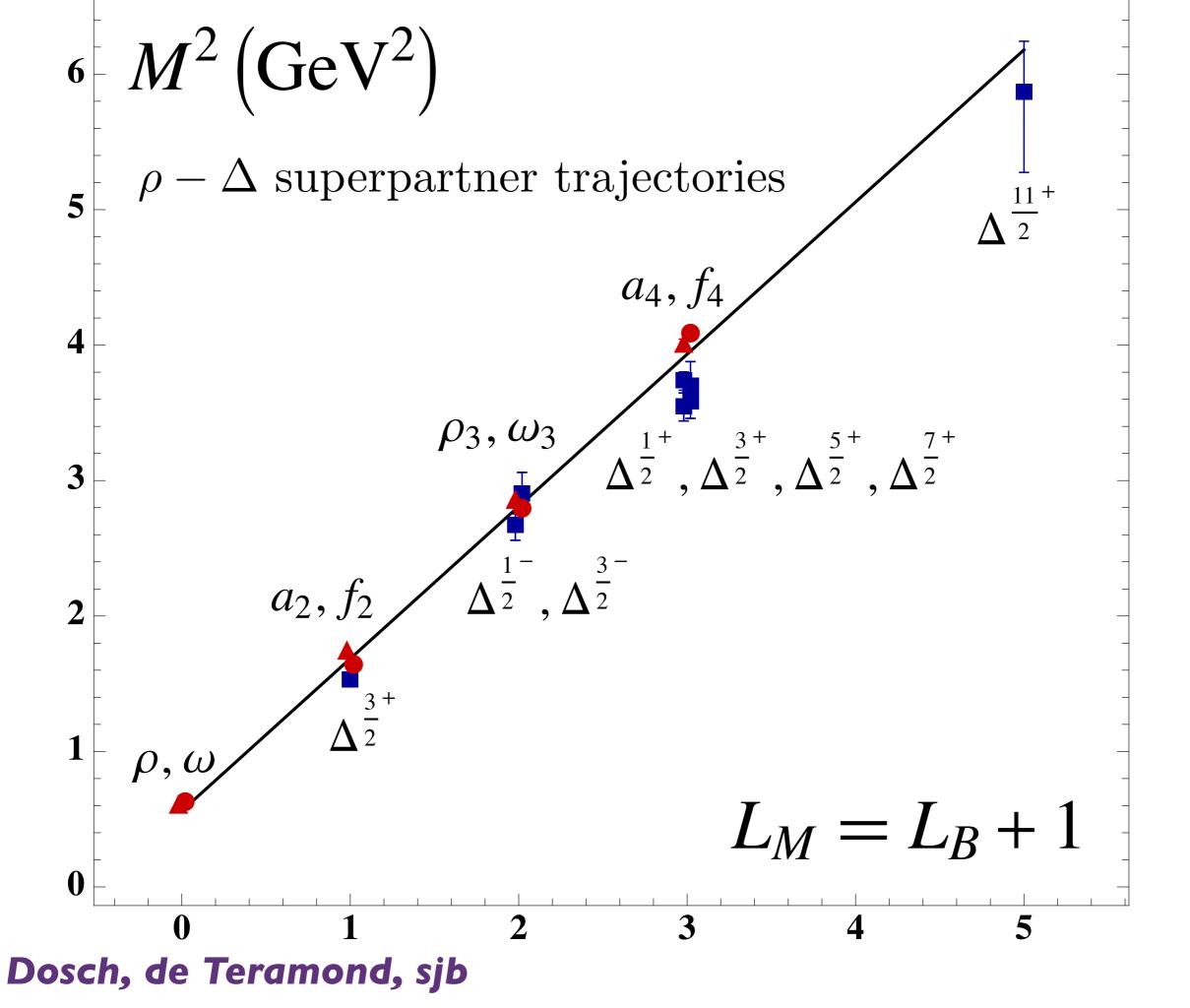
$$\frac{M_{meson}^2}{M_{nucleon}^2} = \frac{n + L_M}{n + L_B + 1}$$

Meson-Baryon
Mass Degeneracy
for L_M=L_B+1

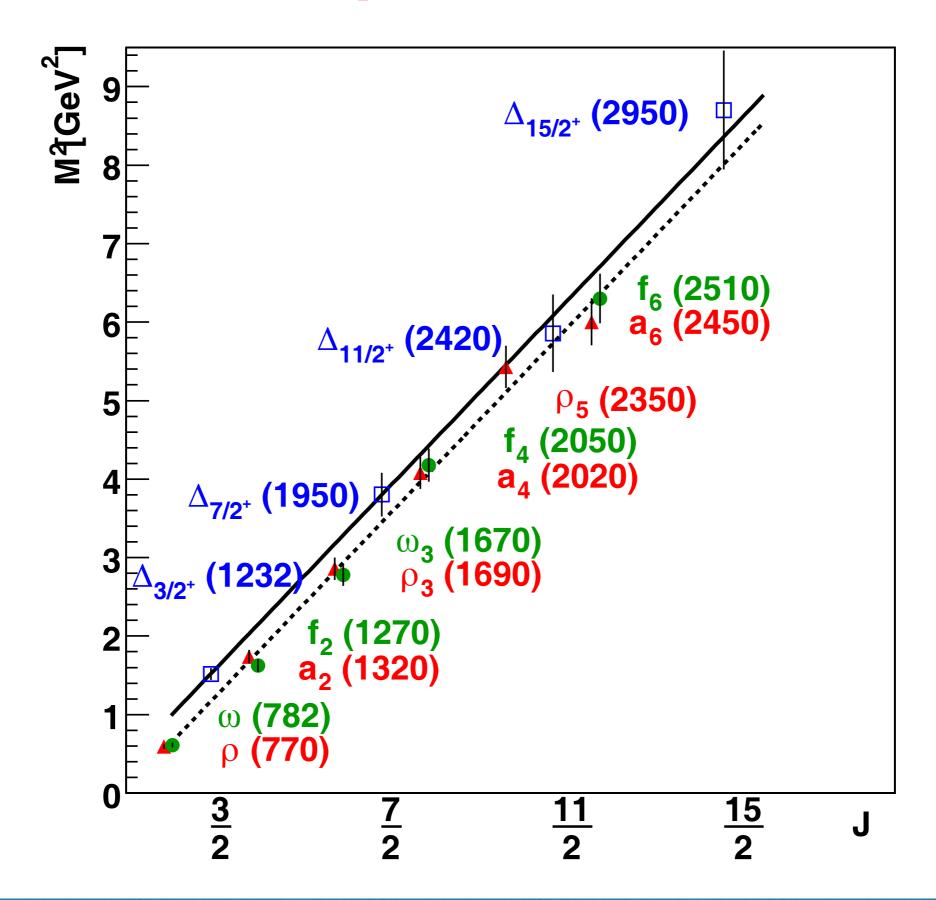
Superconformal AdS Light-Front Holographic QCD (LFHQCD): Identical meson and baryon spectra!



S=0, I=1 Meson is superpartner of S=1/2, I=1 Baryon



E. Klempt and B. Ch. Metsch



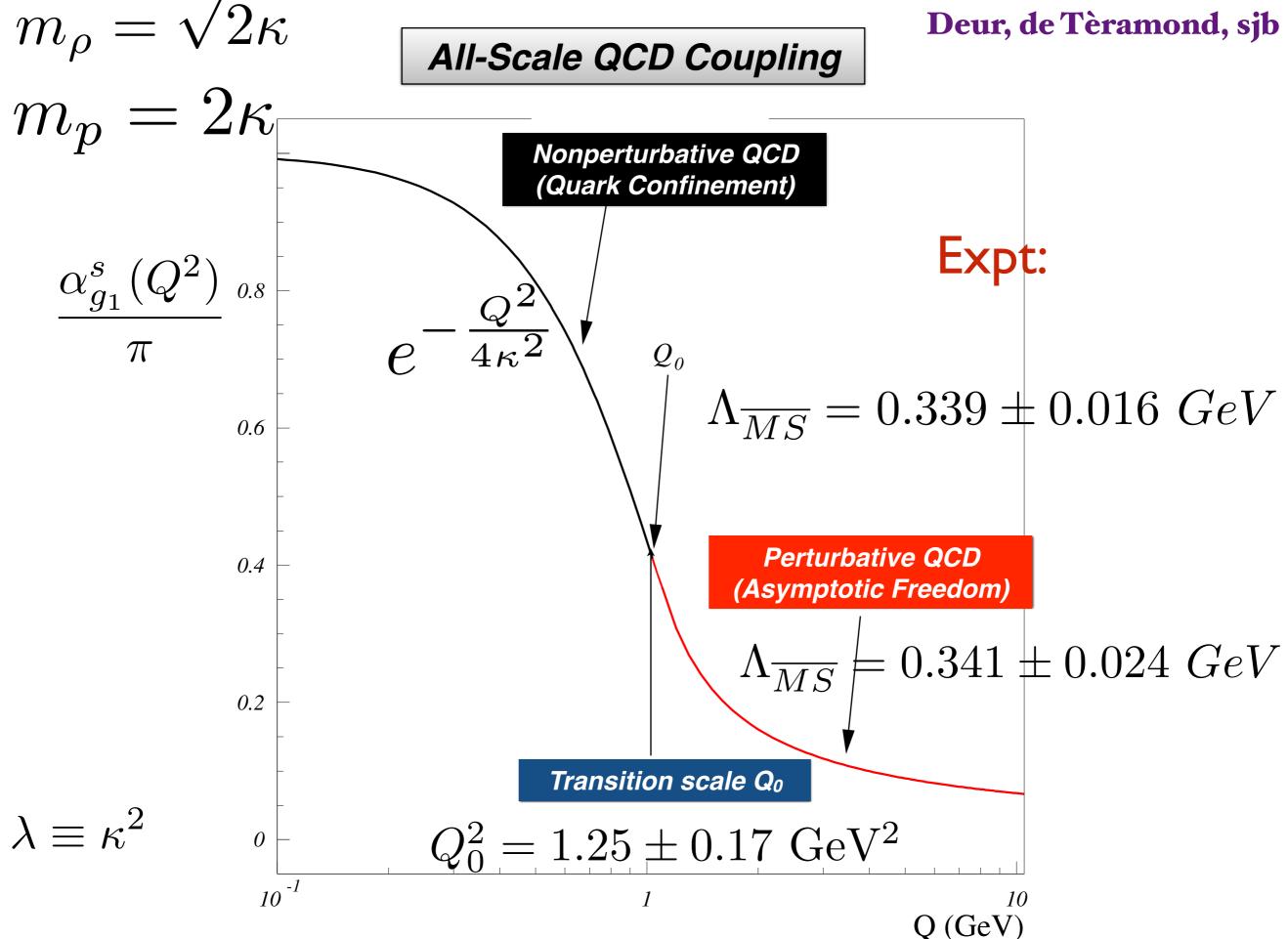
Some Features of AdS/QCD

- Regge spectroscopy—same slope in n,L for mesons,
- Chiral features for m_q =0: m_{π} =0, chiral-invariant proton
- Hadronic LFWFs
- Counting Rules
- ullet Connection between hadron masses and $\Lambda_{\overline{MS}}$

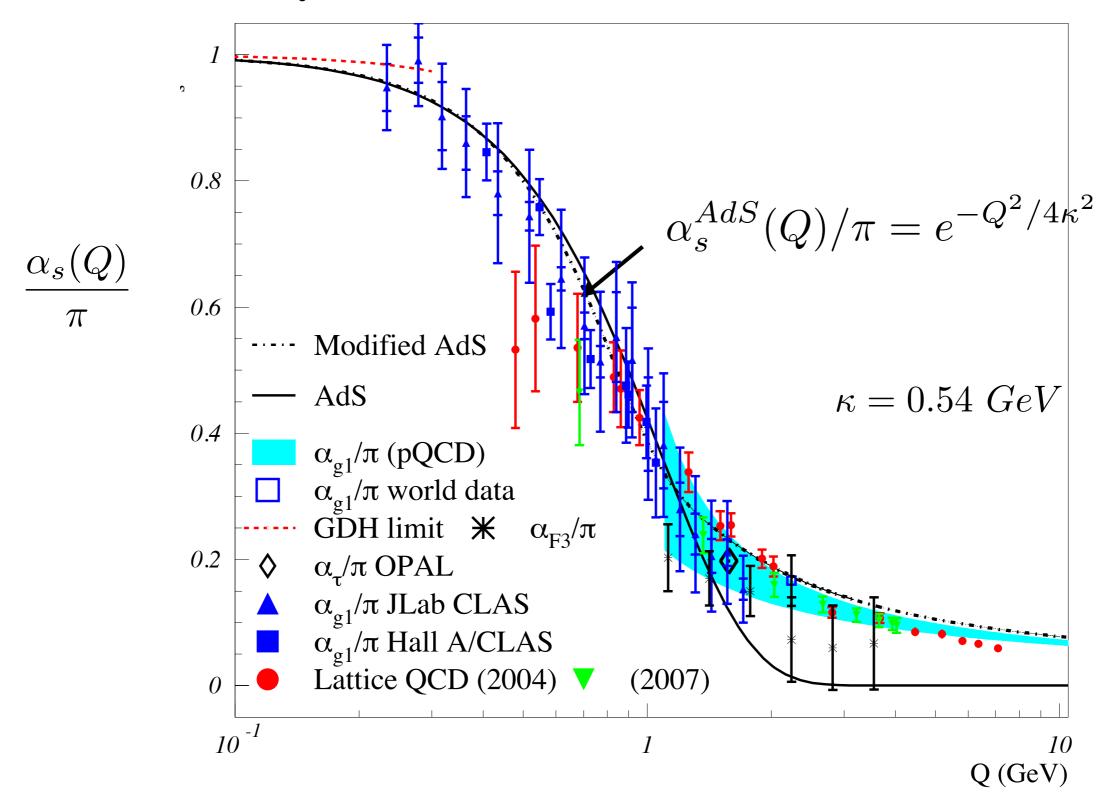
Superconformal AdS Light-Front Holographic QCD (LFHQCD)

Meson-Baryon Mass Degeneracy for L_M=L_B+1





Analytic, defined at all scales, IR Fixed Point



AdS/QCD dilaton captures the higher twist corrections to effective charges for Q < 1 GeV

$$e^{\varphi} = e^{+\kappa^2 z^2}$$

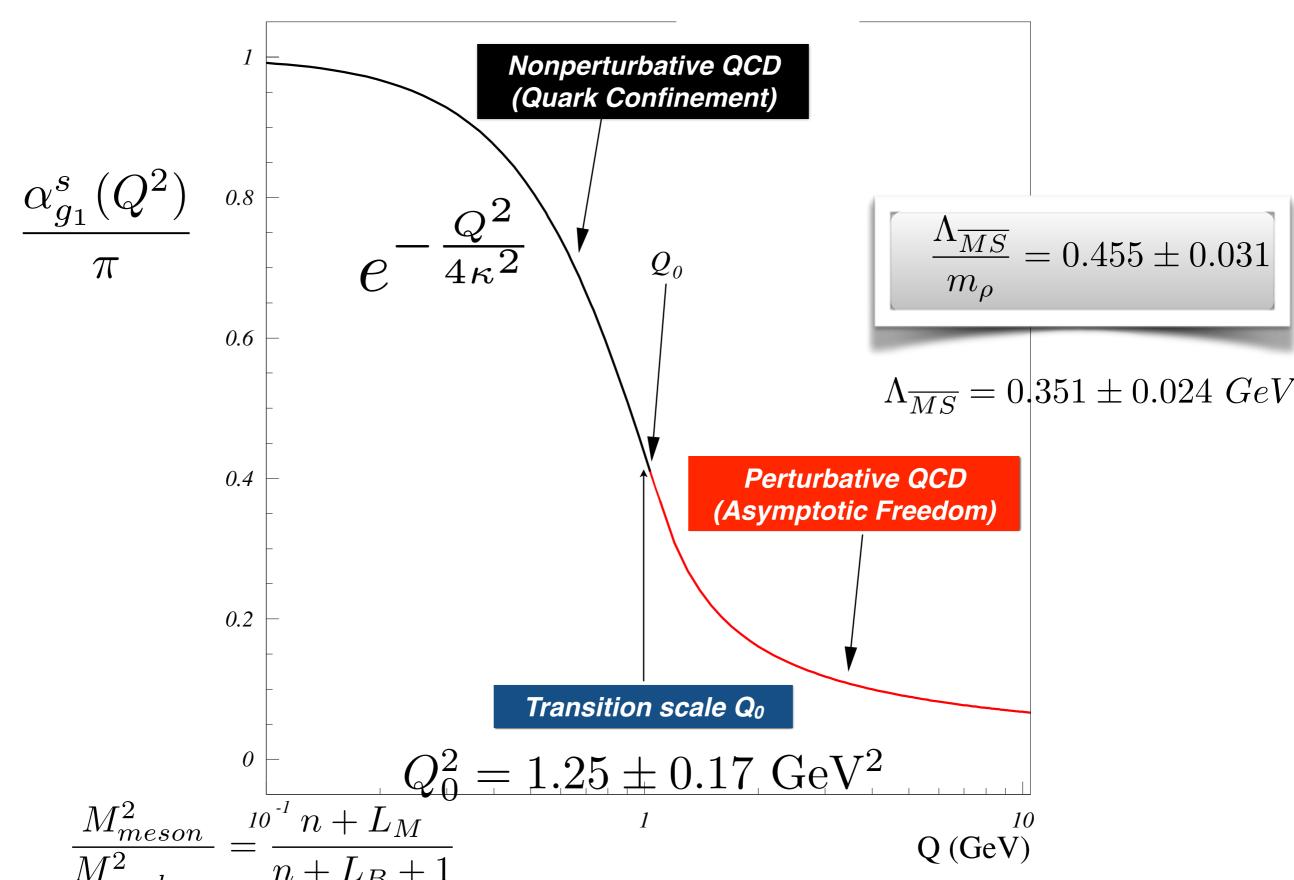
Deur, de Teramond, sjb

$$m_{\rho} = \sqrt{2}\kappa$$

All-Scale QCD Coupling

Deur, de Teramond, sjb

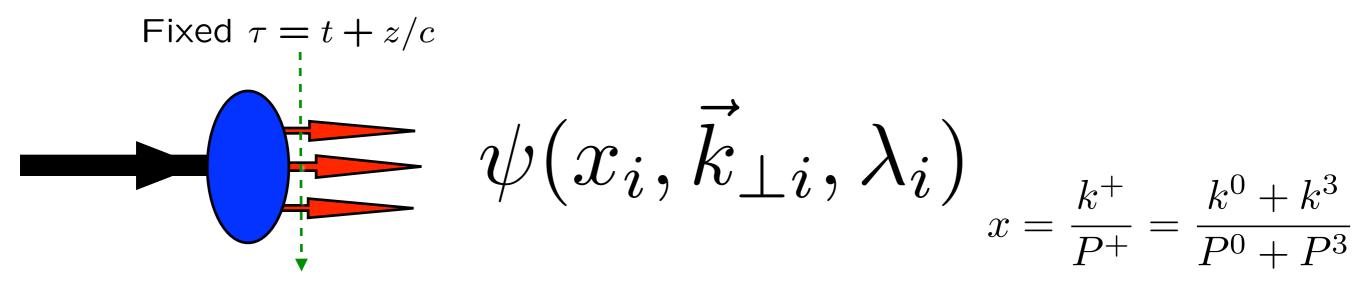
Prediction from AdS/QCD:



Bound States in Relativistic Quantum Field Theory:

Light-Front Wavefunctions

Dirac's Front Form: Fixed $\tau = t + z/c$



Invariant under boosts. Independent of P^{μ}

$$H_{LF}^{QCD}|\psi>=M^2|\psi>$$

Direct connection to QCD Lagrangian

Off-shell in invariant mass

Remarkable new insights from AdS/CFT, the duality between conformal field theory and Anti-de Sitter Space

Each element of flash photograph illuminated at same LF time

$$\tau = t + z/c$$

Causal, frame-independent

Evolve in LF time

$$P^- = i \frac{d}{d\tau}$$

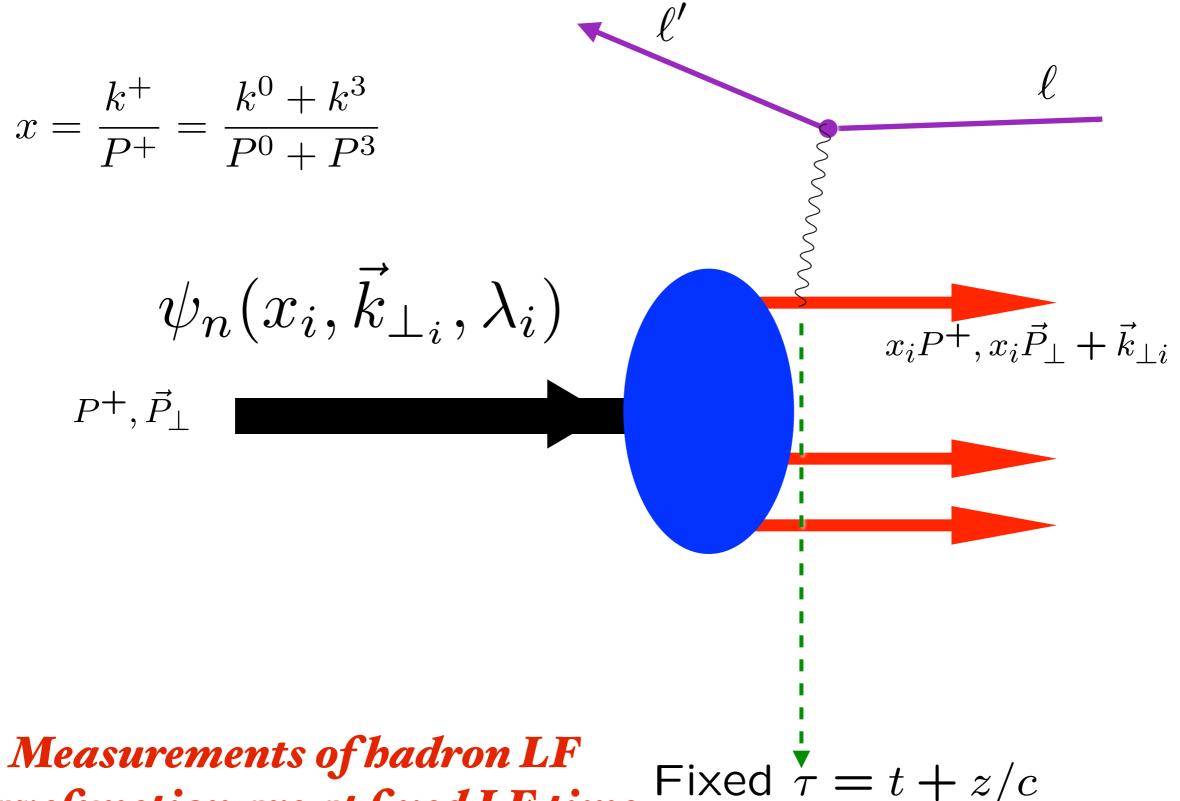
Eigenstate -- independent of au

$$H_{LF} = P^+P^- - \vec{P}_{\perp}^2$$

$$H_{LF}^{QCD}|\Psi_h> = \mathcal{M}_h^2|\Psi_h>$$



HELEN BRADLEY - PHOTOGRAPHY



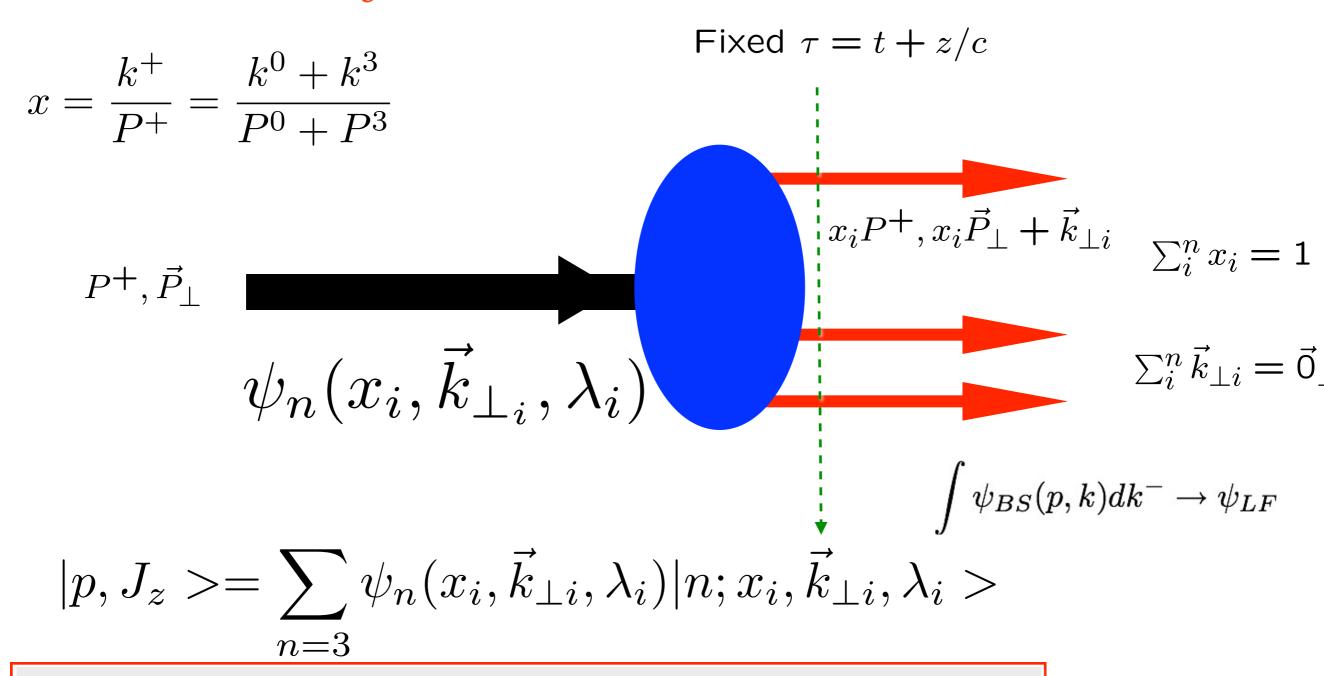
wavefunction are at fixed LF time

Like a flash photograph

$$x_{bj} = x = \frac{k^+}{P^+}$$

Light-Front Wavefunctions: rigorous representation of composite systems in quantum field theory

Eigenstate of LF Hamiltonian



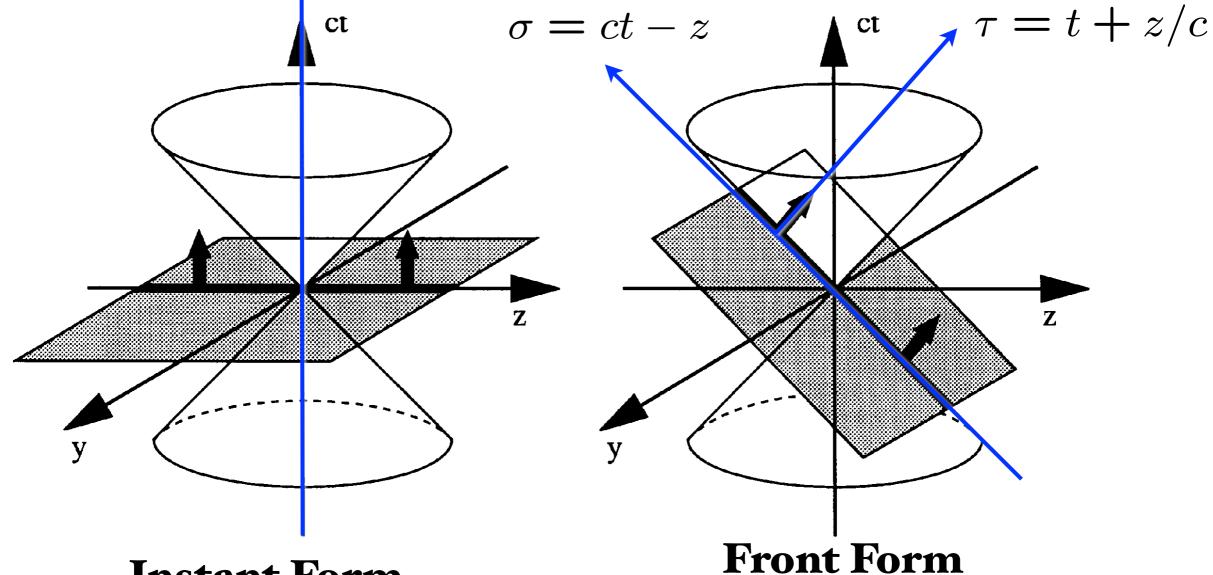
Invariant under boosts! Independent of P^{μ}

Causal, Frame-independent. Creation Operators on Simple Vacuum, Current Matrix Elements are Overlaps of LFWFS

Dirac's Amazing Idea: The "Front Form"

P.A.M Dirac, Rev. Mod. Phys. 21, 392 (1949)

Evolve in Evolve in ordinary time light-front time!



Instant Form

Boost Invariant!

$$= 2p^{+}F(q^{2}) \qquad \text{Interaction}$$

$$q_{\perp}^{2} = Q^{2} = -q^{2} \qquad \qquad \gamma^{*} \qquad \text{Fixed } \tau = t + z/c$$

$$q^{+} = 0 \quad \vec{q}_{\perp} \qquad \qquad \qquad \qquad \text{Form Factors are}$$

$$Q \text{Verlaps of LFWFs}$$

$$x, \vec{k}_{\perp} \qquad \qquad x, \vec{k}_{\perp} + \vec{q}_{\perp} \qquad \qquad p + q$$

$$\psi(x_{i}, \vec{k}_{\perp i}) \qquad \qquad \psi(x_{i}, \vec{k}_{\perp i}')$$

$$\text{struck} \qquad \vec{k}_{\perp i}' = \vec{k}_{\perp i} + (1 - x_{i})\vec{q}_{\perp}$$

$$\text{Drell \&Yan, West}$$

$$\text{Exact LF formula!} \qquad \qquad \text{spectators} \qquad \vec{k}_{\perp i}' = \vec{k}_{\perp i} - x_{i}\vec{q}_{\perp}$$

No comparable formula in instant form

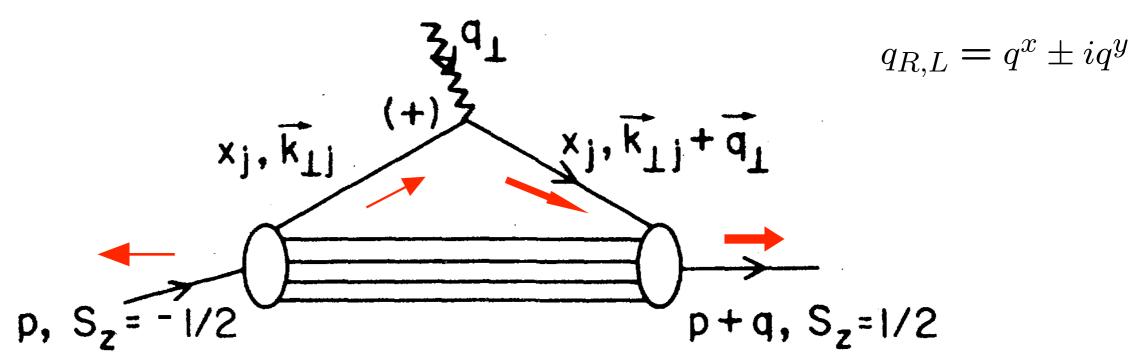
Exact LF Formula for Pauli Form Factor

$$\frac{F_{2}(q^{2})}{2M} = \sum_{a} \int [\mathrm{d}x][\mathrm{d}^{2}\mathbf{k}_{\perp}] \sum_{j} e_{j} \frac{1}{2} \times$$

$$\left[-\frac{1}{q^{L}} \psi_{a}^{\uparrow *}(x_{i}, \mathbf{k}'_{\perp i}, \lambda_{i}) \psi_{a}^{\downarrow}(x_{i}, \mathbf{k}_{\perp i}, \lambda_{i}) + \frac{1}{q^{R}} \psi_{a}^{\downarrow *}(x_{i}, \mathbf{k}'_{\perp i}, \lambda_{i}) \psi_{a}^{\uparrow}(x_{i}, \mathbf{k}_{\perp i}, \lambda_{i}) \right]$$

$$\mathbf{k}'_{\perp i} = \mathbf{k}_{\perp i} - x_{i} \mathbf{q}_{\perp}$$

$$\mathbf{k}'_{\perp j} = \mathbf{k}_{\perp j} + (1 - x_{j}) \mathbf{q}_{\perp}$$
Drell, sjb
$$\mathbf{k}'_{\perp i} = \mathbf{k}_{\perp i} + x_{i} \mathbf{q}_{\perp}$$



Must have $\Delta \ell_z = \pm 1$ to have nonzero $F_2(q^2)$

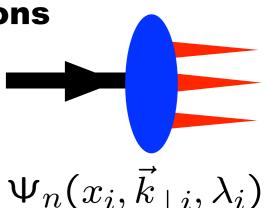
Nonzero Proton Anomalous Moment -->
Nonzero orbital quark angular momentum

Advantages of the Dirac's Front Form for Hadron Physics

- Measurements are made at fixed τ
- Causality is automatic
- Structure Functions are squares of LFWFs
- Form Factors are overlap of LFWFs
- LFWFs are frame-independent: no boosts, no pancakes!
- Same structure function in e p collider and p rest frame
- No dependence on observer's frame
- LF Holography: Dual to AdS space
- LF Vacuum trivial -- no condensates!
- Profound implications for Cosmological Constant



- LF wavefunctions play the role of Schrödinger wavefunctions in Atomic Physics
- LFWFs=Hadron Eigensolutions: Direct Connection to QCD Lagrangian

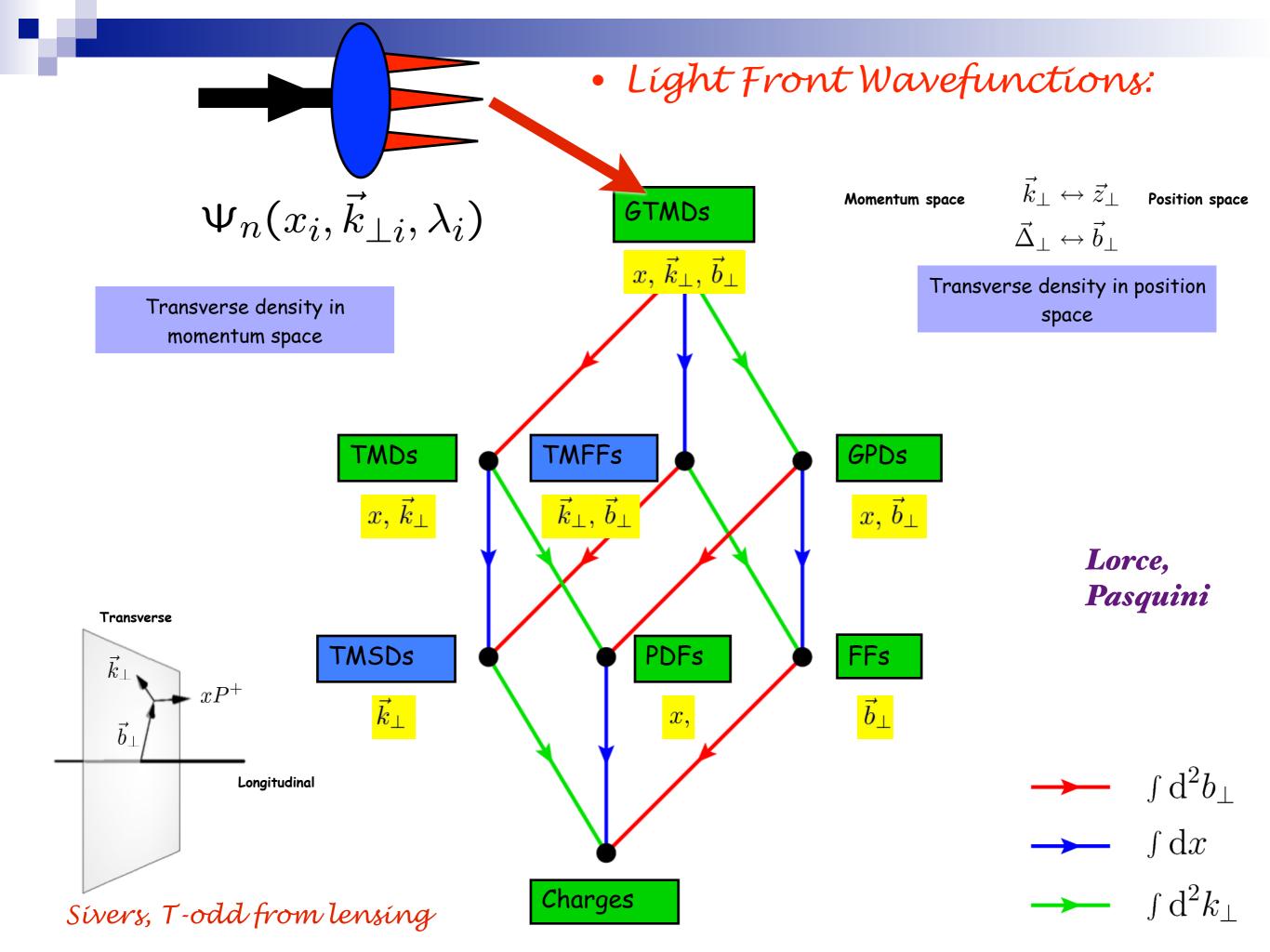


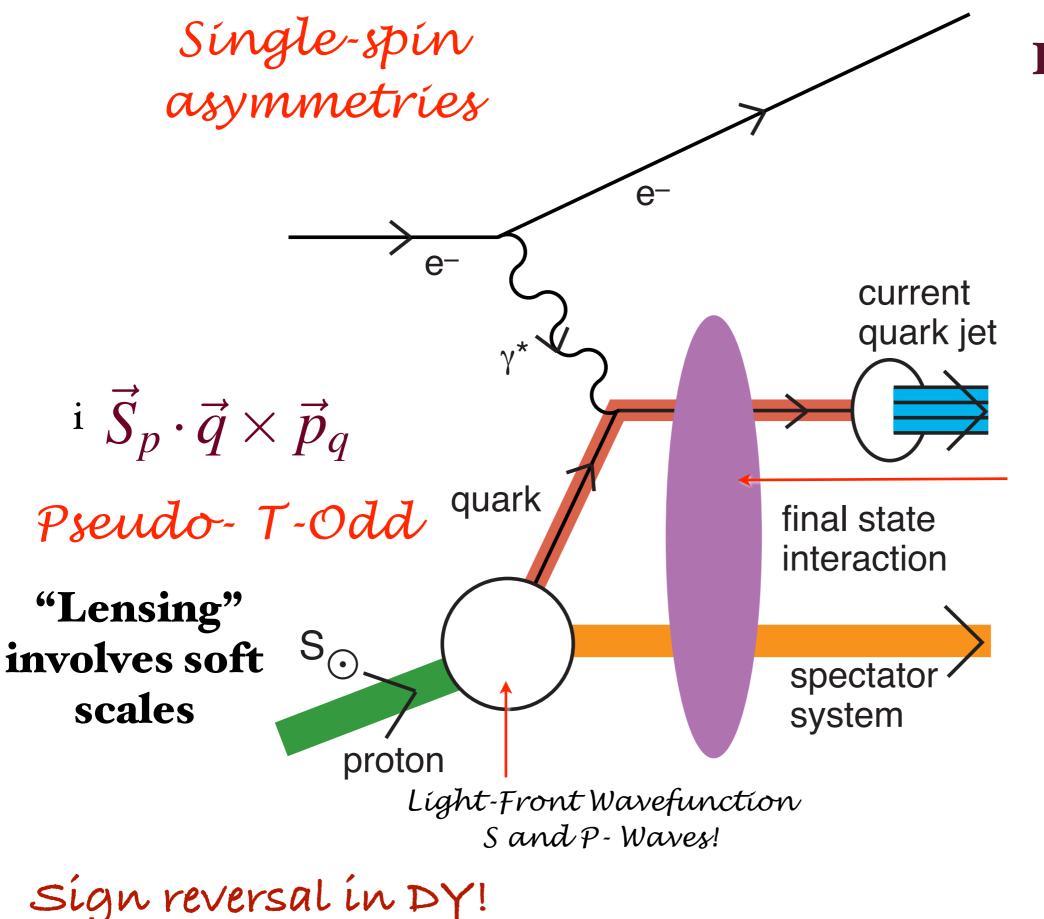
- Relativistic, frame-independent: no boosts, no disc contraction, Melosh built into LF spinors
- Hadronic observables computed from LFWFs: Form factors,
 Structure Functions, Distribution Amplitudes, GPDs, TMDs,
 Weak Decays, modulo `lensing' from ISIs, FSIs
- Cannot compute current matrix elements using instant form from eigensolutions alone -- need to include vacuum currents!
- Hadron Physics without LFWFs is like Biology without DNA!

Stan Brodsky

SLAC

NATIONAL ACCELERATOR LABORATORY





Leading Twist Sivers Effect

Hwang, Schmidt, sjb

Collins, Burkardt, Ji, Yuan. Pasquini, ...

QCD S- and P-Coulomb Phases --Wilson Line

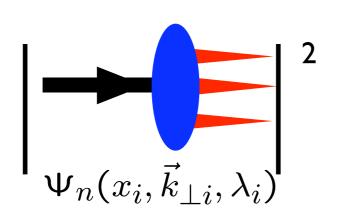
"Lensing Effect"

Leading-Twist Rescattering Violates pQCD Factorization!

Static

Dynamic

- Square of Target LFWFs
- No Wilson Line
- Probability Distributions
- Process-Independent
- T-even Observables
- No Shadowing, Anti-Shadowing
- Sum Rules: Momentum and J^z
- DGLAP Evolution; mod. at large x
- No Diffractive DIS



Modified by Rescattering: ISI & FSI

Contains Wilson Line, Phases

No Probabilistic Interpretation

Process-Dependent - From Collision

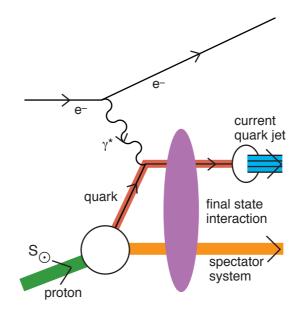
T-Odd (Sivers, Boer-Mulders, etc.)

Shadowing, Anti-Shadowing, Saturation

Sum Rules Not Proven

DGLAP Evolution

Hard Pomeron and Odderon Diffractive DIS



Hwang, Schmidt, sjb,

Mulders, Boer

Qiu, Sterman

Collins, Qiu

Pasquini, Xiao, Yuan, sjb







Goal: An analytic first approximation to QCD

- As Simple as Schrödinger Theory in Atomic Physics
- Relativistic, Frame-Independent, Color-Confining
- Confinement in QCD -- What is the analytic form of the confining interaction?
- What sets the QCD mass scale?
- QCD Running Coupling at all scales
- Hadron Spectroscopy-Regge Trajectories
- Light-Front Wavefunctions
- Form Factors, Structure Functions, Hadronic Observables
- Constituent Counting Rules
- Hadronization at the Amplitude Level
- Insights into QCD Condensates
- Chiral Symmetry
- Systematically improvable

QCD-TNT4
Ilhabela, Brazil
August 31, 2015

Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD

Stan Brodsky



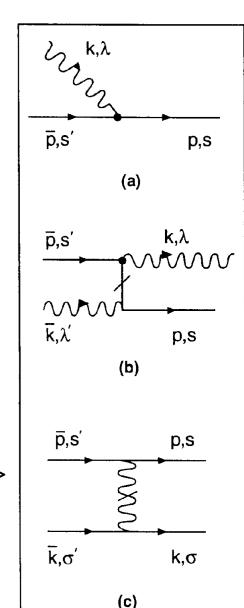
Light-Front QCD

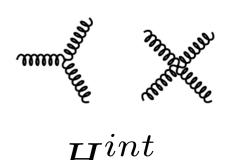
Exact frame-independent formulation of nonperturbative QCD!

$$L^{QCD}
ightarrow H^{QCD}_{LF}$$
 $H^{QCD}_{LF} = \sum_{i} [\frac{m^2 + k_{\perp}^2}{x}]_i + H^{int}_{LF}$
 H^{int}_{LF} : Matrix in Fock Space
 $H^{QCD}_{LF} |\Psi_h> = \mathcal{M}^2_h |\Psi_h>$
 $|p,J_z> = \sum_{n=3} \psi_n(x_i,\vec{k}_{\perp i},\lambda_i)|n;x_i,\vec{k}_{\perp i},\lambda_i>$

Eigenvalues and Eigensolutions give Hadronic Spectrum and Light-Front wavefunctions

LFWFs: Off-shell in P- and invariant mass





Light-Front QCD Heisenberg Equation

$$H_{LC}^{QCD}|\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

DLCQ: Solve QCD(1+1) for any quark mass and flavors

Hornbostel, Pauli, sjb

Ly k, i	
p̄,s′	p,s
(a)
p ,s′	<u>k</u> ,λ
~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	p,s
(1	b)
p̄,s′	p,s
	٠ ٢
¯k,σ′	k,σ
(0	c)

n	Sector	1 q <del>q</del>	2 99	3 q <del>q</del> g	4 q <b>q</b> q <b>q</b>	5 gg g	6 qq gg	7 qq qq g	8 वव वव वव	9 99 99	10 qq gg g	11 वव वव gg	12 qq qq qq g	13 qq qq qq qq
1	qq		+	~	X	•		•	•	•	•	•	•	•
2	<b>gg</b>			~	•	~~~~~		•	•	}-\\\\\\\\\\\\\\\\\\\\\\\\\\\\\\\\\\\\	•	•	•	•
3	qq g	<b>&gt;</b>	<b>&gt;</b>	<u> </u>	~<		~~~~		•	•	+	•	•	•
4	qq qq	<u> </u>	•	<b>&gt;</b>		•		~	XIII	•	•	1	•	•
5	gg g	•	<i>&gt;</i>		•	X	~~<	•	•	~~~~~		•	•	•
6	qq gg	V4. {	<b>*</b>	<b>&gt;</b>		>		~<	•		-<		•	•
7	वव वव व	•	•	***	<b>&gt;-</b>	•	>	+	~~<	•		-<		•
8	qq qq qq	•	•	•	\	•	•	<b>&gt;</b>		•	•		-<	X
9	gg gg	•	\frac{1}{2}	•	•	<i>&gt;</i>		•	•	7/	~-<	•	•	•
10	qq gg g	•	•	7	•	5	>-		•	<b>&gt;</b>		~<	•	•
11	qq qq gg	•	•	•		•	>	<b>&gt;</b>		•	>		~<	•
12	ववे ववे ववे व	•	•	•	•	•	•	<b>**</b>	<b>&gt;</b>	•	•	>		~~
13 q	iā dā dā dā	•	•	•	•	•	•	•	<b>&gt;</b>	•	•	•	<b>&gt;</b>	7

Mínkowskí space; frame-independent; no fermion doubling; no ghosts trivial vacuum

$$|p,S_z>=\sum_{n=3}\Psi_n(x_i,\vec{k}_{\perp i},\lambda_i)|n;\vec{k}_{\perp i},\lambda_i>$$

### sum over states with n=3, 4, ... constituents

The Light Front Fock State Wavefunctions

$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

are boost invariant; they are independent of the hadron's energy and momentum  $P^{\mu}$ .

The light-cone momentum fraction

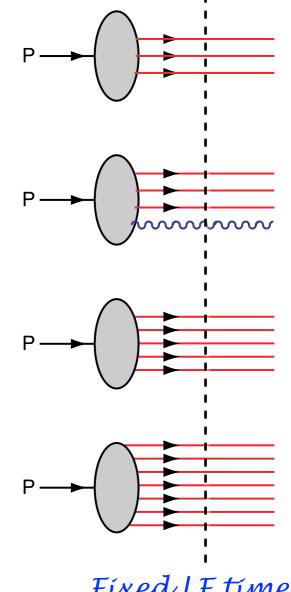
$$x_i = \frac{k_i^+}{p^+} = \frac{k_i^0 + k_i^z}{P^0 + P^z}$$

are boost invariant.

$$\sum_{i}^{n} k_{i}^{+} = P^{+}, \ \sum_{i}^{n} x_{i} = 1, \ \sum_{i}^{n} \vec{k}_{i}^{\perp} = \vec{0}^{\perp}.$$

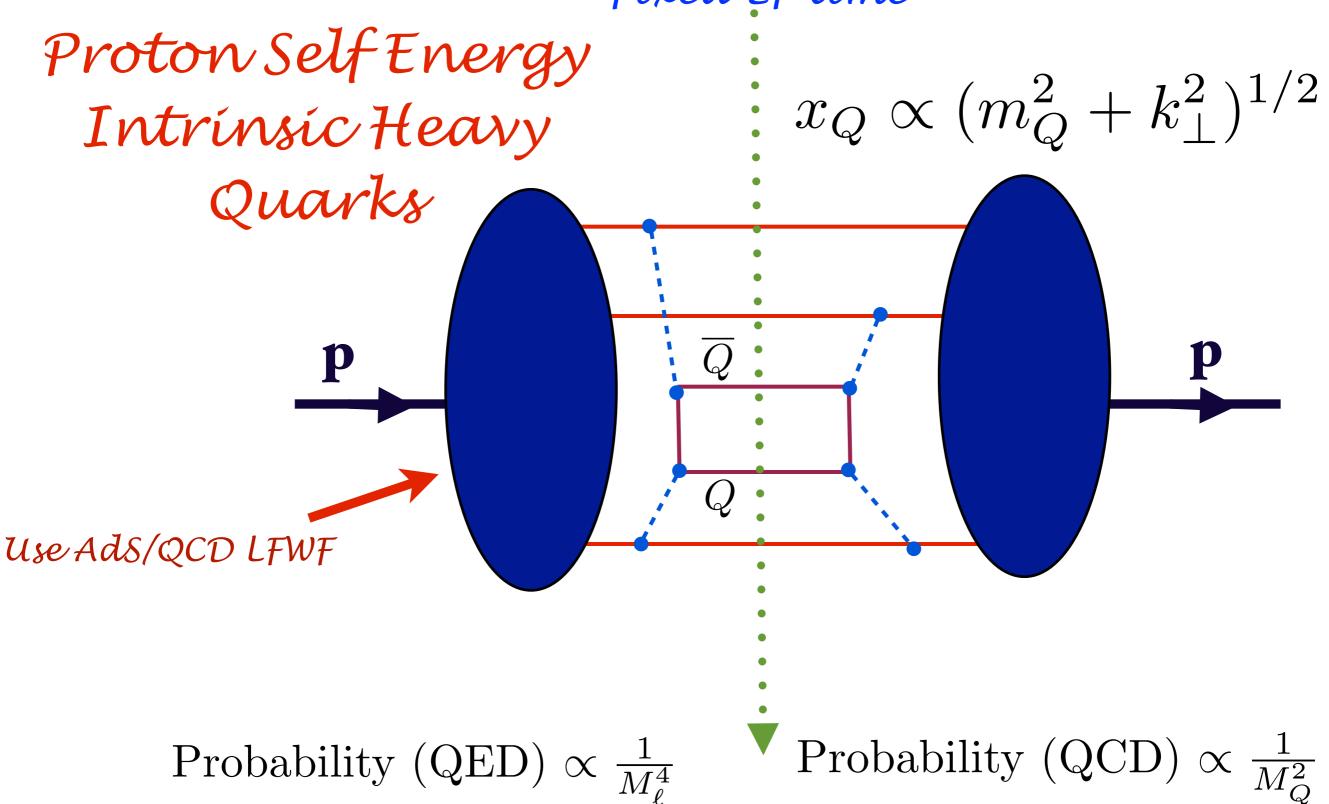
Intrinsic heavy quarks  $\bar{s}(x) \neq s(x)$   $\bar{s}(x) \neq \bar{s}(x) \neq \bar{d}(x)$   $\bar{u}(x) \neq \bar{d}(x)$ 

$$\bar{s}(x) \neq s(x)$$
 $\bar{u}(x) \neq \bar{d}(x)$ 



Fixed LF time

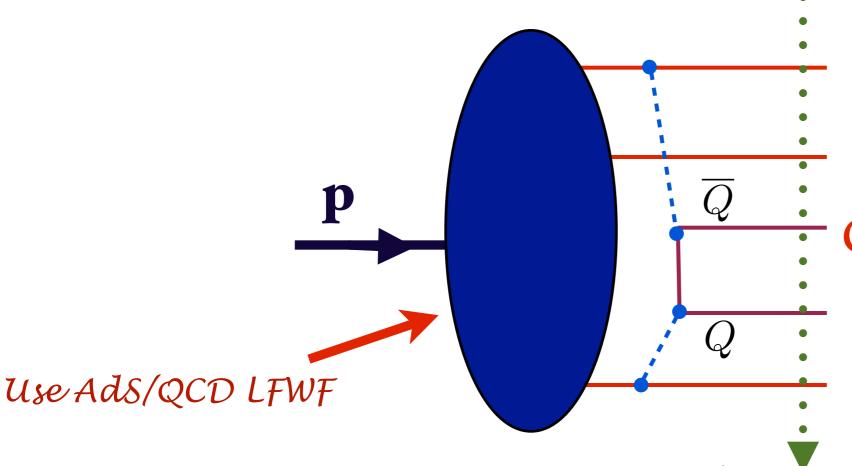
Fixed LF time



Collins, Ellis, Gunion, Mueller, sjb Polyakov, et al.

#### Fixed LF time





QCD predicts
Intrinsic Heavy
Quarks at high x

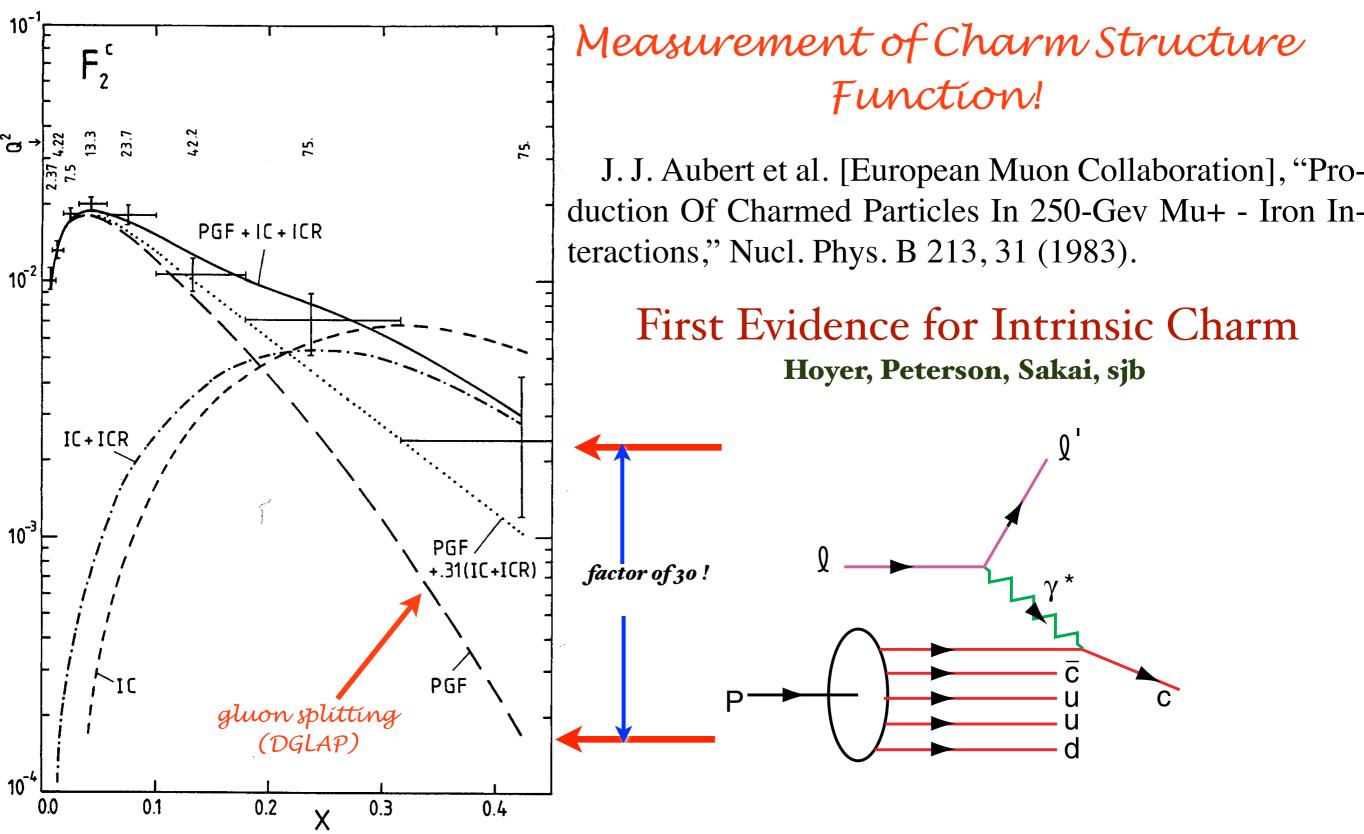
# Minimal offshellness

$$x_Q \propto (m_Q^2 + k_\perp^2)^{1/2}$$

Probability (QED) 
$$\propto \frac{1}{M_{\ell}^4}$$

Probability (QCD) 
$$\propto \frac{1}{M_Q^2}$$

Collins, Ellis, Gunion, Mueller, sjb Polyakov, et al.



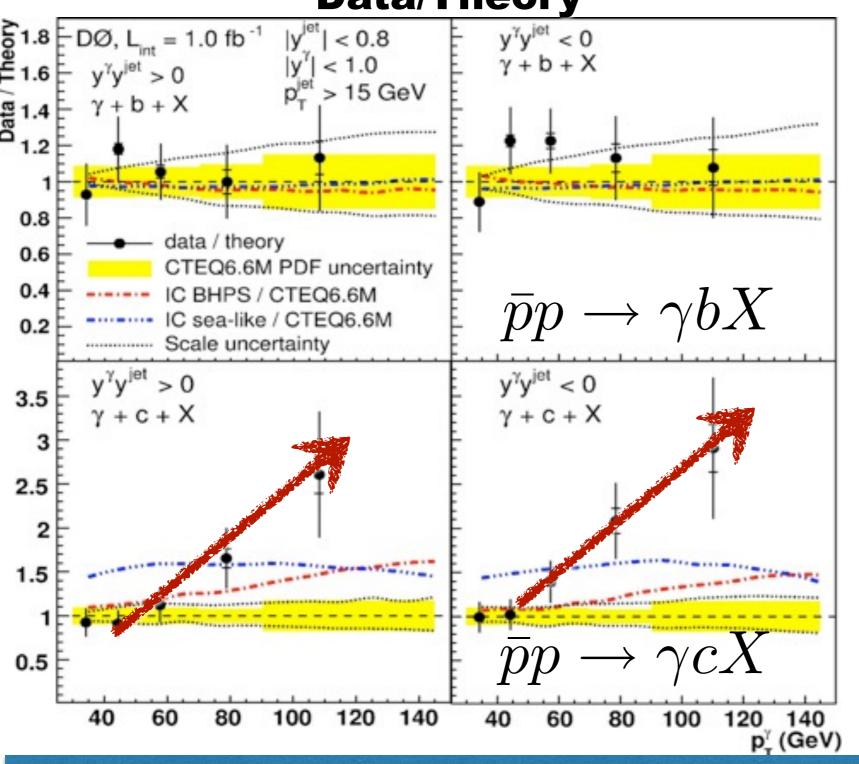
## DGLAP / Photon-Gluon Fusion: factor of 30 too small

Two Components (separate evolution):

$$c(x, Q^2) = c(x, Q^2)_{\text{extrinsic}} + c(x, Q^2)_{\text{intrinsic}}$$

Measurement of  $\gamma + b + X$  and  $\gamma + c + X$  Production Cross Sections in  $p\bar{p}$  Collisions at  $\sqrt{s} = 1.96$  TeV

**Data/Theory** 



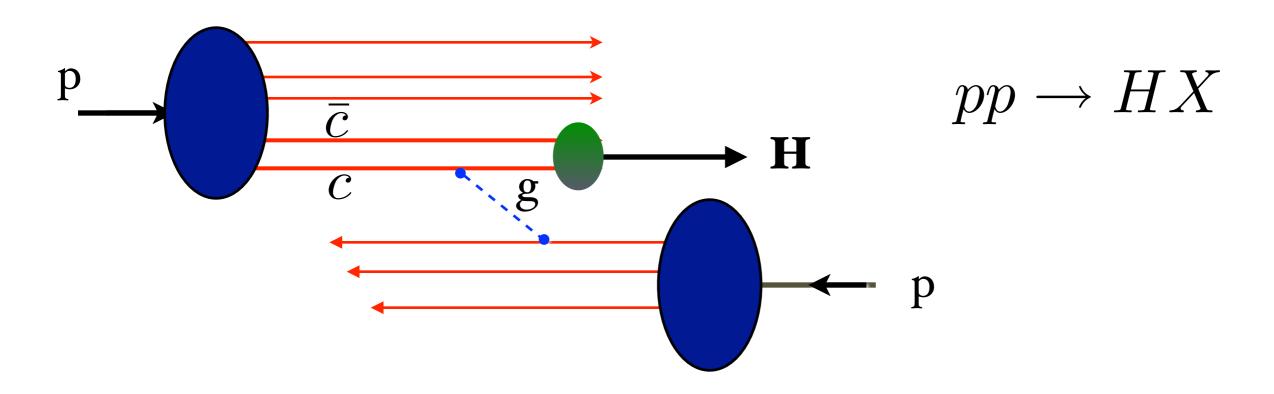
$$\frac{\Delta\sigma(\bar{p}p\to\gamma cX)}{\Delta\sigma(\bar{p}p\to\gamma bX)}$$

Ratio insensitive to gluon PDF, scales

Signal for significant IC at x > 0.1

Consistent with EMC measurement of charm structure function at high x

# Intrinsic Charm Mechanism for Inclusive $High-X_F$ Higgs Production



Also: intrinsic strangeness, bottom, top

Higgs can have > 80% of Proton Momentum!

New production mechanism for Higgs

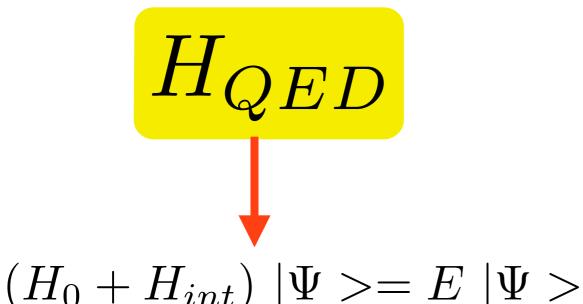
AFTER: Higgs production at threshold!

JLab: Charm production near threshold!

## Need a First Approximation to QCD

# Comparable in simplicity to Schrödinger Theory in Atomic Physics

Relativistic, Frame-Independent, Color-Confining



# QED atoms: positronium and muonium

Coupled Fock states

$$\left[-\frac{\Delta^2}{2m_{red}} + V_{\text{eff}}(\vec{S}, \vec{r})\right] \psi(\vec{r}) = E \ \psi(\vec{r})$$

Effective two-particle equation
Includes Lamb Shift, quantum corrections

$$\left[ -\frac{1}{2m_{\rm red}} \frac{d^2}{dr^2} + \frac{1}{2m_{\rm red}} \frac{\ell(\ell+1)}{r^2} + V_{\rm eff}(r, S, \ell) \right] \psi(r) = E \psi(r)$$

$$V_{eff} \to V_C(r) = -\frac{\alpha}{r}$$

Semiclassical first approximation to QED



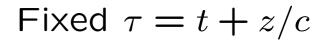
Spherical Basis  $r, heta, \phi$ 

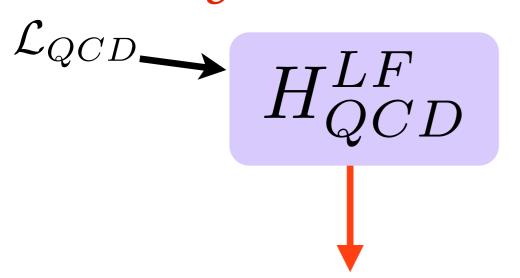
Coulomb potential

Bohr Spectrum

Schrödinger Eq.

## Light-Front QCD





$$(H_{LF}^0 + H_{LF}^I)|\Psi> = M^2|\Psi>$$

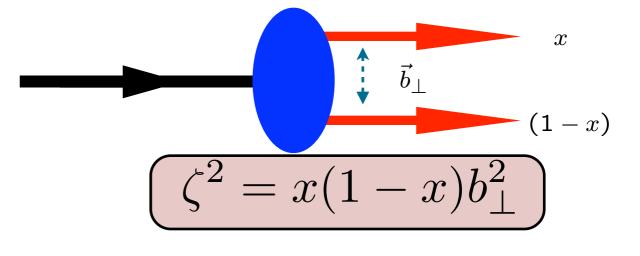
$$\left[\frac{\vec{k}_{\perp}^{2} + m^{2}}{x(1-x)} + V_{\text{eff}}^{LF}\right] \psi_{LF}(x, \vec{k}_{\perp}) = M^{2} \psi_{LF}(x, \vec{k}_{\perp})$$

$$\left[ -\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$

### AdS/QCD:

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

Semiclassical first approximation to QCD



#### Coupled Fock states

Eliminate higher Fock states and retarded interactions

Effective two-particle equation

Azimuthal Basis

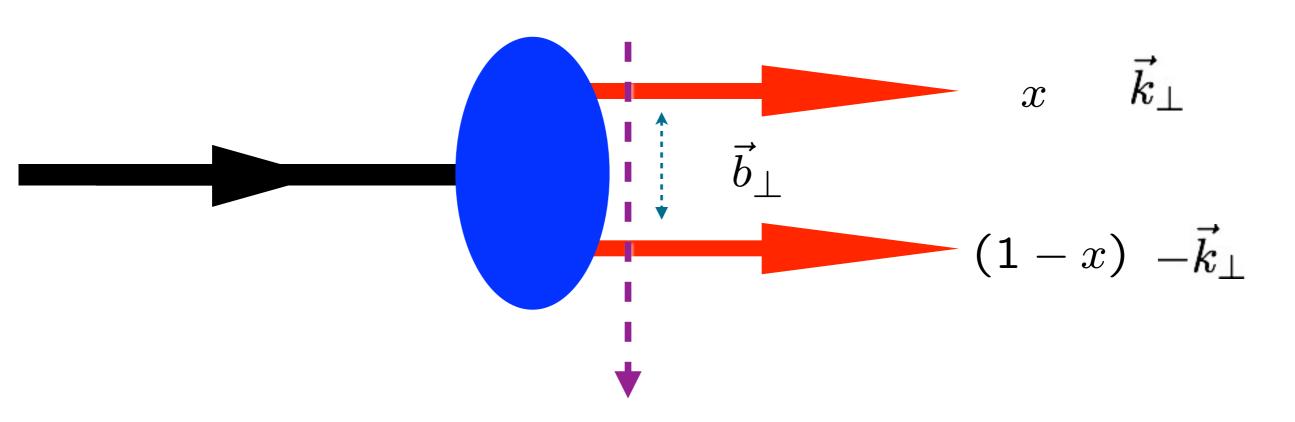
$$\zeta, \phi$$

$$m_q = 0$$

Confining AdS/QCD potential!

Sums an infinite # diagrams

## Fixed $\tau = t + z/c$



$$\zeta^2 \equiv b_\perp^2 x (1-x)$$

Invariant transverse separation

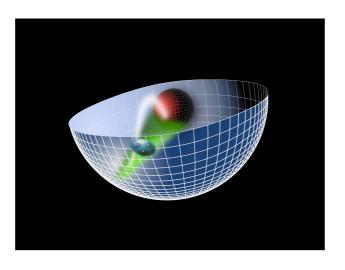
$$\zeta^2$$
 conjugate to  $\frac{k_\perp^2}{x(1-x)}=(p_q+p_{\bar q})^2=\mathcal{M}_{q+\bar q}^2$ 

$$\int dk^- \Psi_{BS}(P,k) \to \psi_{LF}(x,\vec{k}_\perp)$$

de Tèramond, Dosch, sjb

AdS/QCD Soft-Wall Model

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$



$$\zeta^2 = x(1-x)\mathbf{b}_{\perp}^2.$$

Light-Front Holography

$$\left[ -\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



#### Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

 $\kappa \simeq 0.6 \; GeV$ 

### Confinement scale:

$$1/\kappa \simeq 1/3 \ fm$$

- de Alfaro, Fubini, Furlan:
  - Fubini, Rabinovici:

Unique Confinement Potential!

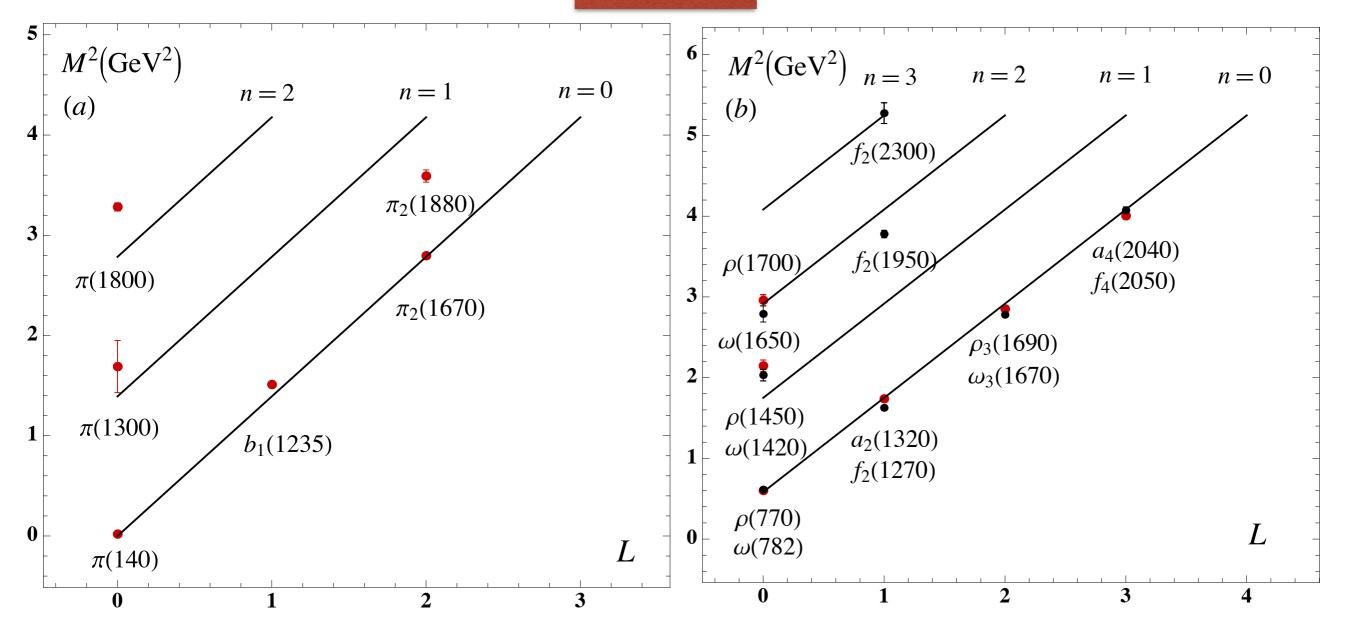
Scale can appear in Hamiltonian and EQM

without affecting conformal invariance of action!

Preserves Conformal Symmetry of the action

$$m_u = m_d = 0$$

## Preview



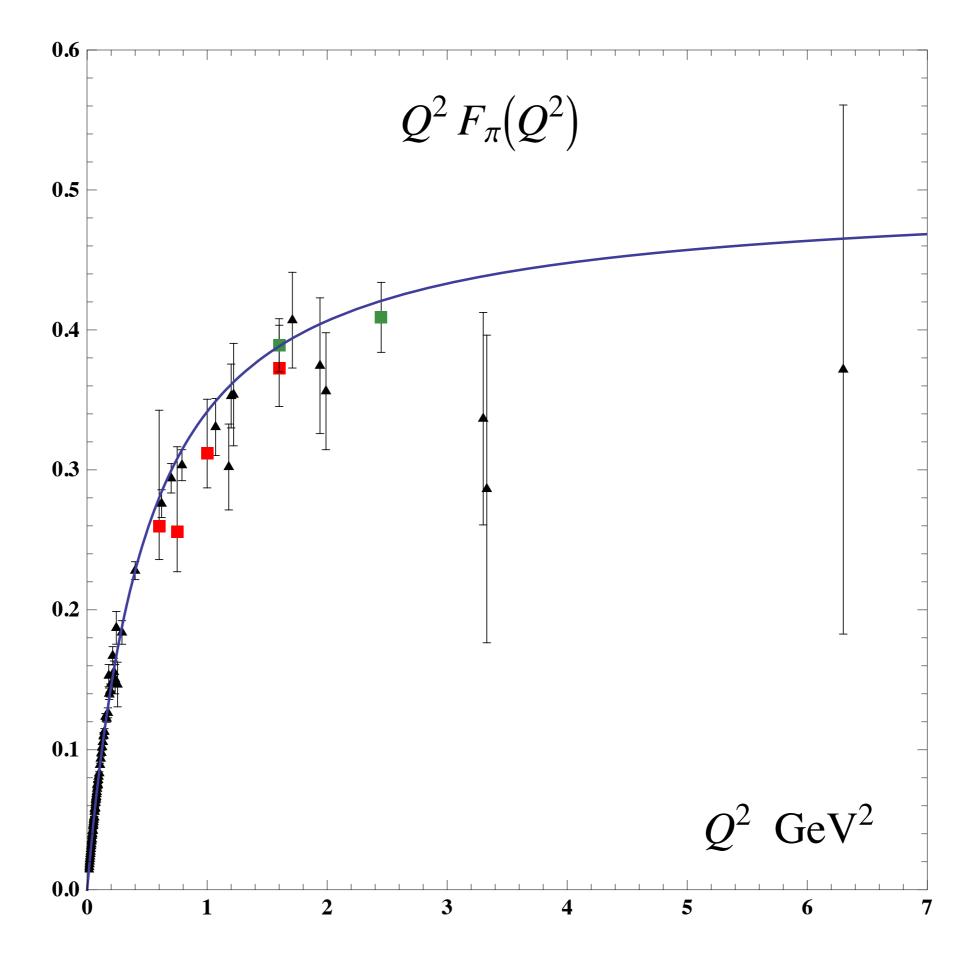
$$M^{2}(n, L, S) = 4\kappa^{2}(n + L + S/2)$$

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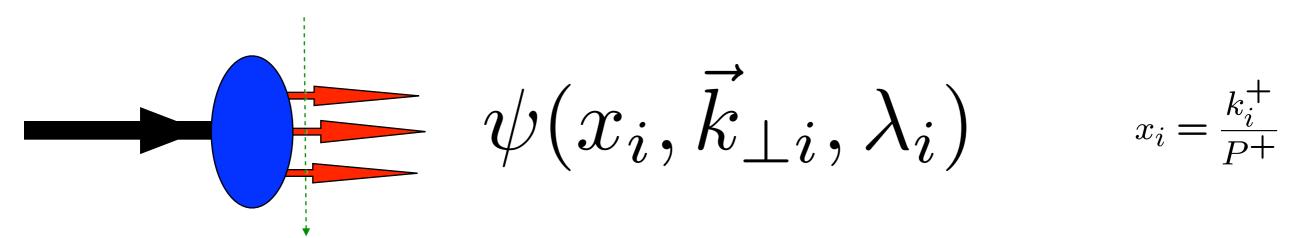
Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD

**Stan Brodsky** 





## Dirac's Front Form: Fixed $\tau = t + z/c$



Invariant under boosts. Independent of P^µ

$$H_{LF}^{QCD}|\psi>=M^2|\psi>$$

Direct connection to QCD Lagrangian

Remarkable new insights from AdS/CFT, the duality between conformal field theory and Anti-de Sitter Space

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Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD

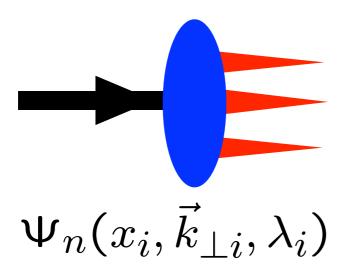
**Stan Brodsky** 

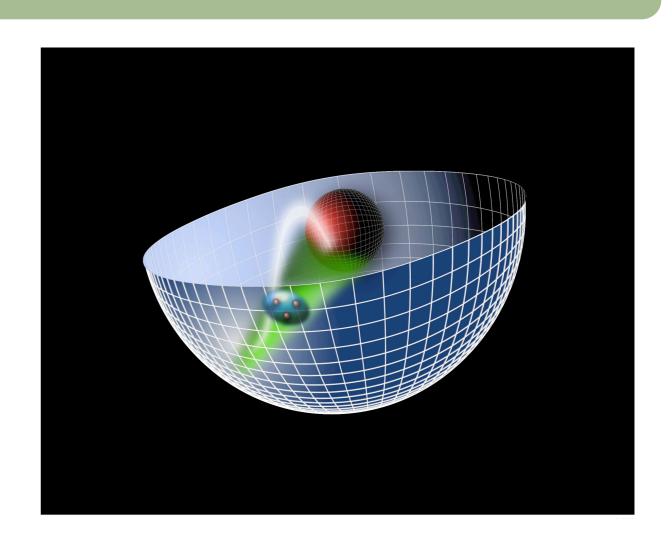


## Light-Front Holography and Non-Perturbative QCD

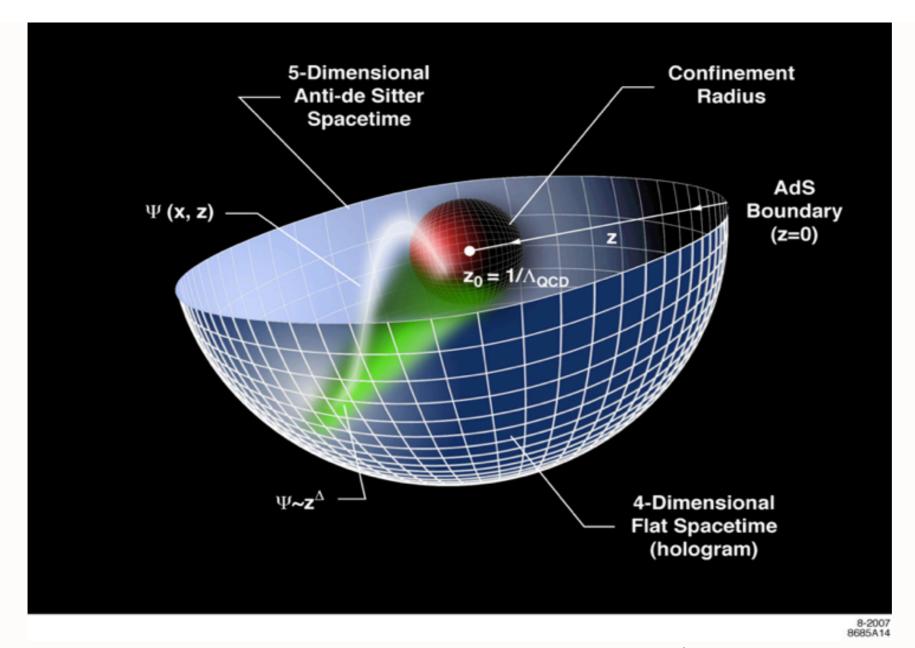
Goal:
Use AdS/QCD duality to construct
a first approximation to QCD

Hadron Spectrum Light-Front Wavefunctions, Form Factors, DVCS, etc





in collaboration with Guy de Teramond and H. Guenter Dosch



Changes in physical length scale mapped to evolution in the 5th dimension z

• Truncated AdS/CFT (Hard-Wall) model: cut-off at  $z_0=1/\Lambda_{\rm QCD}$  breaks conformal invariance and allows the introduction of the QCD scale (Hard-Wall Model) Polchinski and Strassler (2001).

• Smooth cutoff: introduction of a background dilaton field  $\varphi(z)$  – usual linear Regge dependence can be obtained (Soft-Wall Model) Karch, Katz, Son and Stephanov (2006).

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## AdS/CFT

ullet Isomorphism of SO(4,2) of conformal QCD with the group of isometries of AdS space

$$ds^2 = \frac{R^2}{z^2} (\eta_{\mu\nu} dx^\mu dx^\nu - dz^2), \end{area}$$
 invariant measure

 $x^{\mu} \to \lambda x^{\mu}, \ z \to \lambda z$ , maps scale transformations into the holographic coordinate z.

- AdS mode in z is the extension of the hadron wf into the fifth dimension.
- ullet Different values of z correspond to different scales at which the hadron is examined.

$$x^2 \to \lambda^2 x^2, \quad z \to \lambda z.$$

 $x^2 = x_\mu x^\mu$ : invariant separation between quarks

ullet The AdS boundary at z o 0 correspond to the  $Q o \infty$ , UV zero separation limit.

# Dílaton-Modified AdS/QCD

$$ds^{2} = e^{\varphi(z)} \frac{R^{2}}{z^{2}} (\eta_{\mu\nu} x^{\mu} x^{\nu} - dz^{2})$$

- $\bullet$  Soft-wall dilaton profile breaks conformal invariance  $\,e^{\varphi(z)}=e^{+\kappa^2z^2}$
- Color Confinement
- ullet Introduces confinement scale  $\kappa$
- Uses AdS₅ as template for conformal theory

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## Ads Soft-Wall Schrodinger Equation for bound state of two scalar constituents:

$$\left[ -\frac{d^2}{dz^2} - \frac{1 - 4L^2}{4z^2} + U(z) \right] \Phi(z) = \mathcal{M}^2 \Phi(z)$$

$$U(z) = \kappa^4 z^2 + 2\kappa^2 (L + S - 1)$$

Derived from variation of Action for Dilaton-Modified AdS5

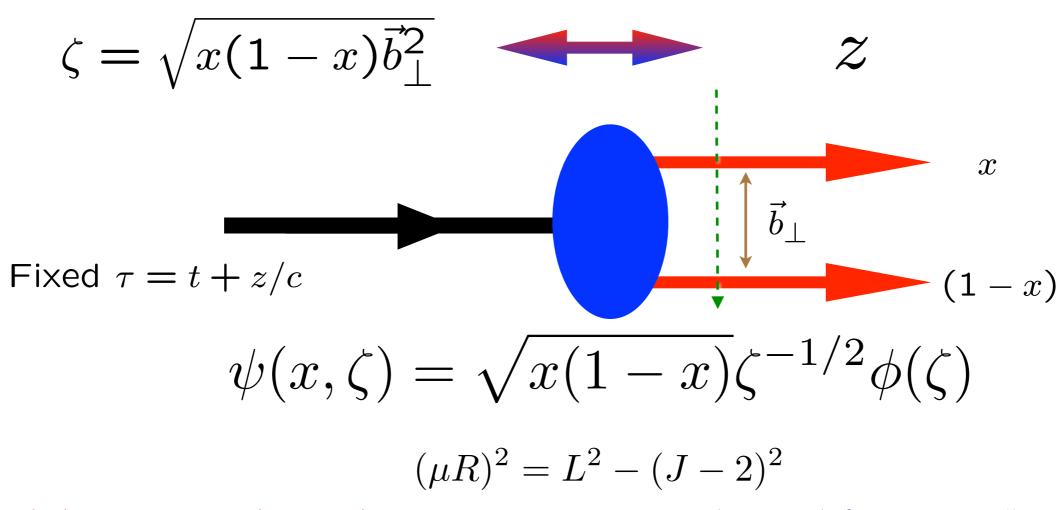
#### Identical to Light-Front Bound State Equation!

$$z \qquad \qquad \zeta = \sqrt{x(1-x)\vec{b}_{\perp}^2}$$



## Light-Front Holographic Dictionary

$$\psi(x,\vec{b}_{\perp})$$
  $\phi(z)$ 

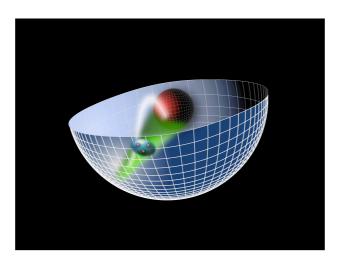


**Light-Front Holography**: Unique mapping derived from equality of LF and AdS formula for EM and gravitational current matrix elements and identical equations of motion

de Tèramond, Dosch, sjb

AdS/QCD Soft-Wall Model

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$



$$\zeta^2 = x(1-x)\mathbf{b}_{\perp}^2.$$

Light-Front Holography

$$\left[ -\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



#### Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

 $\kappa \simeq 0.6 \; GeV$ 

### Confinement scale:

$$1/\kappa \simeq 1/3 \ fm$$

- de Alfaro, Fubini, Furlan:
  - Fubini, Rabinovici:

Unique Confinement Potential!

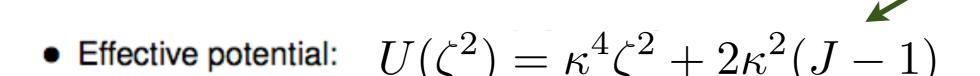
Scale can appear in Hamiltonian and EQM

without affecting conformal invariance of action!

Preserves Conformal Symmetry of the action

#### Meson Spectrum in Soft Wall Model

Píon: Negative term for J=0 cancels positive terms from LFKE and potential



LF WE

$$\left(-\frac{d^2}{d\zeta^2} - \frac{1 - 4L^2}{4\zeta^2} + \kappa^4 \zeta^2 + 2\kappa^2 (J - 1)\right) \phi_J(\zeta) = M^2 \phi_J(\zeta)$$

• Normalized eigenfunctions  $\langle \phi | \phi \rangle = \int d\zeta \, \phi^2(z)^2 = 1$ 

$$\phi_{n,L}(\zeta) = \kappa^{1+L} \, \sqrt{\frac{2n!}{(n+L)!}} \, \zeta^{1/2+L} e^{-\kappa^2 \zeta^2/2} L_n^L(\kappa^2 \zeta^2)$$

Eigenvalues

$$\mathcal{M}_{n,J,L}^2 = 4\kappa^2 \left(n + rac{J+L}{2}
ight)$$

G. de Teramond, H. G. Dosch, sjb

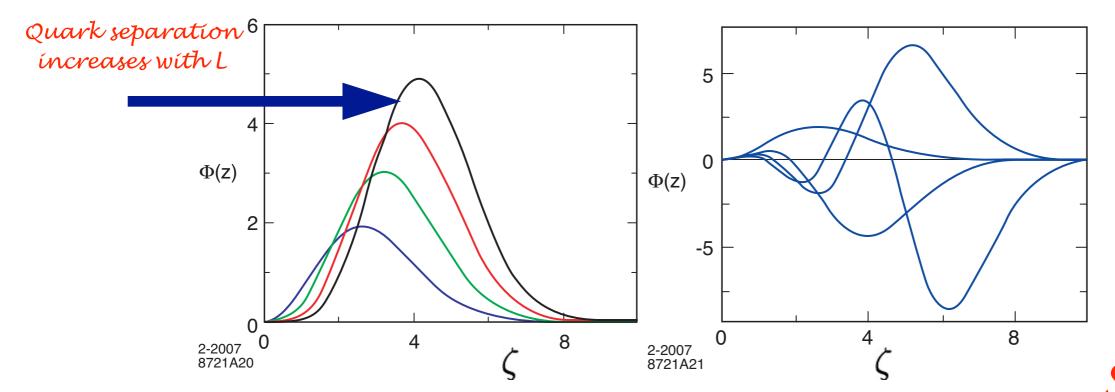
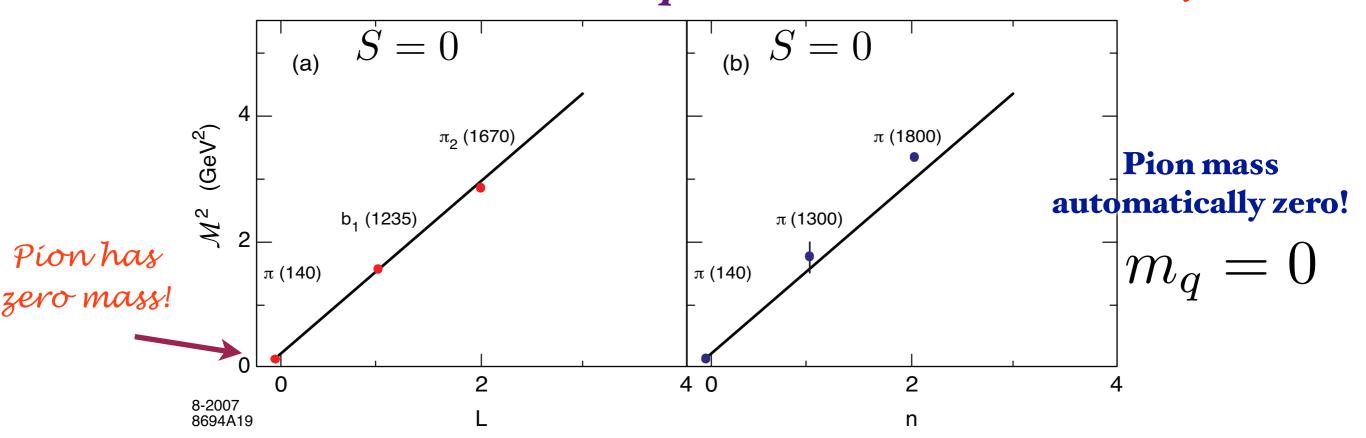


Fig: Orbital and radial AdS modes in the soft wall model for  $\kappa$  = 0.6 GeV . Same slope in n and L!

Soft Wall Model



Light meson orbital (a) and radial (b) spectrum for  $\kappa=0.6$  GeV.

$$ullet$$
  $J=L+S$  ,  $I=1$  meson families  ${\cal M}_{n,L,S}^2=4\kappa^2\,(n+L+S/2)$ 

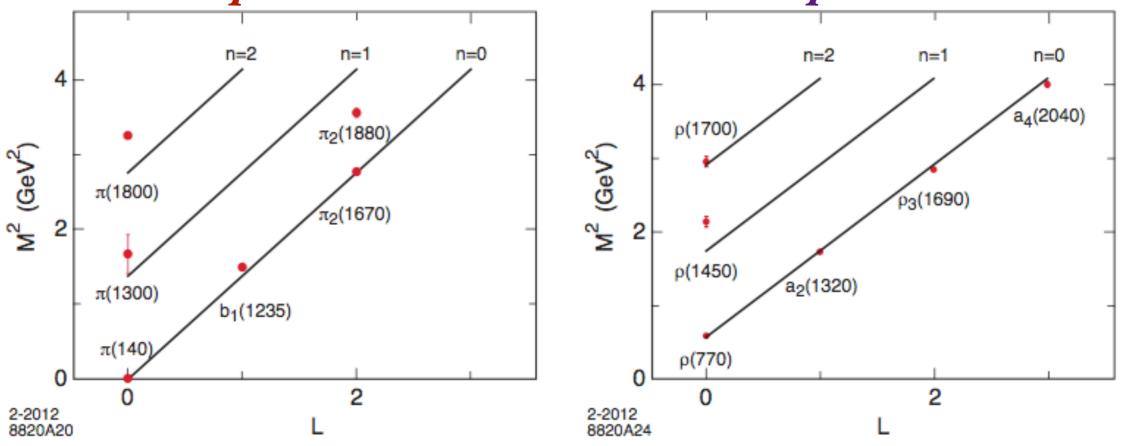
 $m_q = 0$ 

$$\mathcal{M}_{n,L,S}^2 = 4\kappa^2 \left(n + L + S/2\right)$$

$$4\kappa^2$$
 for  $\Delta n=1$   $4\kappa^2$  for  $\Delta L=1$   $2\kappa^2$  for  $\Delta S=1$ 

### Massless pion in Chiral Limit!

#### Same slope in n and L!



I=1 orbital and radial excitations for the  $\pi$  ( $\kappa=0.59$  GeV) and the ho-meson families ( $\kappa=0.54$  GeV)

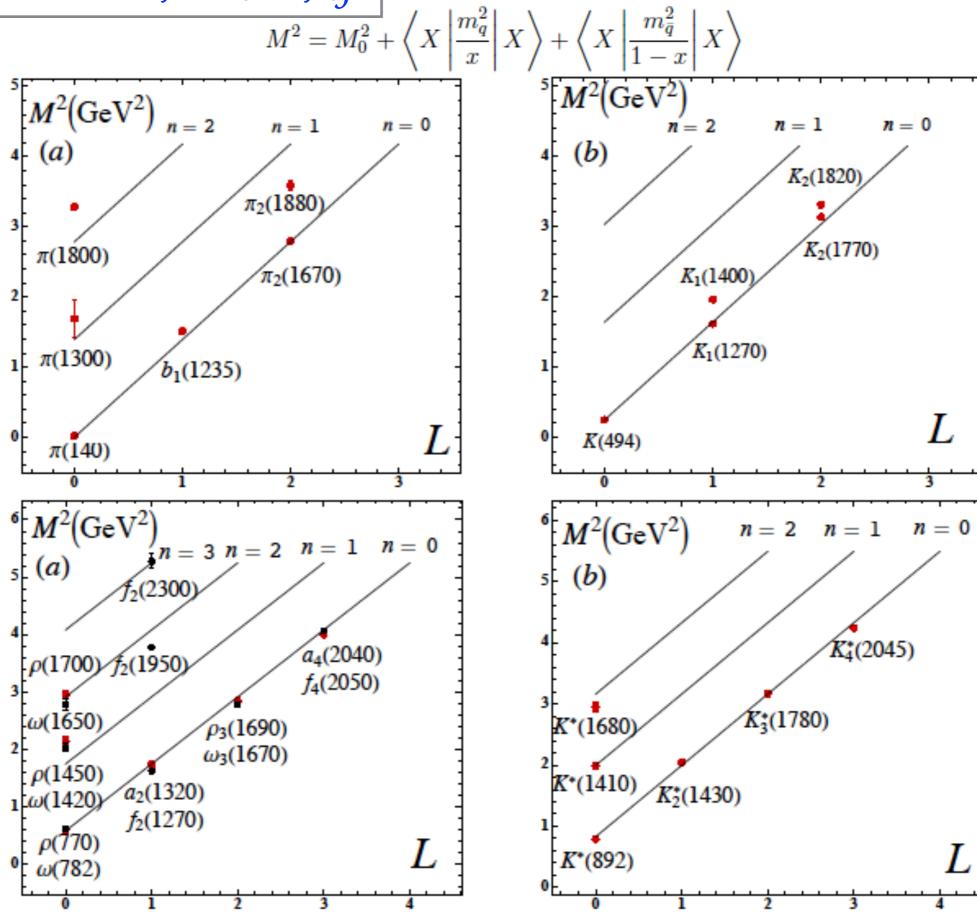
Triplet splitting for the I=1, L=1, J=0,1,2, vector meson a-states

$$\mathcal{M}_{a_2(1320)} > \mathcal{M}_{a_1(1260)} > \mathcal{M}_{a_0(980)}$$

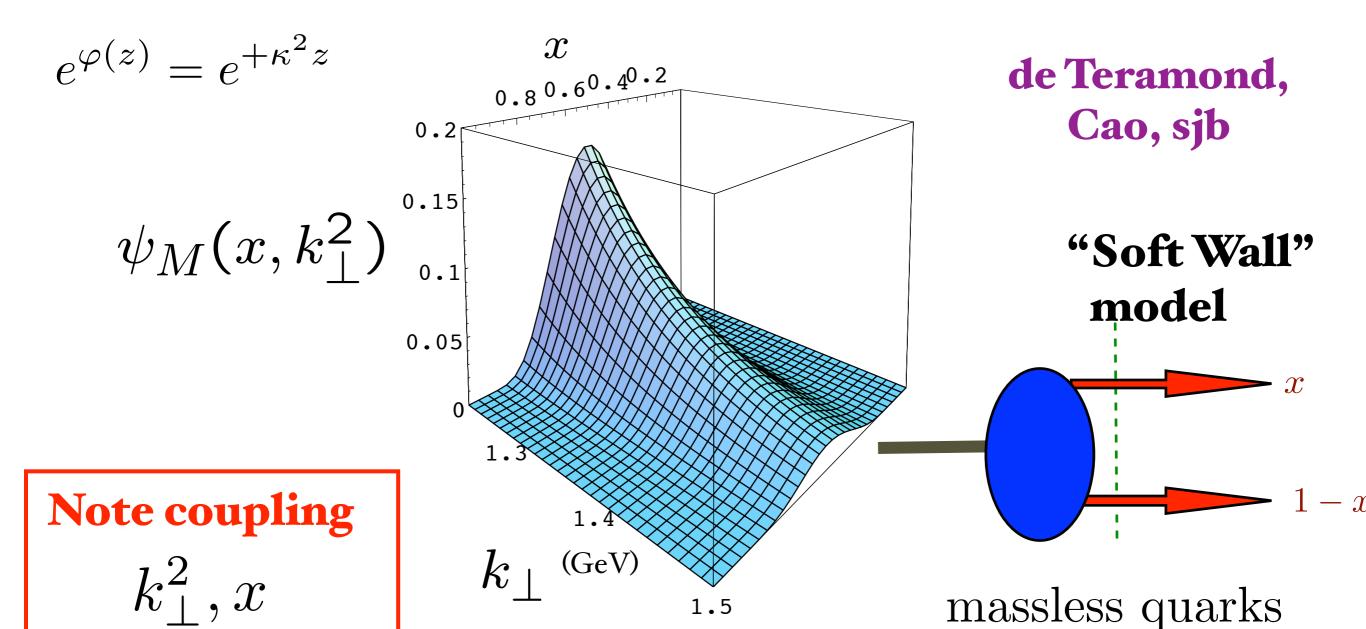
Mass ratio of the ρ and the a₁ mesons: coincides with Weinberg sum rules

G. de Teramond, H. G. Dosch, sjb

De Tèramond, Dosch, sjb



## Prediction from AdS/QCD: Meson LFWF



$$\psi_M(x,k_{\perp}) = \frac{4\pi}{\kappa\sqrt{x(1-x)}} e^{-\frac{k_{\perp}^2}{2\kappa^2x(1-x)}} \quad \left[\phi_{\pi}(x) = \frac{4}{\sqrt{3}\pi} f_{\pi} \sqrt{x(1-x)}\right]$$

$$f_{\pi} = \sqrt{P_{q\bar{q}}} \frac{\sqrt{3}}{8} \kappa = 92.4 \text{ MeV}$$

Same as DSE!

Provides Connection of Confinement to Hadron Structure

## AdS/QCD Holographic Wave Function for the $\rho$ Meson and Diffractive $\rho$ Meson Electroproduction

#### J. R. Forshaw*

Consortium for Fundamental Physics, School of Physics and Astronomy, University of Manchester, Oxford Road, Manchester M13 9PL, United Kingdom

#### R. Sandapen[†]

Département de Physique et d'Astronomie, Université de Moncton, Moncton, New Brunswick E1A3E9, Canada (Received 5 April 2012; published 20 August 2012)

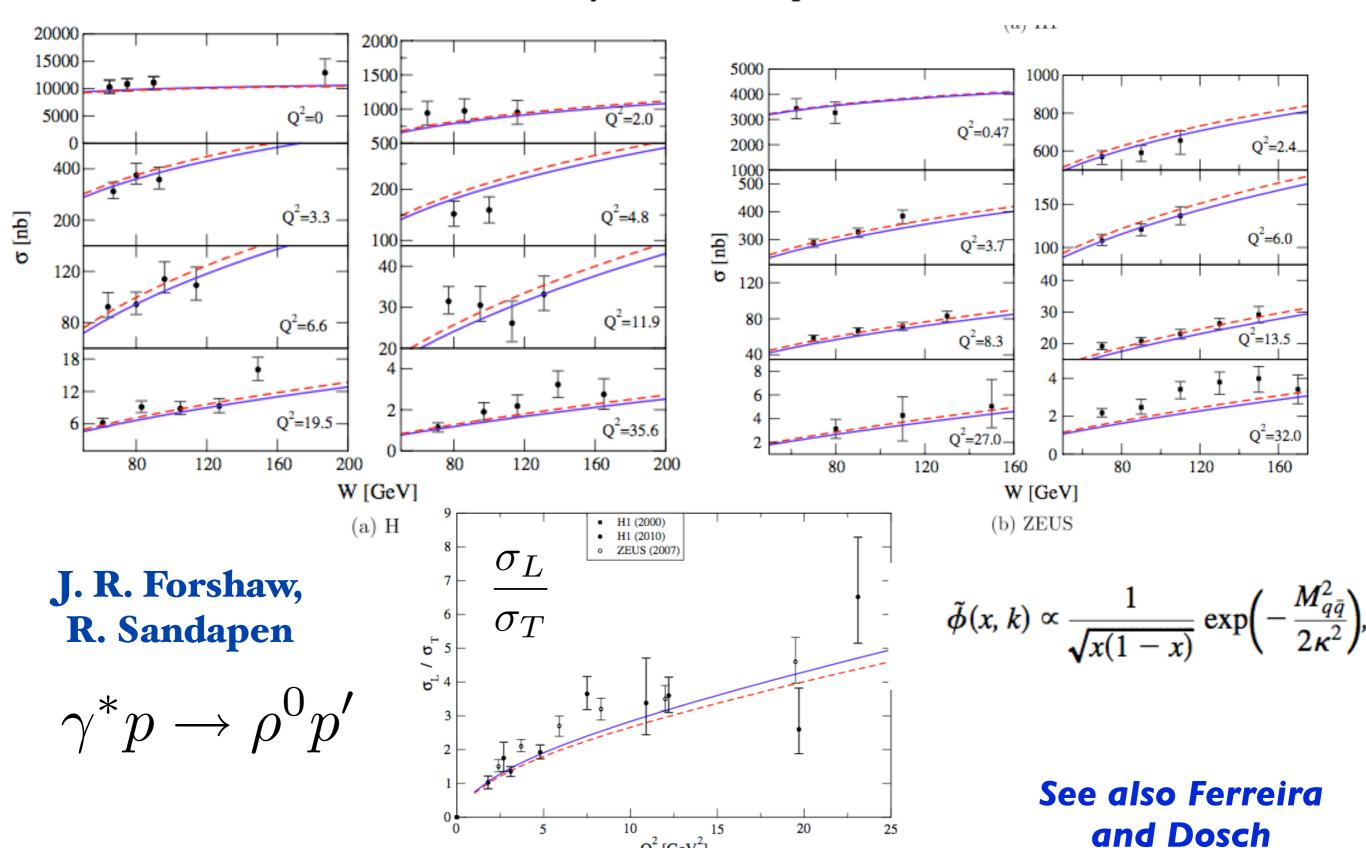
We show that anti-de Sitter/quantum chromodynamics generates predictions for the rate of diffractive  $\rho$ -meson electroproduction that are in agreement with data collected at the Hadron Electron Ring Accelerator electron-proton collider.

$$\psi_M(x, k_{\perp}) = \frac{4\pi}{\kappa \sqrt{x(1-x)}} e^{-\frac{k_{\perp}^2}{2\kappa^2 x(1-x)}}$$

## See also Ferreira and Dosch

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$

#### AdS/QCD Holographic Wave Function for the $\rho$ Meson and Diffractive $\rho$ Meson Electroproduction



 $Q^2 [GeV^2]$ 

#### de Tèramond, Dosch, sjb

# Uniqueness

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$
  $e^{\varphi(z)} = e^{+\kappa^2 z^2}$ 

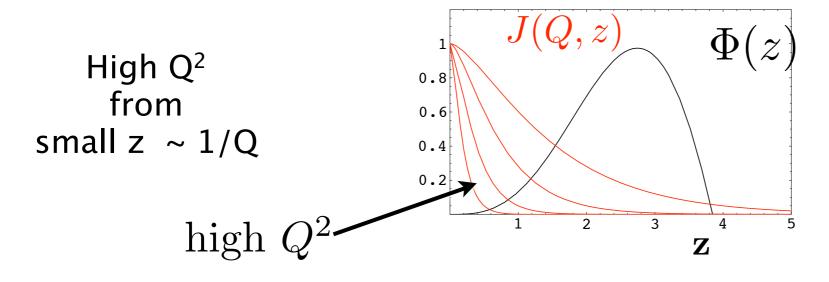
- $\zeta_2$  confinement potential and dilaton profile unique!
- Linear Regge trajectories in n and L: same slope!
- Massless pion in chiral limit! No vacuum condensate!
- Conformally invariant action for massless quarks retained despite mass scale
- Same principle, equation of motion as de Alfaro, Furlan, Fubini,
   Conformal Invariance in Quantum Mechanics Nuovo Cim. A34 (1976)
   569

#### **Hadron Form Factors from AdS/QCD**

Propagation of external perturbation suppressed inside AdS.

$$J(Q,z) = zQK_1(zQ)$$

$$F(Q^2)_{I\to F} = \int \frac{dz}{z^3} \Phi_F(z) J(Q, z) \Phi_I(z)$$



Polchinski, Strassler de Teramond, sjb

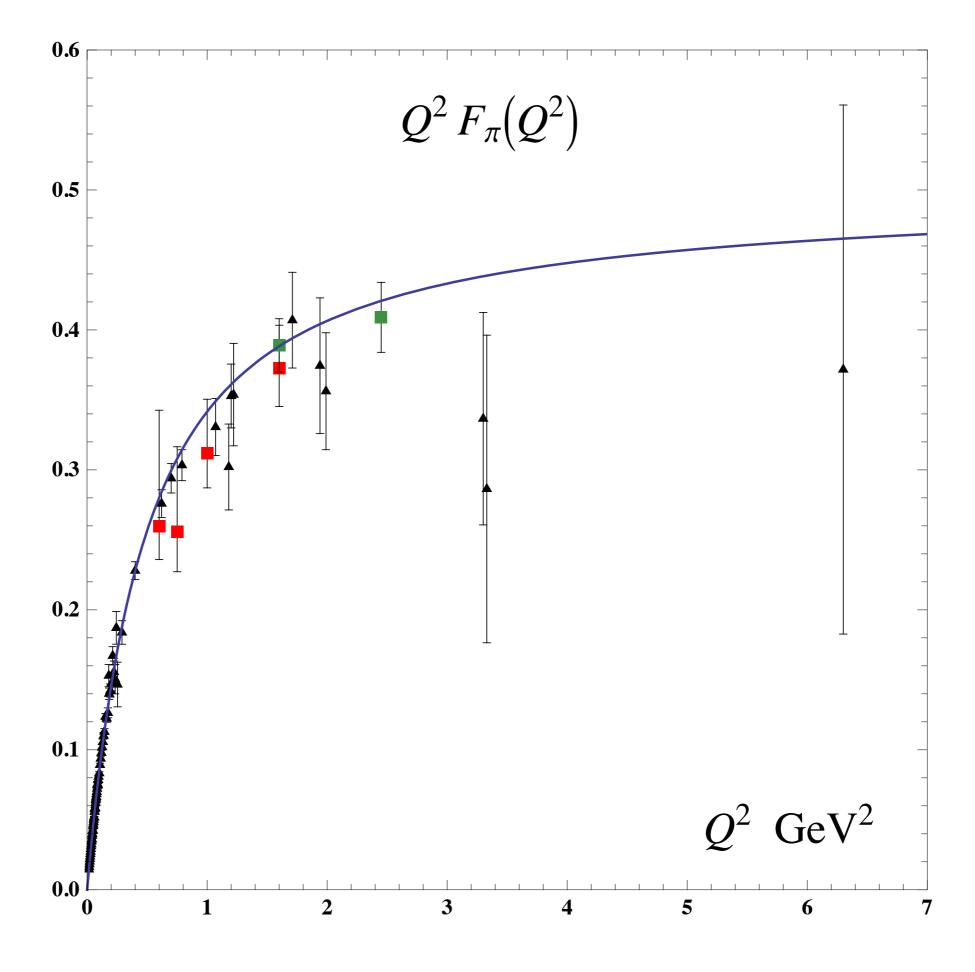
Consider a specific AdS mode  $\Phi^{(n)}$  dual to an n partonic Fock state  $|n\rangle$ . At small z,  $\Phi^{(n)}$ scales as  $\Phi^{(n)} \sim z^{\Delta_n}$ . Thus:

$$F(Q^2) \to \left[\frac{1}{Q^2}\right]^{7-1}$$

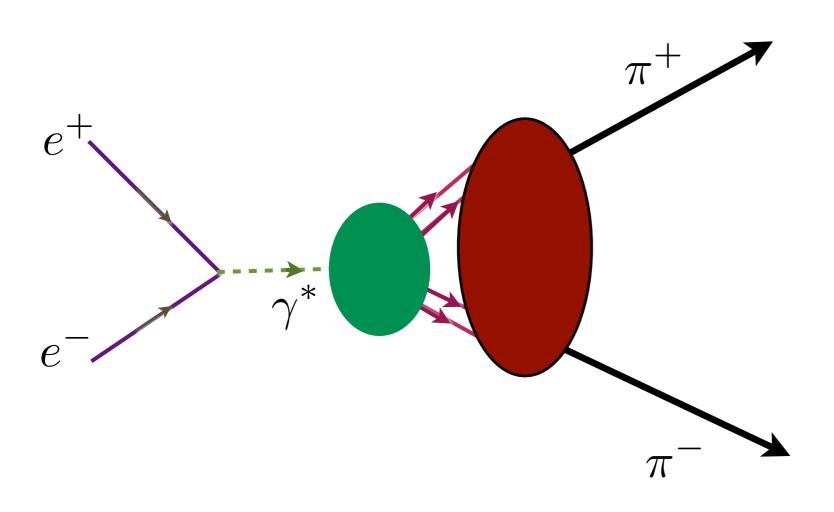
where 
$$\tau = \Delta_n - \sigma_n$$
,  $\sigma_n = \sum_{i=1}^n \sigma_i$ .

 $F(Q^2) \to \left\lceil \frac{1}{O^2} \right\rceil^{\tau-1}, \qquad \begin{array}{c} \text{Dimensional Quark Counting Rules:} \\ \text{General result from} \end{array}$ **AdS/CFT and Conformal Invariance** 

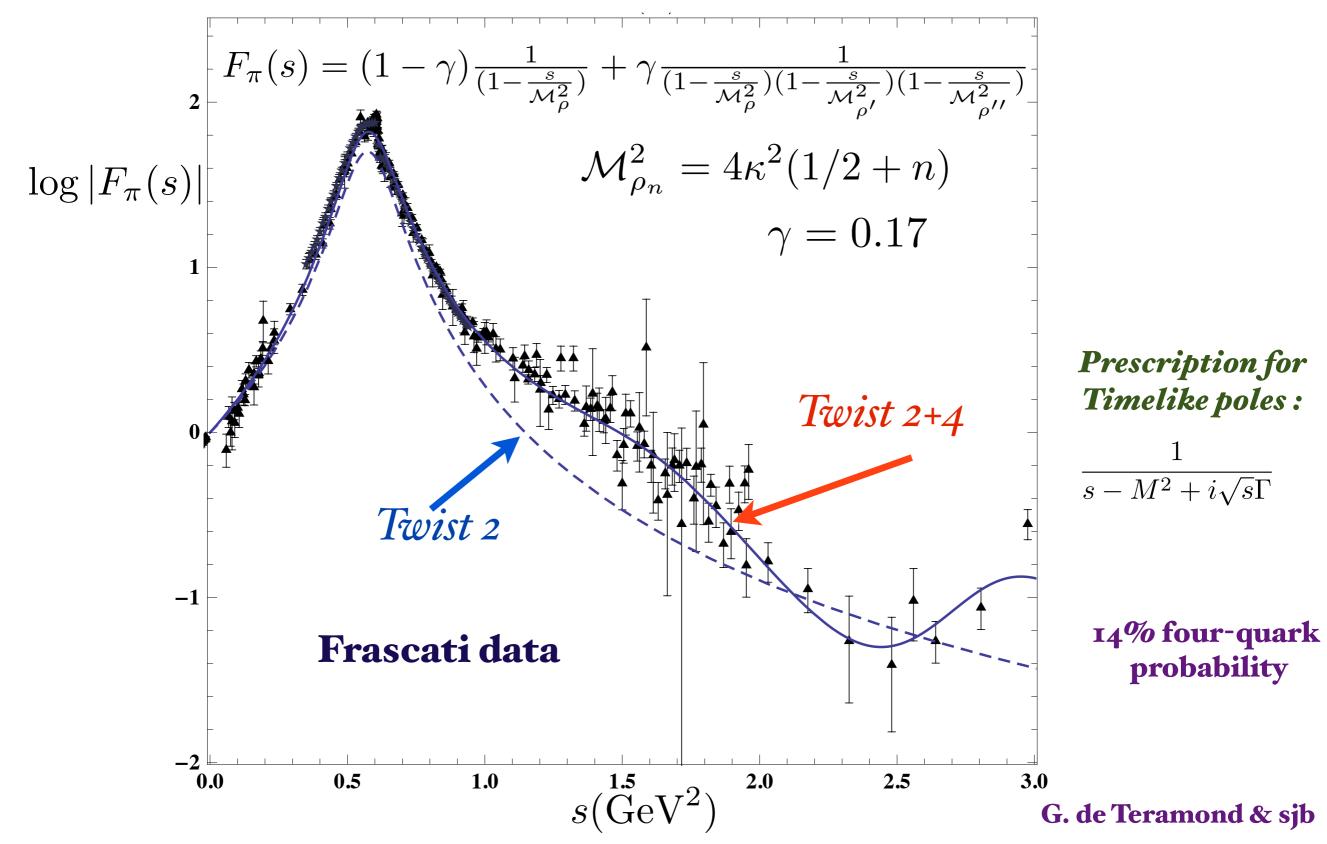
Twist 
$$\tau = n + L$$



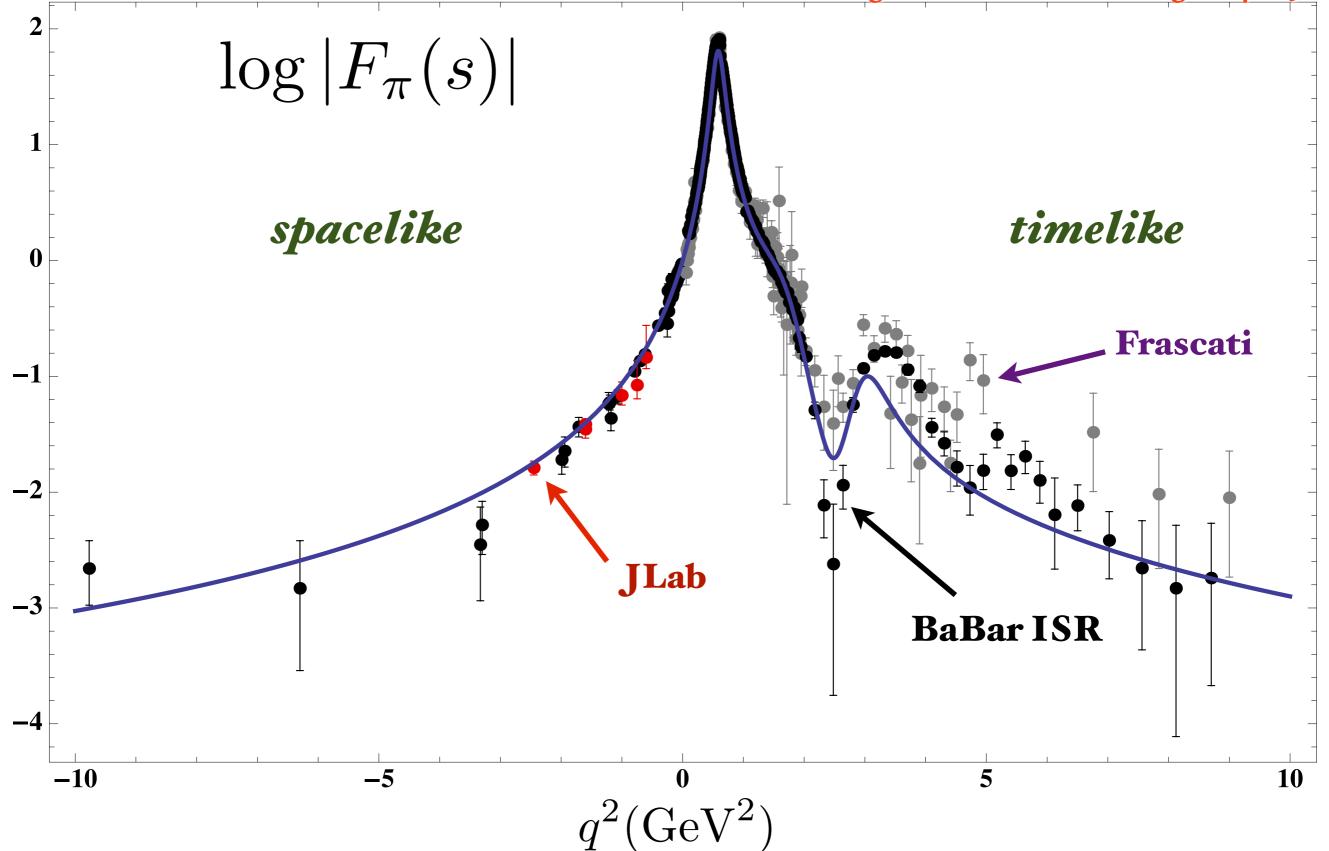
# Dressed soft-wall current brings in higher Fock states and more vector meson poles



# Timelike Pion Form Factor from AdS/QCD and Light-Front Holography



Pion Form Factor from AdS/QCD and Light-Front Holography



## Remarkable Features of Light-Front Schrödinger Equation

- Relativistic, frame-independent
- QCD scale appears unique LF potential
- Reproduces spectroscopy and dynamics of light-quark hadrons with one parameter
- Zero-mass pion for zero mass quarks!
- Regge slope same for n and L -- not usual HO
- Splitting in L persists to high mass -- contradicts conventional wisdom based on breakdown of chiral symmetry
- Phenomenology: LFWFs, Form factors, electroproduction
- Extension to heavy quarks

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

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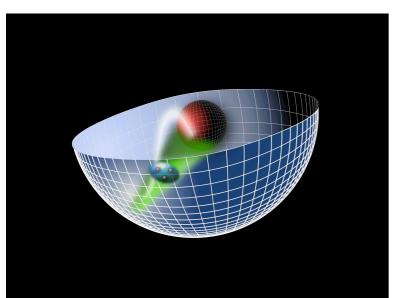
Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD



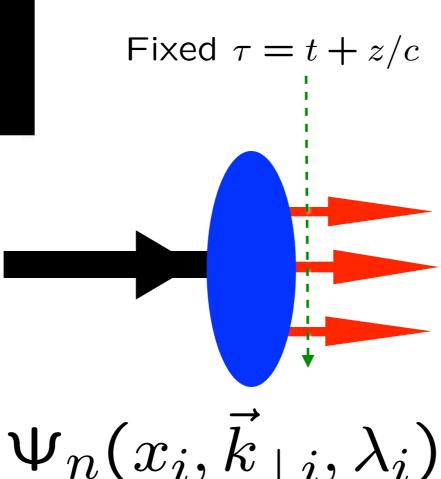


 $\phi(z)$ 

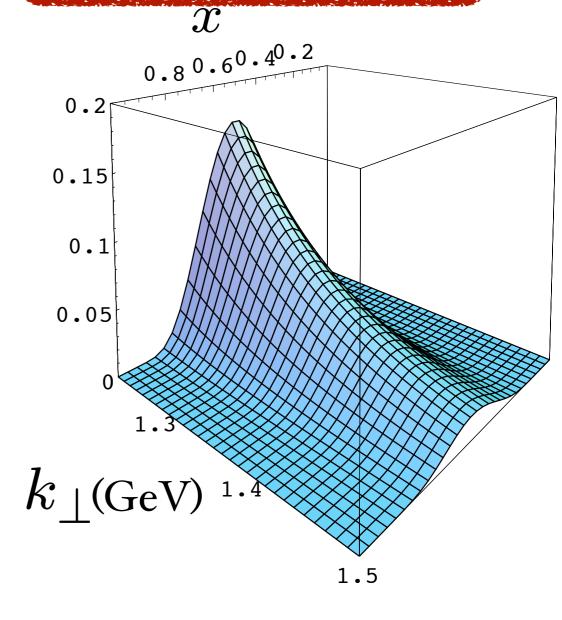
### AdS5: Conformal Template for QCD



· Light-Front Holography



Duality of AdS₅ with LF Hamiltonian Theory



• Light Front Wavefunctions:

Light-Front Schrödinger Equation
Spectroscopy and Dynamics

# QCD Lagrangian

$$\mathcal{L}_{QCD} = -\frac{1}{4} Tr(G^{\mu\nu} G_{\mu\nu}) + \sum_{f=1}^{n_f} i \bar{\Psi}_f D_{\mu} \gamma^{\mu} \Psi_f + \sum_{f=1}^{n_f} I_f \bar{\Psi}_f \Psi_f$$

$$iD^{\mu} = i\partial^{\mu} - gA^{\mu} \qquad G^{\mu\nu} = \partial^{\mu}A^{\mu} - \partial^{\nu}A^{\mu} - g[A^{\mu}, A^{\nu}]$$

### Classical Chiral Lagrangian is Conformally Invariant

Where does the QCD Mass Scale  $\Lambda_{\rm QCD}$  come from?

How does color confinement arise?

ode Alfaro, Fubini, Furlan:

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

Unique confinement potential!

#### de Alfaro, Fubini, Furlan

$$G|\psi(\tau)>=i\frac{\partial}{\partial\tau}|\psi(\tau)>$$

$$G=uH+vD+wK$$

$$G=H_{\tau}=\frac{1}{2}\left(-\frac{d^2}{dx^2}+\frac{g}{r^2}+\frac{4uw-v^2}{4}x^2\right)$$

Retains conformal invariance of action despite mass scale!

$$4uw - v^2 = \kappa^4 = [M]^4$$

Identical to LF Hamiltonian with unique potential and dilaton!

 $\begin{bmatrix} -\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \end{bmatrix} \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$   $U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L+S-1)$ 

## What determines the QCD mass scale $\Lambda_{QCD}$ ?

- Mass scale does not appear in the QCD Lagrangian (massless quarks)
- $\bullet$  Dimensional Transmutation? Requires external constraint such as  $~\alpha_s(M_Z)$
- dAFF: Confinement Scale K appears spontaneously via the Hamiltonian:  $G=uH+vD+wK \qquad 4uw-v^2=\kappa^4=[M]^4$
- The confinement scale regulates infrared divergences, connects  $\Lambda_{\mathcal{QCD}}$  to the confinement scale  $\kappa$
- Only dimensionless mass ratios (and M times R) predicted
- Mass and time units [GeV] and [sec] from physics external to QCD
- New feature: bounded frame-independent relative time between constituents

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August 31, 2015

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**Stan Brodsky** 



## dAFF: New Time Variable

$$\tau = \frac{2}{\sqrt{4uw - v^2}} \arctan\left(\frac{2tw + v}{\sqrt{4uw - v^2}}\right),\,$$

- Identify with difference of LF time  $\Delta x^+/P^+$  between constituents
- Finite range
- Measure in Double-Parton Processes

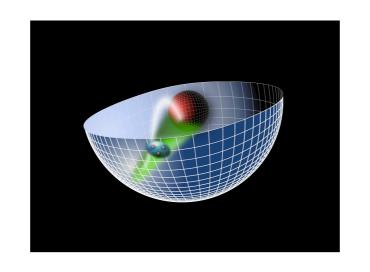






de Tèramond, Dosch, sjb

AdS/QCD Soft-Wall Model



$$\zeta^2 = x(1-x)\mathbf{b}_{\perp}^2.$$

Light-Front Holography

$$\left[ -\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



#### Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

 $\kappa \simeq 0.6 \; GeV$ 

### Confinement scale:

$$1/\kappa \simeq 1/3 \ fm$$

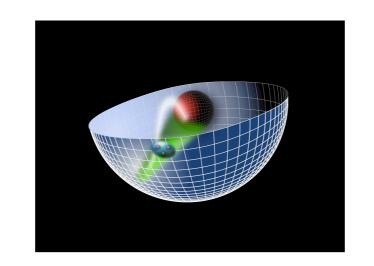
de Alfaro, Fubini, Furlan:

Unique Confinement Potential!

Conformal Symmetry of the action

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

AdS/QCD Soft-Wall Model



Light-Front Holography

### Semi-Classical Approximation to QCD

Relativistic, frame-independent
Unique color-confining potential
Zero mass pion for massless quarks
Regge trajectories with equal slopes in n and L
Light-Front Wavefunctions

Light-Front Schrödinger Equation

Conformal Symmetry of the action

### Superconformal Algebra

de Teramond Dosch and SJB

1+1

$$\{\psi,\psi^+\}=1$$

two anti-commuting fermionic operators

$$\psi=rac{1}{2}(\sigma_1-i\sigma_2), \quad \psi^+=rac{1}{2}(\sigma_1+i\sigma_2)$$
 Realization as Pauli Matrices

$$Q = \psi^{+}[-\partial_{x} + W(x)], \quad Q^{+} = \psi[\partial_{x} + W(x)], \qquad W(x) = \frac{f}{x}$$
(Conformal)

$$S = \psi^+ x, \quad S^+ = \psi x$$

Introduce new spinor operators

$${Q, Q^+} = 2H, {S, S^+} = 2K$$

$$Q \simeq \sqrt{H}, \quad S \simeq \sqrt{K}$$

$${Q,Q} = {Q^+,Q^+} = 0, [Q,H] = [Q^+,H] = 0$$

### Superconformal Algebra

# Baryon Equation

Consider 
$$R_w = Q + wS;$$

w: dimensions of mass squared

$$G = \{R_w, R_w^+\} = 2H + 2w^2K + 2wfI - 2wB \qquad 2B = \sigma_3$$

Retains Conformal Invariance of Action

Fubini and Rabinovici

### New Extended Hamiltonian G is diagonal:

$$G_{11} = \left(-\partial_x^2 + w^2 x^2 + 2wf - w + \frac{4(f + \frac{1}{2})^2 - 1}{4x^2}\right)$$

$$G_{22} = \left(-\partial_x^2 + w^2 x^2 + 2wf + w + \frac{4(f - \frac{1}{2})^2 - 1}{4x^2}\right)$$

Identify 
$$f - \frac{1}{2} = L_B$$
,  $w = \kappa^2$ 

Eigenvalue of G:  $M^2(n,L) = 4\kappa^2(n+L_B+1)$ 

# LF Holography

### Baryon Equation

 $x \to \zeta$ 

$$\left(-\frac{d^2}{d\zeta^2} + \lambda_B^2 \zeta^2 + 2\lambda_B (L_B + 1) + \frac{4L_B^2 - 1}{4\zeta^2}\right) \psi_J^+ = M^2 \psi_J^+,$$

$$\left(-\frac{d^2}{d\zeta^2} + \lambda_B^2 \zeta^2 + 2\lambda_B L_B + \frac{4(L_B + 1)^2 - 1}{4\zeta^2}\right) \psi_J^- = M^2 \psi_J^-.$$



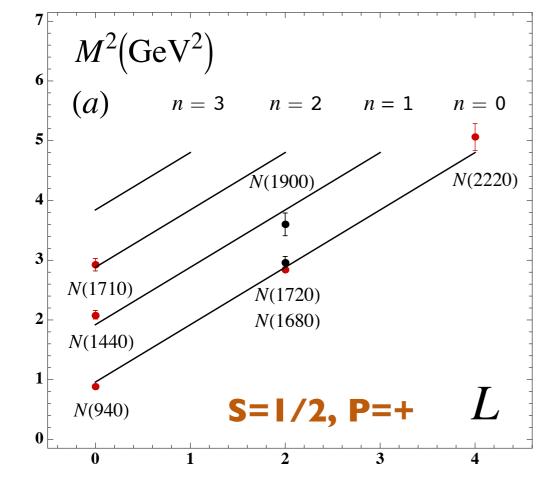
$$M_B^2(N,L_B) = 4\lambda_B(n+L_B+1)$$
 S=1/2, P=+

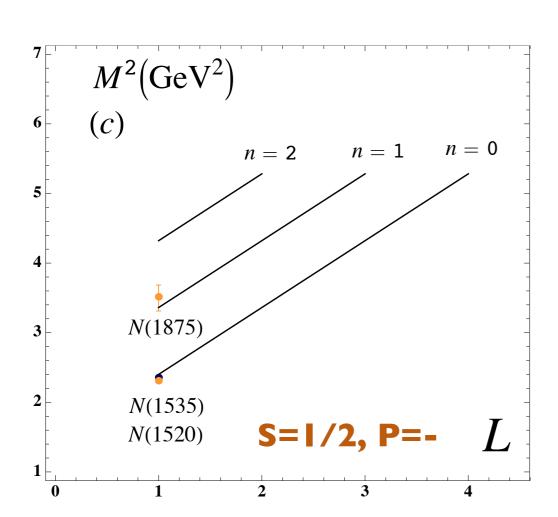
$$(-\frac{d^2}{d\zeta^2}+\lambda_M^2\zeta^2+2\lambda_M(J-1)+\frac{4L_M^2-1}{4\zeta^2})\phi_J=M^2\phi_J$$

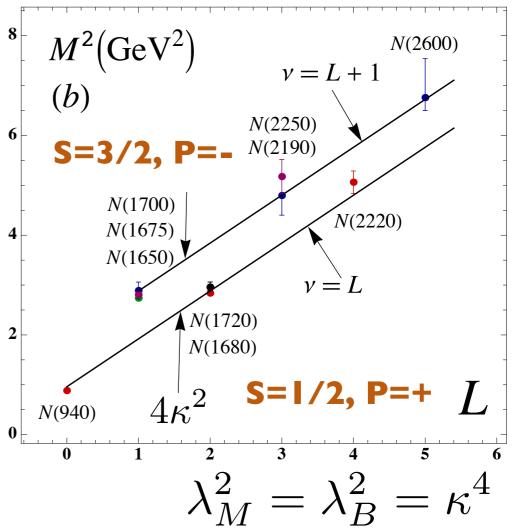
$$M_M^2(N, L_M, S = 0) = 4\lambda_M(n + L_M)$$

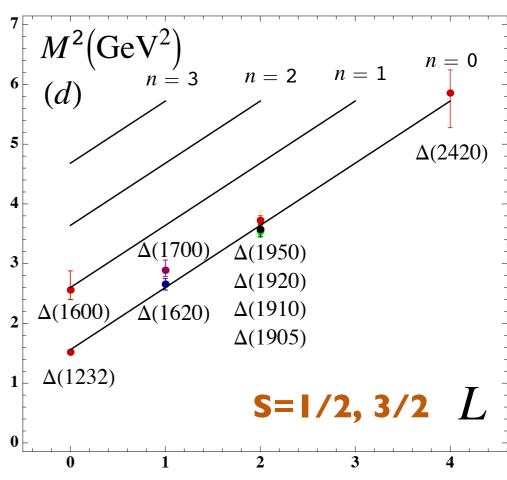
S=0, I=1 Meson is superpartner of S=1/2, I=1 Baryon

Meson-Baryon Degeneracy for L_M=L_B+1

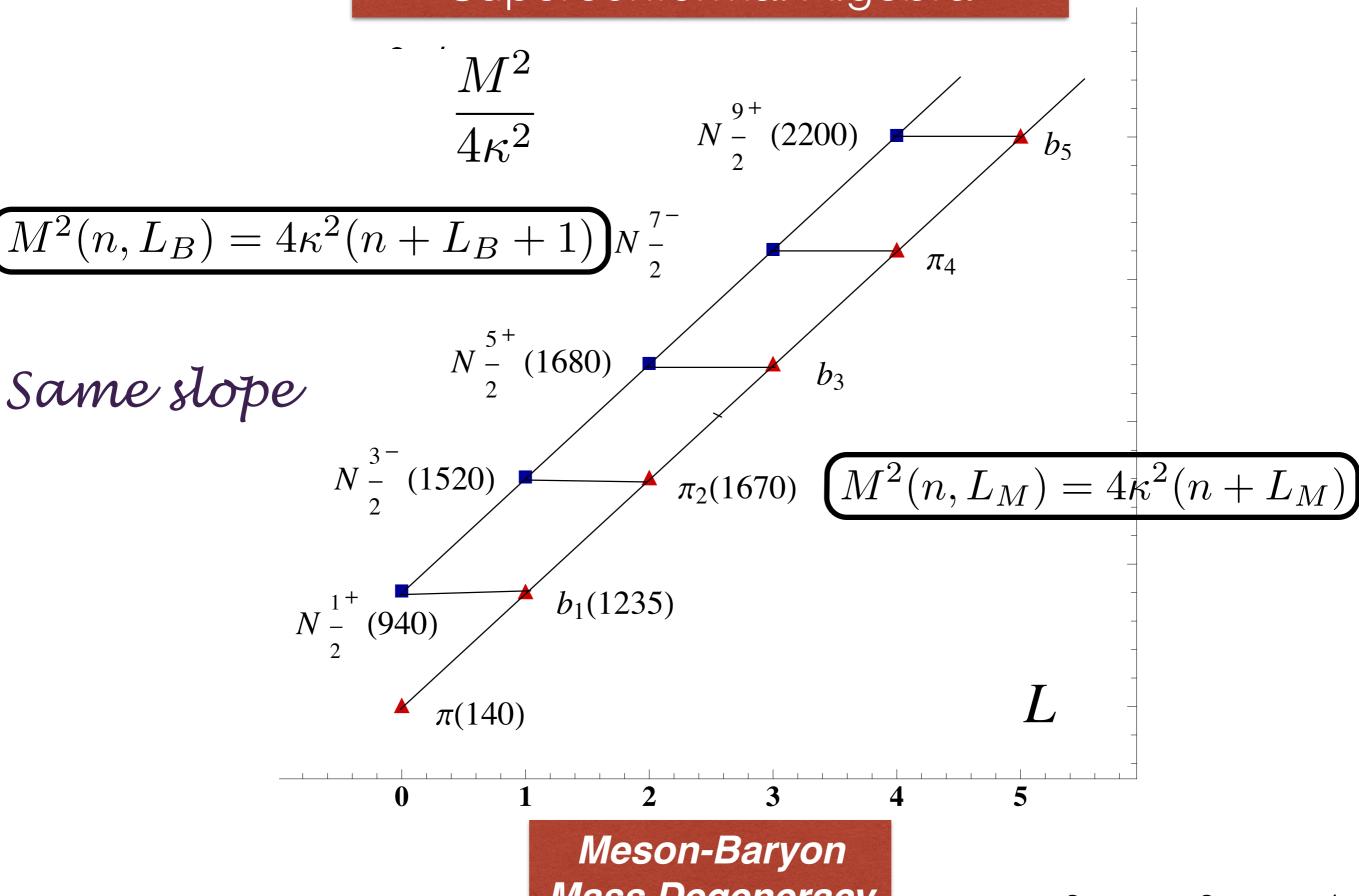








# Superconformal Algebra

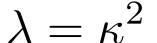


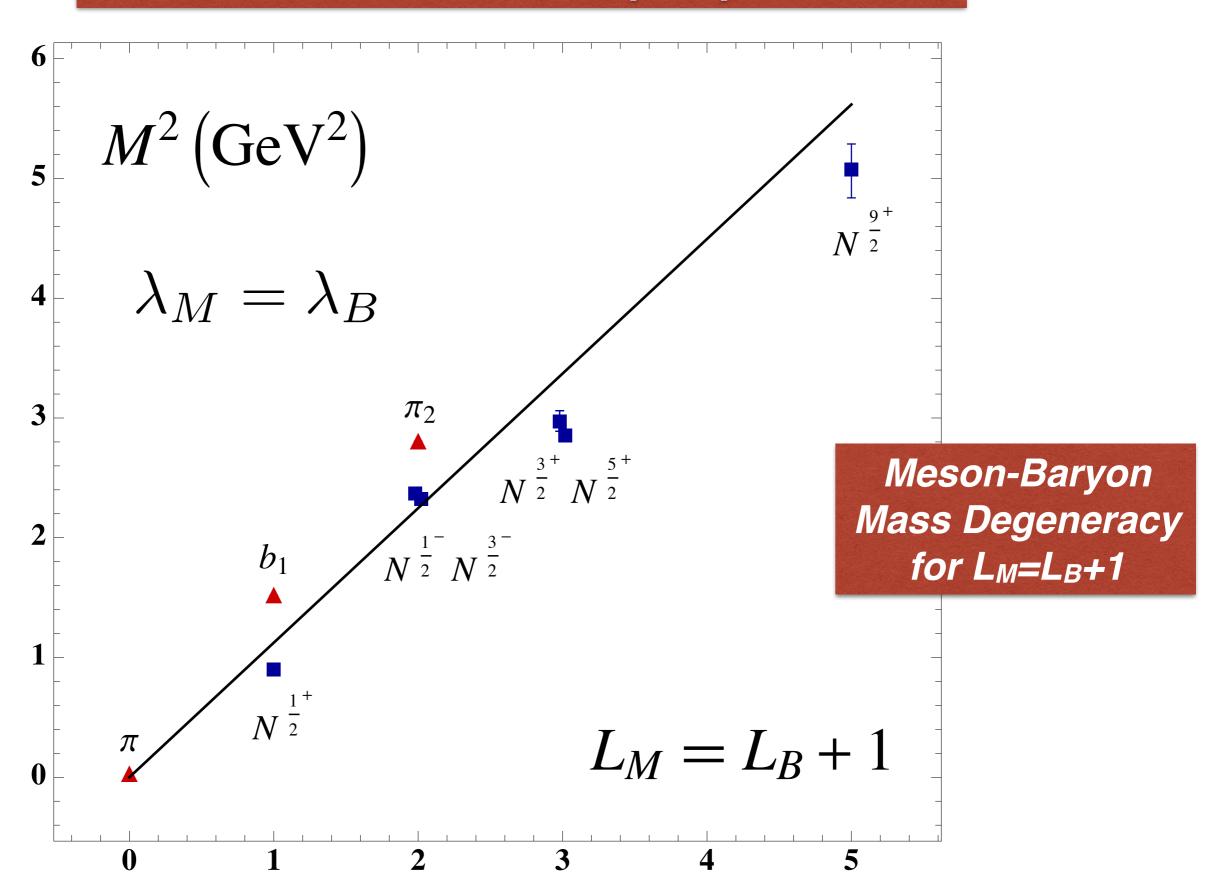
Meson-Baryon
Mass Degeneracy
for L_M=L_B+1

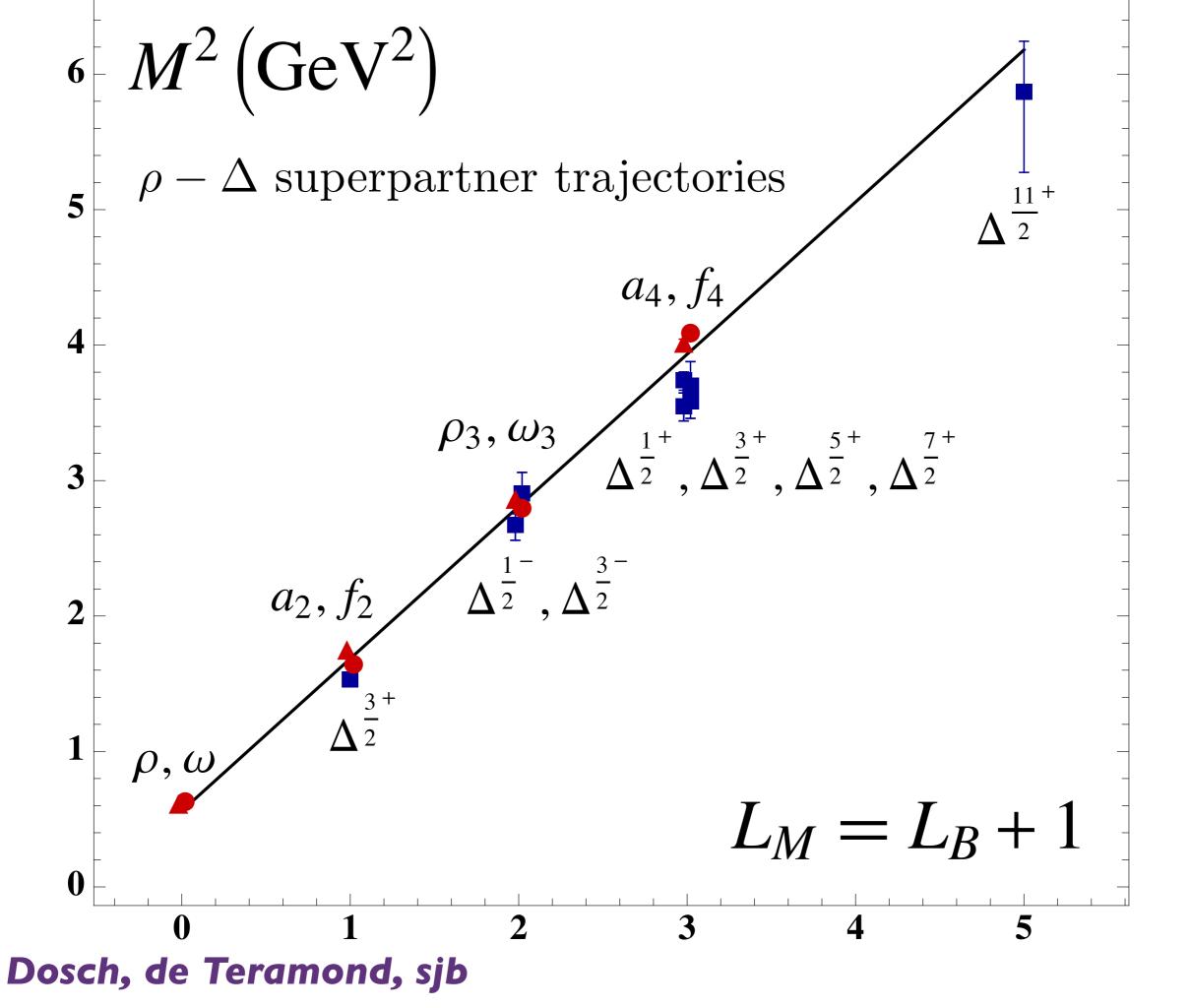
$$\lambda_M^2 = \lambda_B^2 = \kappa^4$$

# Superconformal AdS Light-Front Holographic QCD (LFHQCD):

Identical meson and baryon spectra!







# Chiral Features of Soft-Wall AdS/QCD Model

- Boost Invariant
- Trivial LF vacuum! No condensate, but consistent with GMOR
- Massless Pion
- Hadron Eigenstates (even the pion) have LF Fock components of different Lz
- Proton: equal probability  $S^z = +1/2, L^z = 0; S^z = -1/2, L^z = +1$   $J^z = +1/2 : < L^z > = 1/2, < S^z_q > = 0$
- Self-Dual Massive Eigenstates: Proton is its own chiral partner.
- Label State by minimum L as in Atomic Physics
- Minimum L dominates at short distances
- AdS/QCD Dictionary: Match to Interpolating Operator Twist at z=o.

  No mass -degenerate parity partners!

# Features of Supersymmetric Equations

 J =L+S baryon simultaneously satisfies both equations of G with L, L+1 for same mass eigenvalue

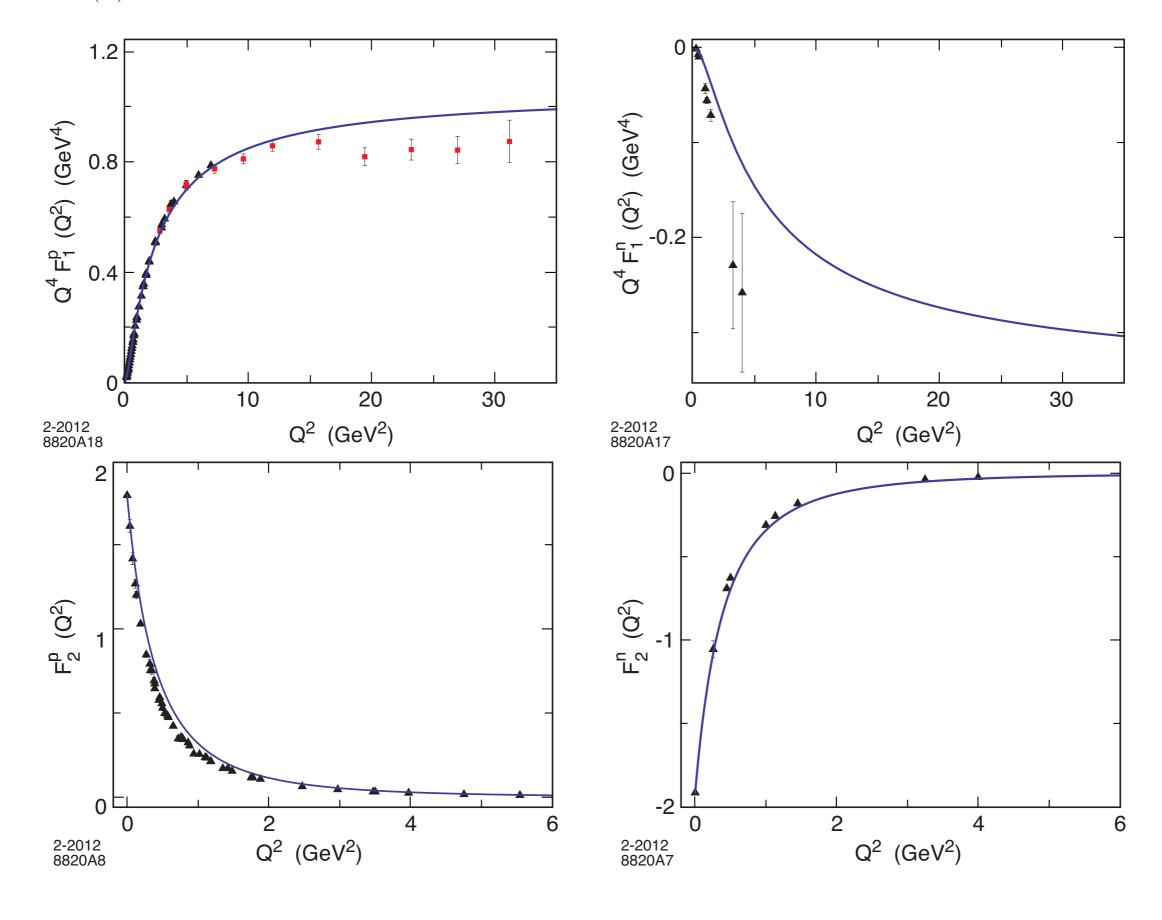
• 
$$J^z = L^z + 1/2 = (L^z + 1) - 1/2$$

$$S^z = \pm 1/2$$

- Baryon spin carried by quark orbital angular momentum: <J^z> =L^z+1/2
- Mass-degenerate meson "superpartner" with
   L_M=L_B+1. "Shifted meson-baryon Duality"

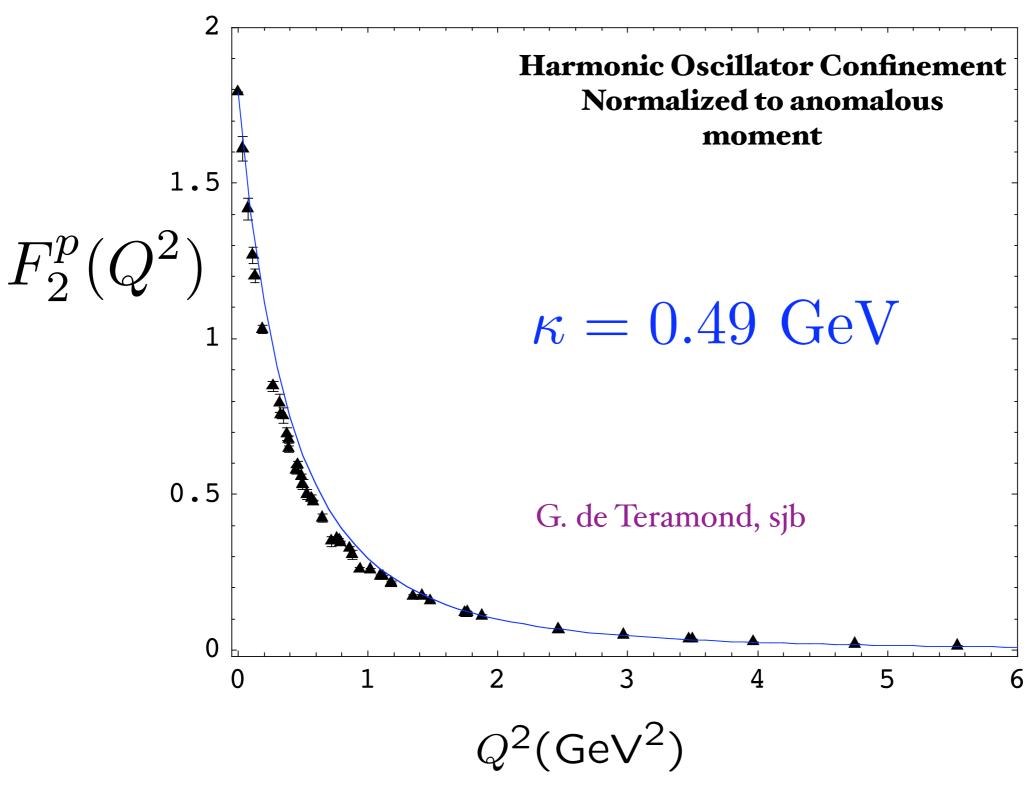
Meson and baryon have same  $\kappa$  !





# Spacelike Pauli Form Factor

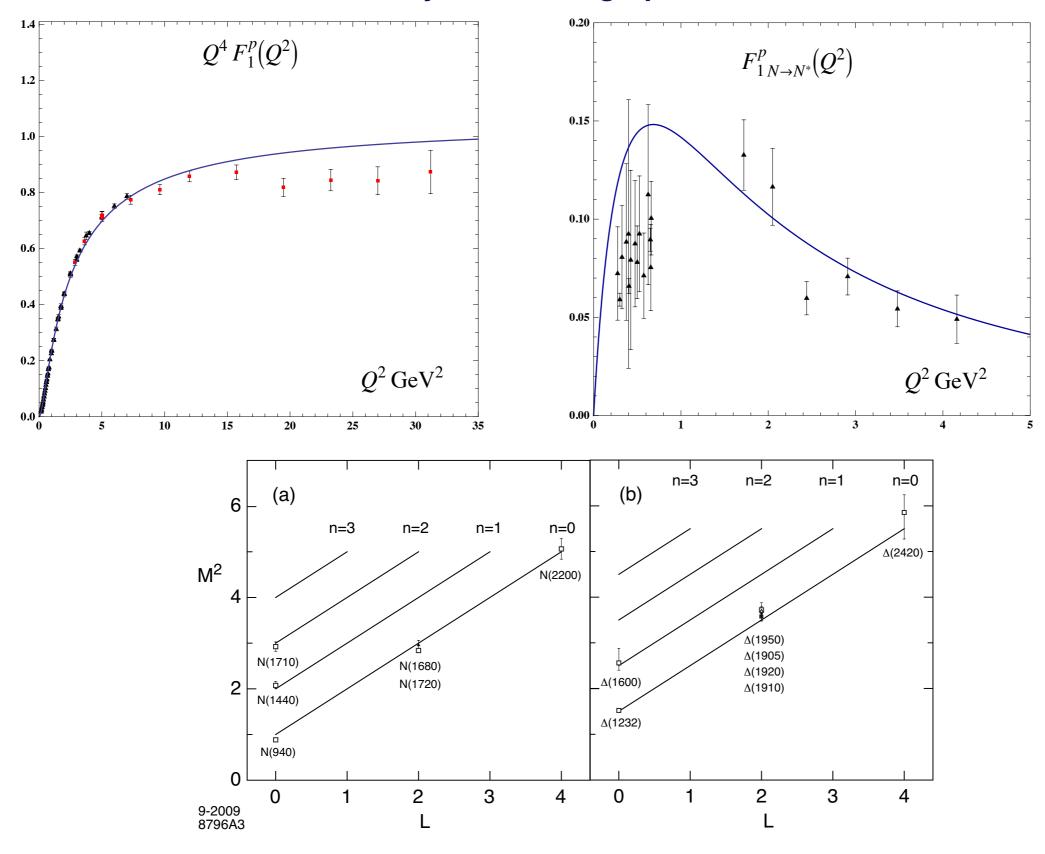
From overlap of L = 1 and L = 0 LFWFs



### Predict hadron spectroscopy and dynamics

#### **Excited Baryons in Holographic QCD**

#### G. de Teramond & sjb



#### **Nucleon Transition Form Factors**

- Compute spin non-flip EM transition  $N(940) \to N^*(1440)$ :  $\Psi^{n=0,L=0}_+ \to \Psi^{n=1,L=0}_+$
- Transition form factor

$$F_{1N \to N^*}^{p}(Q^2) = R^4 \int \frac{dz}{z^4} \Psi_{+}^{n=1,L=0}(z) V(Q,z) \Psi_{+}^{n=0,L=0}(z)$$

 $\bullet \ \ \text{Orthonormality of Laguerre functions} \quad \left( F_1{}^p_{N \to N^*}(0) = 0, \quad V(Q=0,z) = 1 \right)$ 

$$R^{4} \int \frac{dz}{z^{4}} \Psi_{+}^{n',L}(z) \Psi_{+}^{n,L}(z) = \delta_{n,n'}$$

Find

$$F_{1N \to N^*}^{p}(Q^2) = \frac{2\sqrt{2}}{3} \frac{\frac{Q^2}{M_P^2}}{\left(1 + \frac{Q^2}{\mathcal{M}_{\rho}^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho''}^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho''}^2}\right)}$$

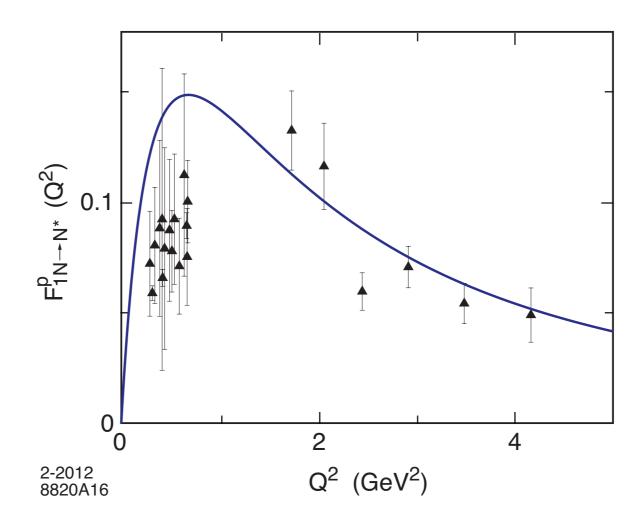
with  $\mathcal{M}_{
ho_n}^{\ 2} o 4\kappa^2(n+1/2)$ 

de Teramond, sjb

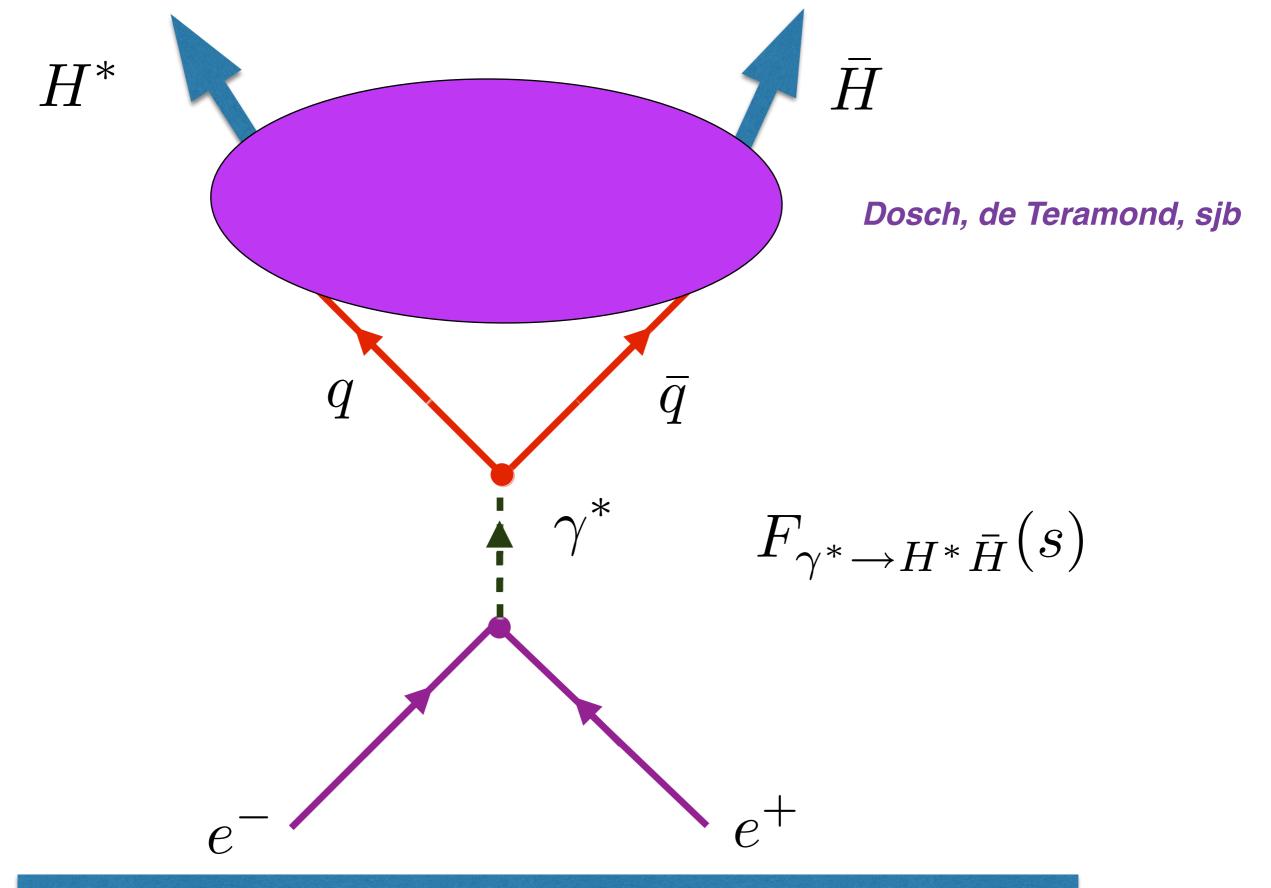
Consistent with counting rule, twist 3

#### **Nucleon Transition Form Factors**

$$F_{1 N \to N^*}^p(Q^2) = \frac{\sqrt{2}}{3} \frac{\frac{Q^2}{\mathcal{M}_{\rho}^2}}{\left(1 + \frac{Q^2}{\mathcal{M}_{\rho}^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho''}^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho''}^2}\right)}$$



Proton transition form factor to the first radial excited state. Data from JLab



Prediction from Super Conformal AdS/QCD: Same Form Factors for H=M and H=B if  $L_M=L_B+I$ 

# Bjorken sum rule defines effective charge $\alpha_{q1}(Q^2)$

$$\alpha_{g1}(Q^2)$$

$$\int_0^1 dx [g_1^{ep}(x, Q^2) - g_1^{en}(x, Q^2)] \equiv \frac{g_a}{6} [1 - \frac{\alpha_{g1}(Q^2)}{\pi}]$$

- Can be used as standard QCD coupling
- Well measured
- Asymptotic freedom at large Q²
- Computable at large Q² in any pQCD scheme
- Universal β₀, β₁

### Running Coupling from Modified AdS/QCD

#### Deur, de Teramond, sjb

ullet Consider five-dim gauge fields propagating in AdS $_5$  space in dilaton background  $arphi(z)=\kappa^2z^2$ 

$$S = -\frac{1}{4} \int d^4x \, dz \, \sqrt{g} \, e^{\varphi(z)} \, \frac{1}{g_5^2} \, G^2$$

Flow equation

$$\frac{1}{g_5^2(z)} = e^{\varphi(z)} \frac{1}{g_5^2(0)}$$
 or  $g_5^2(z) = e^{-\kappa^2 z^2} g_5^2(0)$ 

where the coupling  $g_5(z)$  incorporates the non-conformal dynamics of confinement

- YM coupling  $\alpha_s(\zeta) = g_{YM}^2(\zeta)/4\pi$  is the five dim coupling up to a factor:  $g_5(z) \to g_{YM}(\zeta)$
- Coupling measured at momentum scale Q

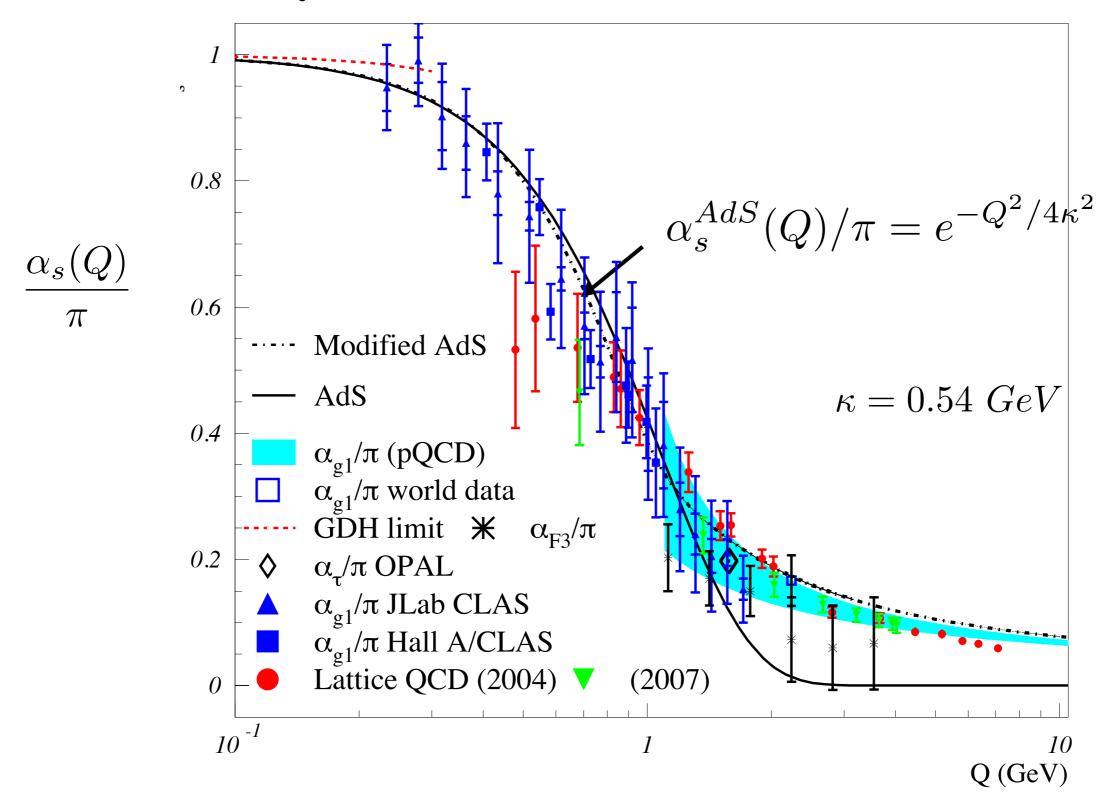
$$\alpha_s^{AdS}(Q) \sim \int_0^\infty \zeta d\zeta J_0(\zeta Q) \,\alpha_s^{AdS}(\zeta)$$

Solution

$$\alpha_s^{AdS}(Q^2) = \alpha_s^{AdS}(0) e^{-Q^2/4\kappa^2}$$

 $\alpha_s^{AdS}(Q^2)=\alpha_s^{AdS}(0)\,e^{-Q^2/4\kappa^2}.$  where the coupling  $\alpha_s^{AdS}$  incorporates the non-conformal dynamics of confinement

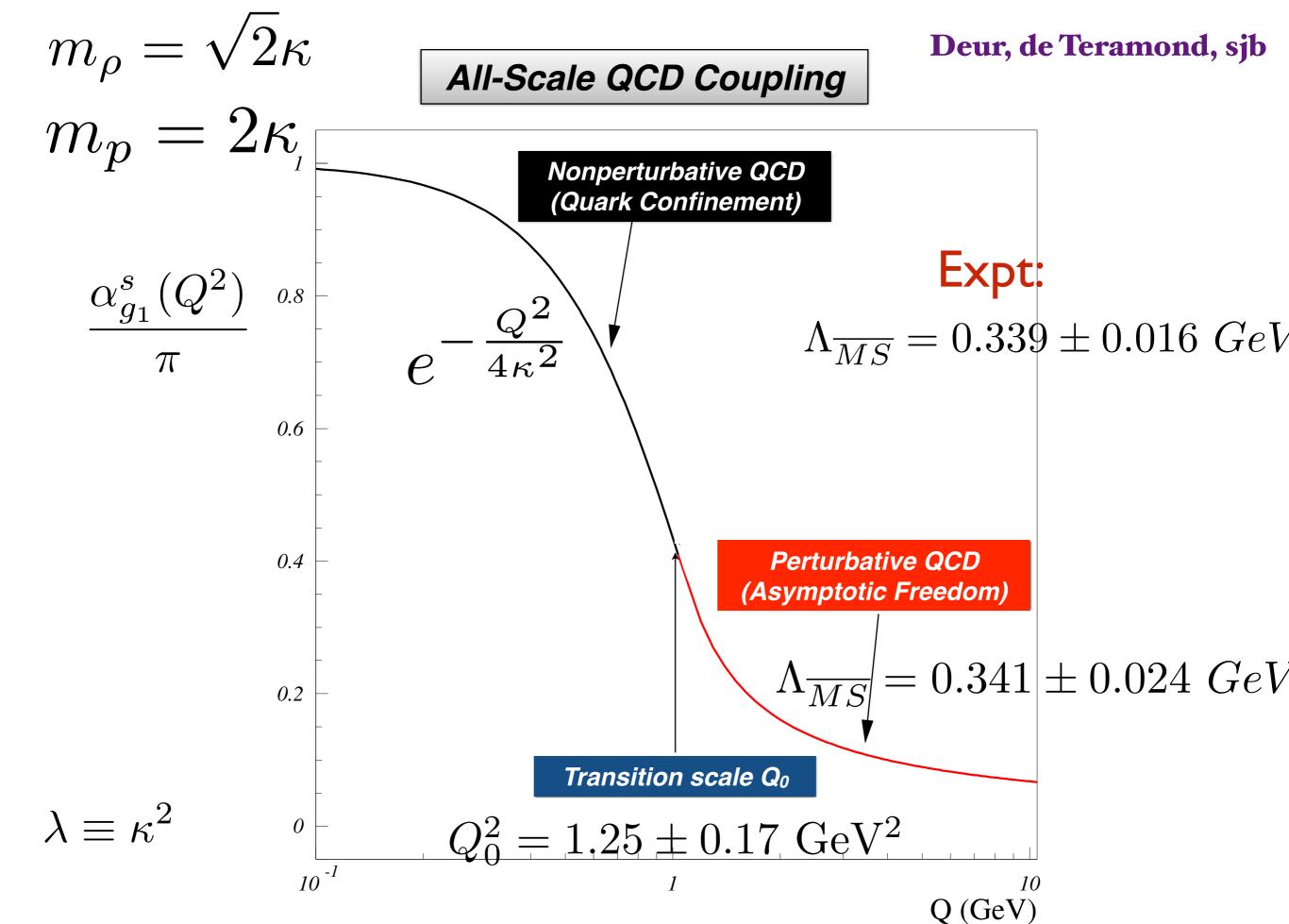
### Analytic, defined at all scales, IR Fixed Point

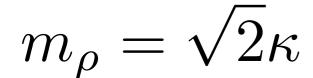


AdS/QCD dilaton captures the higher twist corrections to effective charges for Q < 1 GeV

$$e^{\varphi} = e^{+\kappa^2 z^2}$$

Deur, de Teramond, sjb

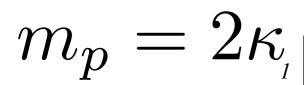




#### All-Scale QCD Coupling

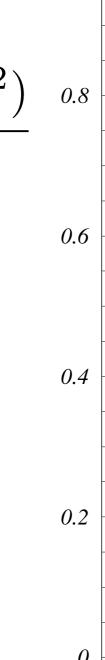
Nonperturbative QCD

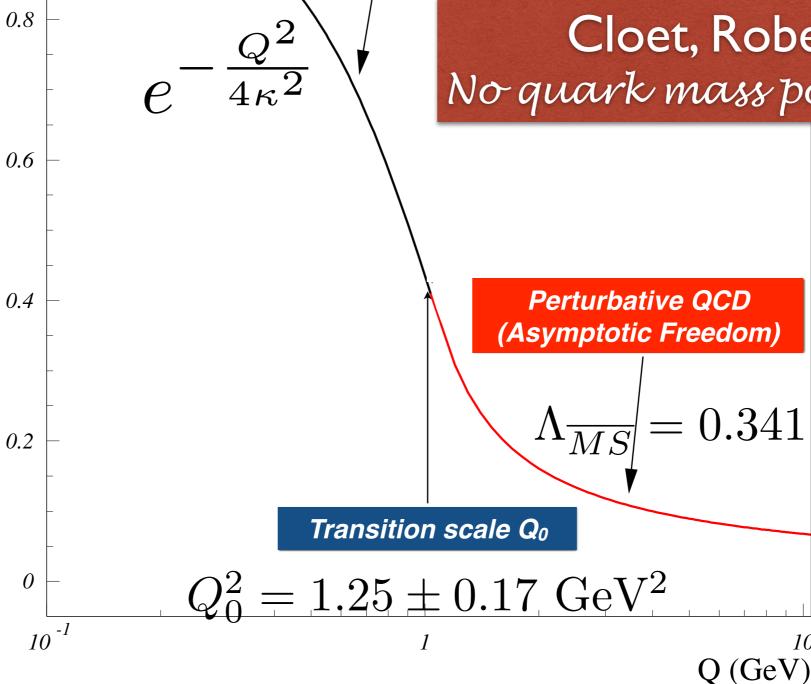
(Quark Confinement)



$$\frac{lpha_{g_1}^s(Q^2)}{\pi}$$
 0.8

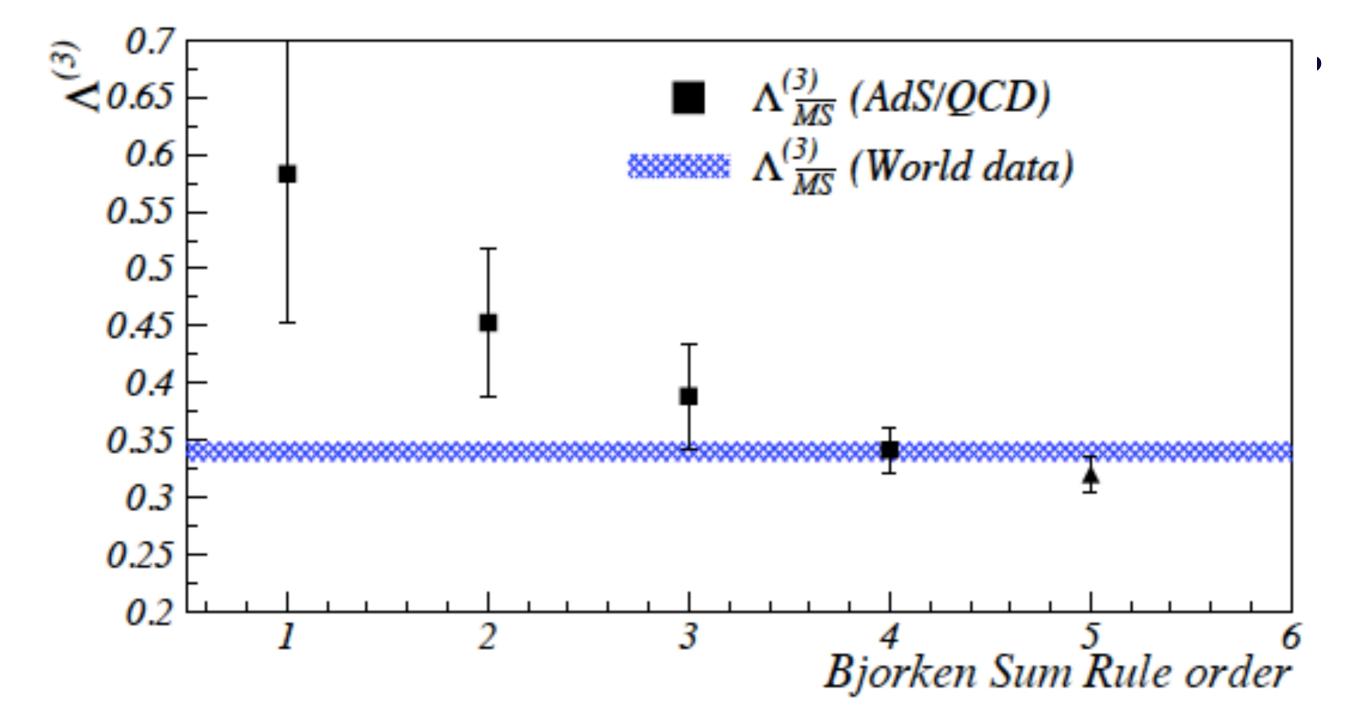
 $\lambda \equiv \kappa^2$ 





No quark mass pole in DSE

$$\Lambda_{\overline{MS}} = 0.341 \pm 0.024 \; GeV$$



$$\Lambda_{\overline{MS}} = 0.341 \ GeV = 0.440 m_{\rho} = 0.622 \kappa$$

Connect  $\Lambda_{\overline{MS}}$  to hadron masses!

Experiment:  $M_{o} = 0.7753 \pm 0.0003 \; GeV$ 

# Interpretation of Mass Scale K

- Does not affect conformal symmetry of QCD action
- Self-consistent regularization of IR divergences
- Determines all mass and length scales for zero quark mass
- ullet Compute scheme-dependent  $\Lambda_{\overline{MS}}$  determined in terms of  $\,\,\mathcal{K}$
- Value of  $\kappa$  itself not determined place holder
- Need external constraint such as  $f_{\pi}$

# AdS/QCD and Light-Front Holography

$$\mathcal{M}_{n,J,L}^2 = 4\kappa^2 \left(n + \frac{J+L}{2}\right)$$

- Zero mass pion for m_q = o (n=J=L=o)
- Regge trajectories: equal slope in n and L
- Form Factors at high Q²: Dimensional counting  $[Q^2]^{n-1}F(Q^2) \to \text{const}$
- Space-like and Time-like Meson and Baryon Form Factors
- Running Coupling for NPQCD  $\alpha_s$

$$\alpha_s(Q^2) \propto e^{-\frac{Q^2}{4\kappa^2}}$$

Meson Distribution Amplitude

$$\phi_{\pi}(x) \propto f_{\pi} \sqrt{x(1-x)}$$



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#### de Tèramond, Dosch, Deur, sjb

# Features of AdS/QCD

- Color confining potential  $\kappa^4\zeta^2$  and universal mass scale from dilaton  $e^{\phi(z)} = e^{\kappa^2z^2} \qquad \alpha_s(Q^2) \propto \exp{-Q^2/4\kappa^2}$
- Dimensional transmutation  $\Lambda_{\overline{MS}} \leftrightarrow \kappa \leftrightarrow m_H$
- Chiral Action remains conformally invariant despite mass scale DAFF
- Light-Front Holography: Duality of AdS and frame-independent LF QCD
- Reproduces observed Regge spectroscopy same slope in n, L, and J for mesons and baryons
- Massless pion for massless quark
- Supersymmetric meson-baryon dynamics and spectroscopy:  $L_{M=L_B+1}$

• Dynamics: LFWFs, Form Factors, GPDs

Superconformal Algebra Fubini and Rabinovici

# An analytic first approximation to QCD

AdS/QCD + Light-Front Holography

- As Simple as Schrödinger Theory in Atomic Physics
- LF radial variable ζ conjugate to invariant mass squared
- Relativistic, Frame-Independent, Color-Confining
- Unique confining potential!
- QCD Coupling at all scales: Essential for Gauge Link phenomena
- Hadron Spectroscopy and Dynamics from one parameter
- Wave Functions, Form Factors, Hadronic Observables, Constituent Counting Rules
- Insight into QCD Condensates: Zero cosmological constant!
- Systematically improvable with DLCQ-BLFQ Methods

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August 31, 2015

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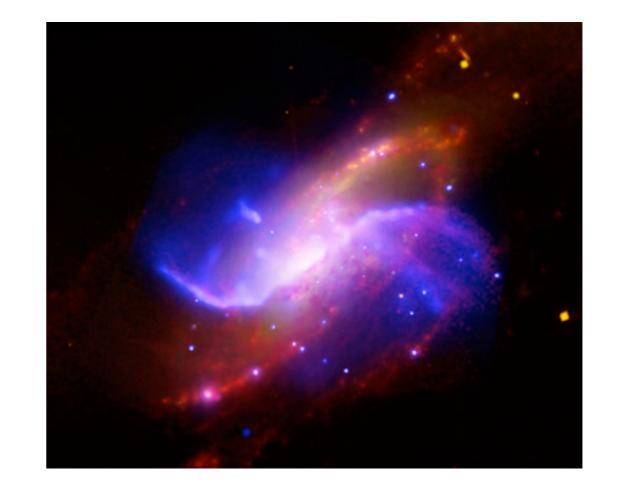


# Future Directions for AdS/QCD

- Hadronization at the Amplitude Level
- Diffractive dissociation of pion and proton to jets
- Identify the factorization Scale for ERBL, DGLAP evolution: Qo
- Compute Tetraquark Spectroscopy Sequentially
- Update SU(6) spin-flavor symmetry
- Heavy Quark States: Supersymmetry, not conformal
- Compute higher Fock states; e.g. Intrinsic Heavy Quarks
- Nuclear States Hidden Color
- Basis LF Quantization

We view the universe as light reaches us along the light-front at fixed

$$\tau = t + z/c$$



Front Form Vacuum Describes the Empty, Causal Universe

# "One of the gravest puzzles of theoretical physics"

DARK ENERGY AND
THE COSMOLOGICAL CONSTANT PARADOX

A. ZEE

Department of Physics, University of California, Santa Barbara, CA 93106, USA
Kavil Institute for Theoretical Physics, University of California,
Santa Barbara, CA 93106, USA
zee@kitp.ucsb.edu

$$(\Omega_{\Lambda})_{QCD} \sim 10^{45}$$

$$\Omega_{\Lambda} = 0.76(expt)$$

$$(\Omega_{\Lambda})_{EW} \sim 10^{56}$$

# Extraordinary conflict between the conventional definition of the vacuum in quantum field theory and cosmology

Elements of the solution:

(A) Light-Front Quantization: causal, frame-independent vacuum (B) New understanding of QCD "Condensates" (C) Higgs Light-Front Zero Mode

# Two Definitions of Vacuum State

### Instant Form: Lowest Energy Eigenstate of Instant-Form Hamiltonian

$$H|\psi_0>=E_0|\psi_0>, E_0=\min\{E_i\}$$

## Eigenstate defined at one time t over all space; Acausal! Frame-Dependent

#### Front Form: Lowest Invariant Mass Eigenstate of Light-Front Hamiltonian

$$H_{LF}|\psi_0>_{LF}=M_0^2|\psi_0>_{LF}, M_0^2=0.$$

Frame-independent eigenstate at fixed LF time \tau = t+z/c within causal horizon

Frame-independent description of the causal physical universe!

# Light-Front vacuum can simulate empty universe Shrock, Tandy, Roberts, sjb

- Independent of observer frame
- Causal
- Lowest invariant mass state M= o.
- Trivial up to k⁺=o zero modes-- already normal-ordering
- Higgs theory consistent with trivial LF vacuum (Srivastava, sjb)
- QCD and AdS/QCD: "In-hadron" condensates (Maris, Tandy Roberts) -- GMOR satisfied.
- QED vacuum; no loops
- Zero cosmological constant from QED, QCD







# Standard Model on the Light-Front

- Same phenomenological predictions
- Higgs field has three components
- Real part creates Higgs particle
- Imaginary part (Goldstone) become longitudinal components of W, Z
- Higgs VEV of instant form becomes k+=0 LF zero mode!
- Analogous to a background static classical Zeeman or Stark Fields
- Zero contribution to  $T^{\mu}_{\mu}$ ; zero coupling to gravity



# Goals

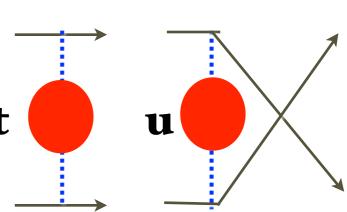
- Test QCD to maximum precision at the LHC
- Maximize sensitivity to new physics
- High precision determination of fundamental parameters
- Determine renormalizations scales without ambiguity
- Eliminate scheme dependence

Predictions for physical observables cannot depend on theoretical conventions such as the renormalization scheme

# Electron-Electron Scattering in QED

$$\mathcal{M}_{ee \to ee}(++;++) = \frac{8\pi s}{t} \alpha(t) + \frac{8\pi s}{u} \alpha(u)$$

- Two separate physical scales: t, u = photon virtuality
- Gauge Invariant. Dressed photon propagator



- Sums all vacuum polarization, non-zero beta terms into running coupling. This is the purpose of the running coupling!
- If one chooses a different initial scale, one must sum an infinite number of graphs -- but always recover same result!
- Number of active leptons correctly set
- Analytic: reproduces correct behavior at lepton mass thresholds
- No renormalization scale ambiguity!

# $\delta$ - $\mathcal{R}$ enormalization Scheme ( $\mathcal{R}_{\delta}$ scheme)

In dim. reg.  $1/\epsilon$  poles come in powers of [Bollini & Gambiagi, 't Hooft & Veltman, '72]

$$\ln\frac{\mu^2}{\Lambda^2} + \frac{1}{\epsilon} + c$$

In the modified minimal subtraction scheme (MS-bar) one subtracts together with the pole a constant [Bardeen, Buras, Duke, Muta (1978) on DIS results]:

$$\ln(4\pi) - \gamma_E$$

This corresponds to a shift in the scale:

$$\mu_{\overline{\rm MS}}^2 = \mu^2 \exp(\ln 4\pi - \gamma_E)$$

A finite subtraction from infinity is arbitrary. Let's make use of this!

Subtract an arbitrary constant and keep it in your calculation:  $\mathcal{R}_{\delta}$ -scheme

M. Mojaza, Xing-Gang Wu, sjb

$$\ln(4\pi) - \gamma_E - \delta,$$

$$\mu_{\delta}^2 = \mu_{\overline{MS}}^2 \exp(-\delta) = \mu^2 \exp(\ln 4\pi - \gamma_E - \delta)$$

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August 31, 2015

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# Exposing the Renormalization Scheme Dependence

#### Observable in the $\mathcal{R}_{\delta}$ -scheme:

$$\rho_{\delta}(Q^2) = r_0 + r_1 a(\mu) + [r_2 + \beta_0 r_1 \delta] a(\mu)^2 + [r_3 + \beta_1 r_1 \delta + 2\beta_0 r_2 \delta + \beta_0^2 r_1 \delta^2] a(\mu)^3 + \cdots$$

$$\mathcal{R}_0 = \overline{\text{MS}}$$
,  $\mathcal{R}_{\ln 4\pi - \gamma_E} = \text{MS}$   $\mu^2 = \mu_{\overline{\text{MS}}}^2 \exp(\ln 4\pi - \gamma_E)$ ,  $\mu_{\delta_2}^2 = \mu_{\delta_1}^2 \exp(\delta_2 - \delta_1)$ 

Note the divergent 'renormalon series'  $n!\beta^n\alpha_s^n$ 

#### Renormalization Scheme Equation

$$\frac{d\rho}{d\delta} = -\beta(a)\frac{d\rho}{da} \stackrel{!}{=} 0 \longrightarrow PMC$$

$$\rho_{\delta}(Q^2) = r_0 + r_1 a_1(\mu_1) + (r_2 + \beta_0 r_1 \delta_1) a_2(\mu_2)^2 + [r_3 + \beta_1 r_1 \delta_1 + 2\beta_0 r_2 \delta_2 + \beta_0^2 r_1 \delta_1^2] a_3(\mu_3)^3$$

The  $\delta_k^p a^n$ -term indicates the term associated to a diagram with  $1/\epsilon^{n-k}$  divergence for any p. Grouping the different  $\delta_k$ -terms, one recovers in the  $N_c \to 0$  Abelian limit the dressed skeleton expansion.

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# Special Degeneracy in PQCD

There is nothing special about a particular value for  $\,\delta$  , thus for any  $\,\delta$ 

$$\rho(Q^{2}) = r_{0,0} + r_{1,0}a(Q) + [r_{2,0} + \beta_{0}\underline{r_{2,1}}]a(Q)^{2} + [r_{3,0} + \beta_{1}\underline{r_{2,1}} + 2\beta_{0}\underline{r_{3,1}} + \beta_{0}^{2}\underline{r_{3,2}}]a(Q)^{3} + [r_{4,0} + \beta_{2}\underline{r_{2,1}} + 2\beta_{1}\underline{r_{3,1}} + \frac{5}{2}\beta_{1}\beta_{0}\underline{r_{3,2}} + 3\beta_{0}r_{4,1} + 3\beta_{0}^{2}r_{4,2} + \beta_{0}^{3}r_{4,3}]a(Q)^{4}$$

According to the principal of maximum conformality we must set the scales such to absorb all 'renormalon-terms', i.e. non-conformal terms

$$\rho(Q^{2}) = r_{0,0} + r_{1,0}a(Q) + (\beta_{0}a(Q)^{2} + \beta_{1}a(Q)^{3} + \beta_{2}a(Q)^{4} + \cdots)r_{2,1}$$

$$+ (\beta_{0}^{2}a(Q)^{3} + \frac{5}{2}\beta_{1}\beta_{0}a(Q)^{4} + \cdots)r_{3,2} + (\beta_{0}^{3} + \cdots)r_{4,3}$$

$$+ r_{2,0}a(Q)^{2} + 2a(Q)(\beta_{0}a(Q)^{2} + \beta_{1}a(Q)^{3} + \cdots)r_{3,1}$$

$$+ \cdots$$

$$r_{1,0}a(Q_1) = r_{1,0}a(Q) - \beta(a)r_{2,1} + \frac{1}{2}\beta(a)\frac{\partial\beta}{\partial a}r_{3,2} + \dots + \frac{(-1)^n}{n!}\frac{d^{n-1}\beta}{(d\ln\mu^2)^{n-1}}r_{n+1,n}$$

$$r_{2,0}a(Q_2)^2 = r_{2,0}a(Q)^2 - 2a(Q)\beta(a)r_{3,1} + \cdots$$

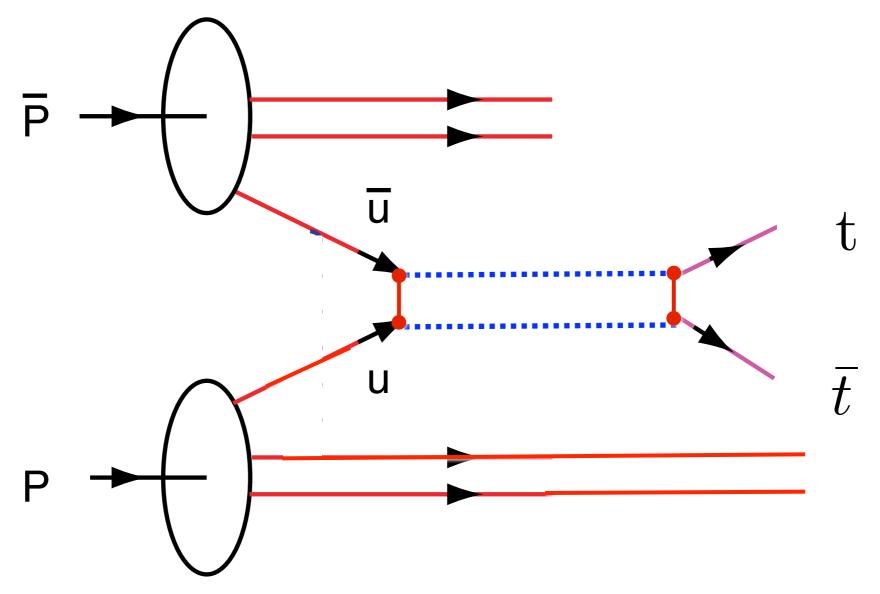
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Implications for the  $\bar{p}p \to t\bar{t}X$  asymmetry at the Tevatron



Interferes with Born term.

Small value of renormalization scale increases asymmetry, just as in QED

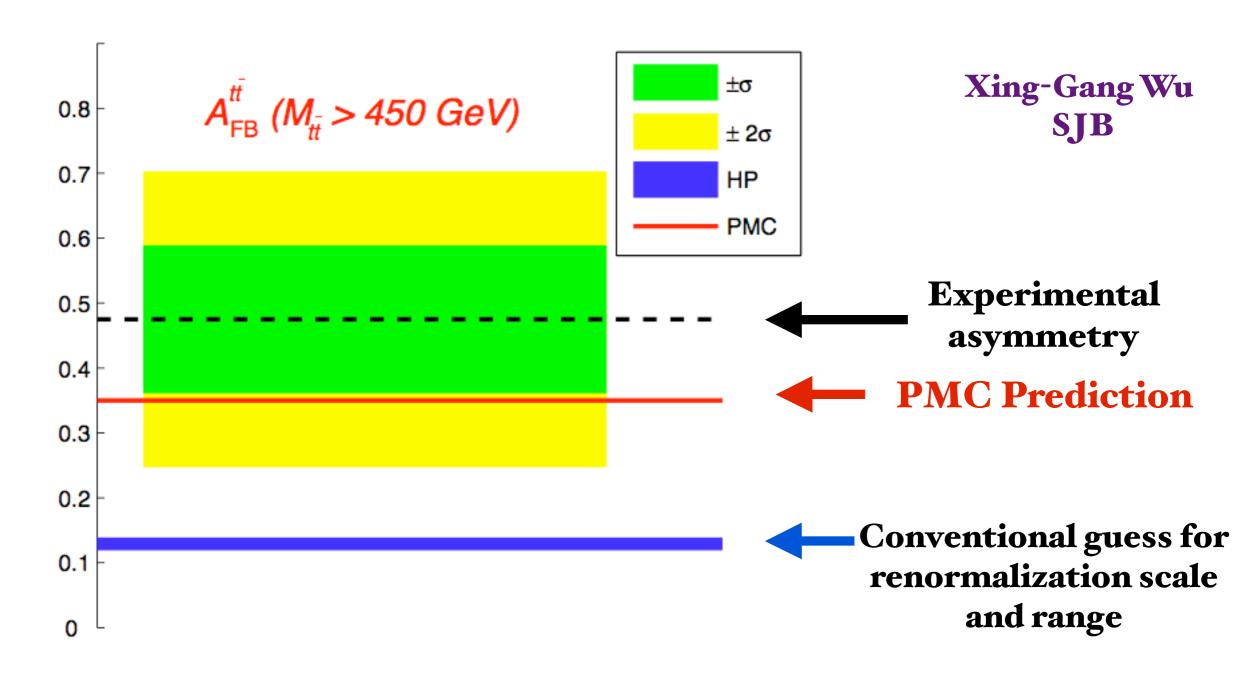
Xing-Gang Wu, sjb Stan Brodsky

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Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD



The Renormalization Scale Ambiguity for Top-Pair Production Eliminated Using the 'Principle of Maximum Conformality' (PMC)

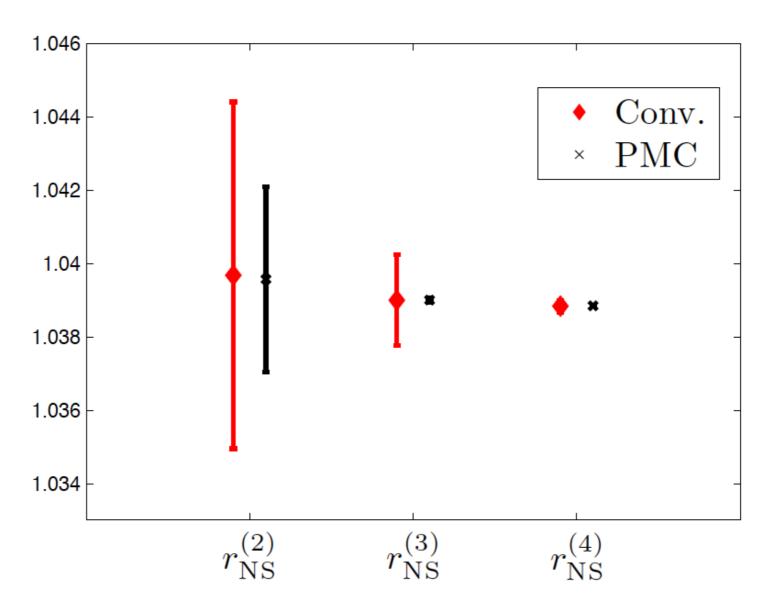


Top quark forward-backward asymmetry predicted by pQCD NNLO within 1  $\sigma$  of CDF/D0 measurements using PMC/BLM scale setting

# Reanalysis of the Higher Order Perturbative QCD corrections to Hadronic Z Decays using the Principle of Maximum Conformality

S-Q Wang, X-G Wu, sjb

P.A. Baikov, K.G. Chetyrkin, J.H. Kuhn, and J. Rittinger, Phys. Rev. Lett. 108, 222003 (2012).



The values of  $r_{\rm NS}^{(n)}=1+\sum_{i=1}^n C_i^{\rm NS} a_s^i$  and their errors  $\pm |C_n^{\rm NS} a_s^n|_{\rm MAX}$ . The diamonds and the crosses are for conventional (Conv.) and PMC scale settings, respectively. The central values assume the initial scale choice  $\mu_r^{\rm init}=M_Z$ .

# What is PMC?

# Principle of Maximum Conformality

Choose renormalization scheme; e.g.  $\alpha_s^R(\mu_R^{\text{init}})$ 

Choose  $\mu_R^{init}$ ; arbitrary initial renormalization scale

Identify  $\{\beta_i^R\}$  – terms using  $n_f$  – terms

 $through \ the \ PMC-BLM \ correspondence \ principle$ 

order-by-order

Shift scale of  $\alpha_s$  to  $\mu_R^{\text{PMC}}$  to eliminate  $\{\beta_i^R\}$  – terms

Conformal Series

Xing-Gang Wu, Matin Mojaza Leonardo di Giustino, SJB

PMC-BLM - one

Phys. Rev. Lett. 109, 042002 (2012)

 $R_{\delta}$ -scheme – two

Phys. Rev. Lett. 110, 192001 (2013)

Eliminate β-terms

Result is independent of  $\mu_R^{\text{init}}$  and scheme at fixed order

 $n_f$  dependence of pQCD series does not uniquely identify the  $\beta$  terms

# Features of BLM/PMC

- Predictions are scheme-independent
- Matches conformal series
- Commensurate Scale Relations between observables: Generalized Crewther Relation (Kataev, Lu, Rathsman, sjb)
- No n! Renormalon growth
- New scale at each order; n_F determined at each order
- Multiple Physical Scales Incorporated
- Rigorous: Satisfies all Renormalization Group Principles
- Realistic Estimate of Higher-Order Terms
- Eliminates unnecessary theory error

# QCD Myths

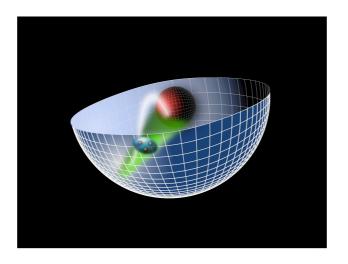
- Anti-Shadowing is Universal
- ISI and FSI are higher twist effects and universal
- High transverse momentum hadrons arise only from jet fragmentation -- baryon anomaly!
- Heavy quarks only from gluon splitting
- Renormalization scale cannot be fixed
- QCD condensates are vacuum effects
- QCD gives 1042 to the cosmological constant



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AdS/QCD Soft-Wall Model

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$



$$\zeta^2 = x(1-x)\mathbf{b}_{\perp}^2.$$

Light-Front Holography

$$\left[ -\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



#### Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

 $\kappa \simeq 0.6 \; GeV$ 

# Confinement scale: $\kappa \simeq$

$$1/\kappa \simeq 1/3 \ fm$$

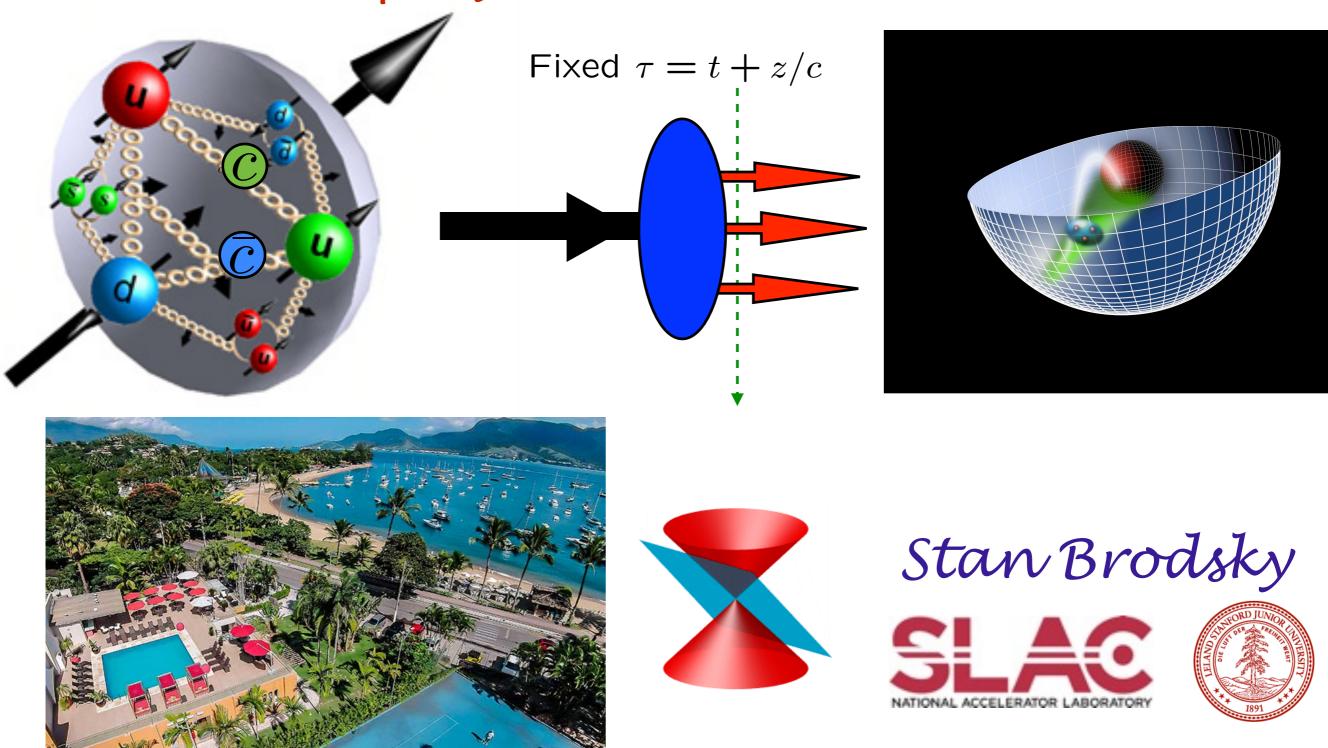
- de Alfaro, Fubini, Furlan:
  - Fubini, Rabinovici:

Unique Confinement Potential!

Preserves Conformal Symmetry of the action

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

# Light-Front Holographic QCD, Color Confinement, and Supersymmetric Features of QCD



QCD-TNT4 Unraveling the Organization of the Tapestry of QCD

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