

# Introduction to Quantum Walks – part 3

Paweł Kurzyński 4-18 June 2024

## **QW** applications



Xiaogang Qiang, Shixin Ma, and Haijing Song arXiv:2404.04178 (2024)

#### QW search algorithms

- QW spread faster than RW ballistic O(t) vs diffusive  $O(\sqrt{t})$
- In some graphs spread is even faster
- This can be used to speed up search

Neil Shenvi, Julia Kempe, and K. Birgitta Whaley Phys. Rev. A 67, 052307 (2003)

• Structured database – graph

Miklos Santha, TAMC, Lecture Notes in Computer Science, vol 4978. Springer, Berlin, Heidelberg (2008)

• Idea: Start QW in a uniform state

$$|\psi(0)
angle = rac{1}{\sqrt{N}}\sum_{x=1}^{N}|x
angle \otimes |anc
angle$$

Find evolution to localise walk in the marked vertex x = m

$$|\psi(t)
angle pprox |m
angle \otimes |anc'
angle$$

## QW search algorithms

#### Grover algorithm as $\mathsf{Q}\mathsf{W}$

Lov K. Grover, Am. J. Phys. 69, 769-777 (2001)

$$U = U_1 U_0, \qquad U_0 |x\rangle = e^{i\pi\delta_{x,m}} |x\rangle, \qquad U_1 = 2|\psi\rangle\langle\psi| - I,$$
$$|\psi\rangle = \frac{1}{\sqrt{n}} \sum_{x=1}^N |x\rangle, \qquad U_0 |x'\rangle = \frac{2}{N} \sum_{x=1}^N |x\rangle - |x'\rangle$$



- In a connected graph the search time is  $O(\sqrt{N})$
- In a structured database (not a connected graph) the walker cannot get straight to the marked vertex
- Structured graph search time is typically  $\sqrt{N} \leq t \leq O(N)$
- Example: search in 1D (chain) does not exhibit quadratic speedup t = O(N)

#### QW as Universal Quantum Computers

CTQW as Universal Quantum Computer

 $H \rightarrow CNOT$ 



Andrew M Childs, Phys. Rev. Lett. 102, 180501 (2009)

#### QW as Universal Quantum Computers

DTQW as Universal Quantum Computer



Neil B Lovett, et. al., Phys. Rev. A 81, 042330 (2010)

QW simulation of solid state systems

- Graph tight binding and crystal structure
- How to introduce electric potentials  $\phi(x)$ ?
- Dirac analogy

$$U = S(I \otimes C) = e^{ip\sigma_z} e^{i\theta\sigma_y} \rightarrow e^{-i(cp\sigma_z + mc^2\sigma_y)t}$$
$$e^{-i(cp\sigma_z + mc^2\sigma_y + \phi(x))t} \rightarrow e^{ip\sigma_z} e^{i\theta\sigma_y} e^{i\phi(x)} = S(I \otimes C)P = U'$$
$$P = \left(\sum_{x = -\infty}^{\infty} e^{i\phi(x)} |x\rangle\langle x|\right) \otimes I$$

• Constant force:  $\phi(x) = \varphi x$ 



D. Bouwmeester, et. al., Phys. Rev. A 61, 013410 (1999)

A. Wojcik, et. al., Phys. Rev. Lett. 93, 180601 (2004)

#### Bloch oscillations



- Constant force  $e^{i\eta x}|k
  angle = |k+\eta
  angle$
- Oscillations due to momentum in the first Brillouin zone  $k \in [-\pi,\pi]$
- Period of Bloch oscillations (unit mass)  $T_{BO} = \frac{2\pi}{n}$
- Period of Bloch oscillations ( $N \times$  unit mass)  $T_{BO} = \frac{2\pi}{N\eta}$

A. Wojcik, et. al., Phys. Rev. Lett. 93, 180601 (2004)

#### Bloch oscillations of two quantum walkers (center of mass)



Z. Lasmar, et. al., Phys. Rev. A 98, 062105 (2018)

#### Magnetic quantum walks (2D)

I. Yalcınkaya and Z. Gedik, Phys. Rev. A 92, 042324 (2015)

• How to introduce vector potentials  $\vec{A}(x, z)$ ?

• 
$$\vec{p} \rightarrow \vec{p} - q\vec{A}$$

• Dirac analogy

$$U = S_{x}(I \otimes C)S_{z}(I \otimes C) = e^{ip_{x}\sigma_{x}}e^{i\theta\sigma_{y}}e^{ip_{z}\sigma_{z}}e^{i\theta\sigma_{y}} \rightarrow e^{-i(cp_{x}\sigma_{x}+cp_{z}\sigma_{z}+mc^{2}\sigma_{y})t}$$
$$e^{-i(c(p_{x}-qA_{x})\sigma_{x}+c(p_{z}-qA_{z})\sigma_{z}+mc^{2}\sigma_{y})t} \rightarrow$$

 $e^{ip_{x}\sigma_{x}}e^{iA_{x}\sigma_{x}}e^{i\theta\sigma_{y}}e^{ip_{z}\sigma_{z}}e^{iA_{x}\sigma_{x}}e^{i\theta\sigma_{y}} = S_{x}M_{x}(I \otimes C)S_{z}M_{z}(I \otimes C) = U'$ 

• Position dependent coin  $(A_x = A_x(x, z), A_z = A_z(x, z))$ :

$$M_{x} = I \otimes C_{x}(x, z), \qquad M_{z} = I \otimes C_{z}(x, z)$$
$$C_{x}(x, z) = e^{iA_{x}\sigma_{x}} = \begin{pmatrix} \cos A_{x}(x, z) & i \sin A_{x}(x, z) \\ i \sin A_{x}(x, z) & \cos A_{x}(x, z) \end{pmatrix}$$
$$C_{z}(x, z) = e^{iA_{z}\sigma_{z}} = \begin{pmatrix} e^{iA_{z}(x, z)} & 0 \\ 0 & e^{-iA_{z}(x, z)} \end{pmatrix}$$

• Constant magnetic field perpendicular to the plane (y-direction):

$$A_x = -\varphi z, \quad A_z = \varphi x$$

## Cyclotron orbit (see video) and Hofstadter butterfly



#### Topological effects

T. Kitagawa, Quant. Inf. Proc. 11, 1107-1148 (2012)





Symmetry protection of topological states









Multipartite quantum walks

- State:  $|x_1,c_1
  angle\otimes|x_2,c_2
  angle$
- Evolution:  $S_1(I \otimes C_1)S_2(I \otimes C_2)I(x_1, x_2)$
- Example (hard core interaction)

$$I(x_1, x_2) = \left(\sum_{x_1, x_2} e^{i\varphi\delta_{x_1, x_2}} |x_1, x_2\rangle \langle x_1, x_2|\right) \otimes I$$

• Creation of bound states



A. Ahlbrecht et al. New J. Phys. 14, 073050 (2012)

#### Borromean states



- State:  $|x_1,c_1
  angle\otimes|x_2,c_2
  angle\otimes|x_3,c_3
  angle$
- Evolution:  $S_1S_2S_3(I \otimes C(x_1, x_2, x_3))$
- Bound state:  $\sum_{x} |x, x, x\rangle \otimes |GHZ\rangle$

Marcin Markiewicz, Marcin Karczewski, Pawel Kurzynski, Quantum 5, 523 (2021)

- One-particle quantum walks: can be simulated classically (computers, coherent light)
- Classical: multipartite walks → cellular automata (Lattice Boltzmann Automata, Conway's Life, etc.)
- Quantum: multipartite walks  $\rightarrow$  quantum cellular automata Pablo Arrighi, arXiv:1904.12956 (2019)
- Cellular automata require powerful classical computers
- Quantum cellular automata require powerful quantum computers!



Nature Communications 3, 882, 2012

#### Atom in optical lattice





position

Cesium ground states hyperfine sublevels

F = 3, 4



Science 325 (2009) 174



PRL 120 (2018) 260501; Science & Applications (2020) 9:7



PRL 121 (2018) 070402





PRX 7 (2017) 031023



PRA 72 (2005) 062317

