



Measuring $\sigma_{p\text{-air}}$ and σ_{pp} with the Pierre Auger Observatory

***Carola Dobrigkeit for the Pierre Auger Collaboration
RETINHA 2014***

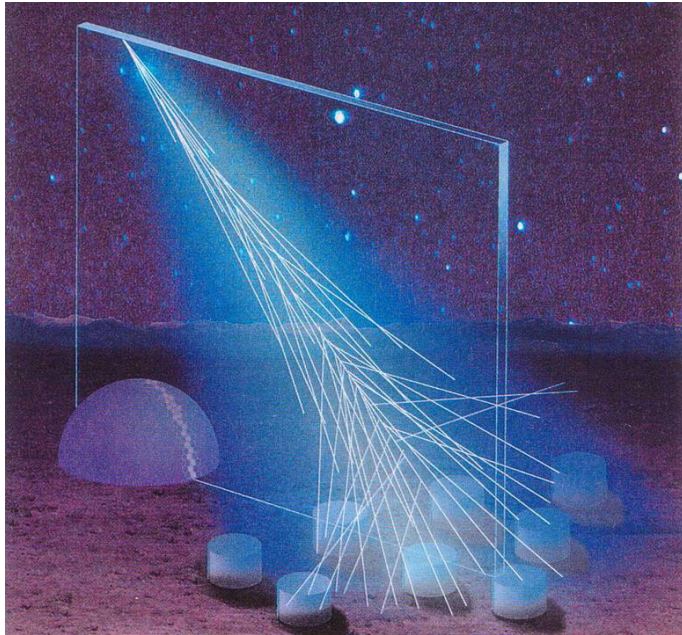
How we detect these UHE cosmic rays?

- **Primary cosmic rays are detected via the extensive air showers they induce by interacting in the atmosphere.**



- **The Pierre Auger Observatory combines two complementary techniques for detection (hybrid detection).**

Hybrid detection:



Combining both techniques allows:

- cross calibration in energy
- better angular resolution

Fluorescence Detector:

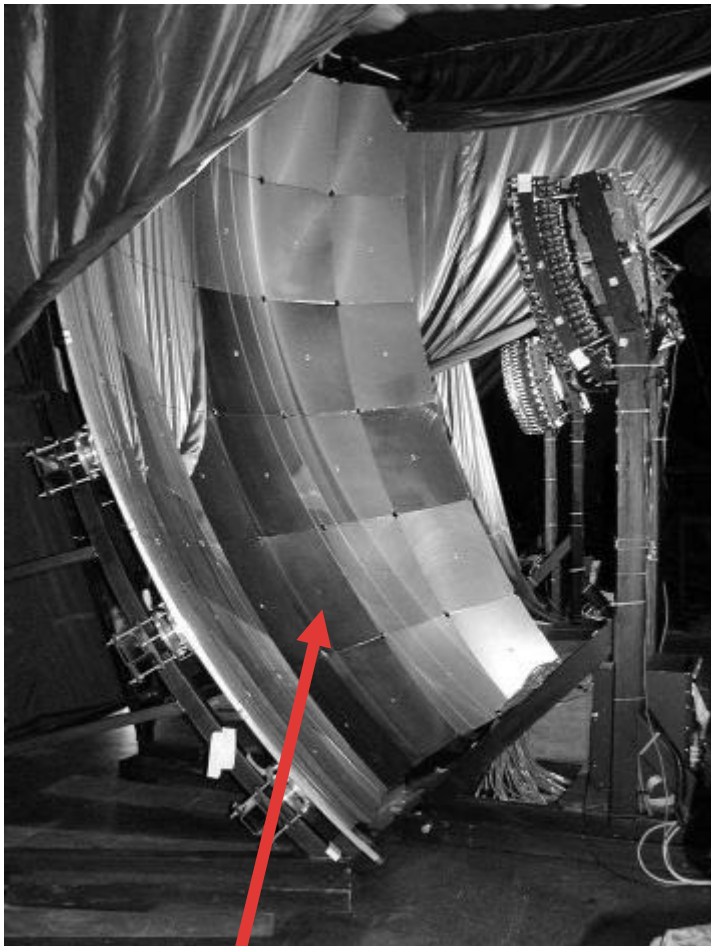
- Almost calorimetric energy measurement
- Longitudinal development
- 10-15% duty cycle
- Complex acceptance calculation

Surface Detector Array:

- 100% duty cycle
- Simple geometrical acceptance
- Extracting primary energy and mass is model dependent

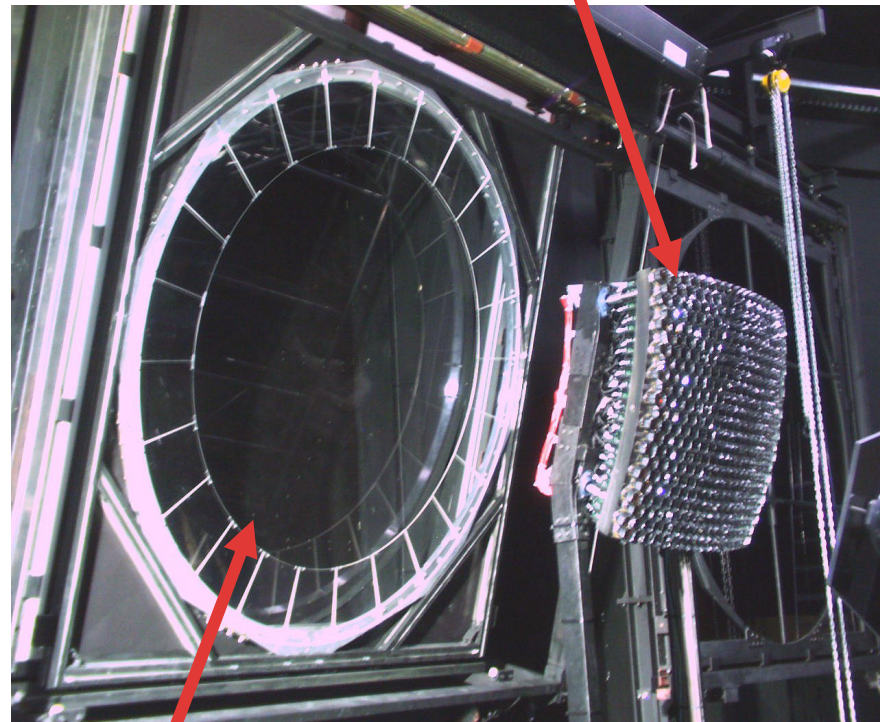


A fluorescence telescope (FD)



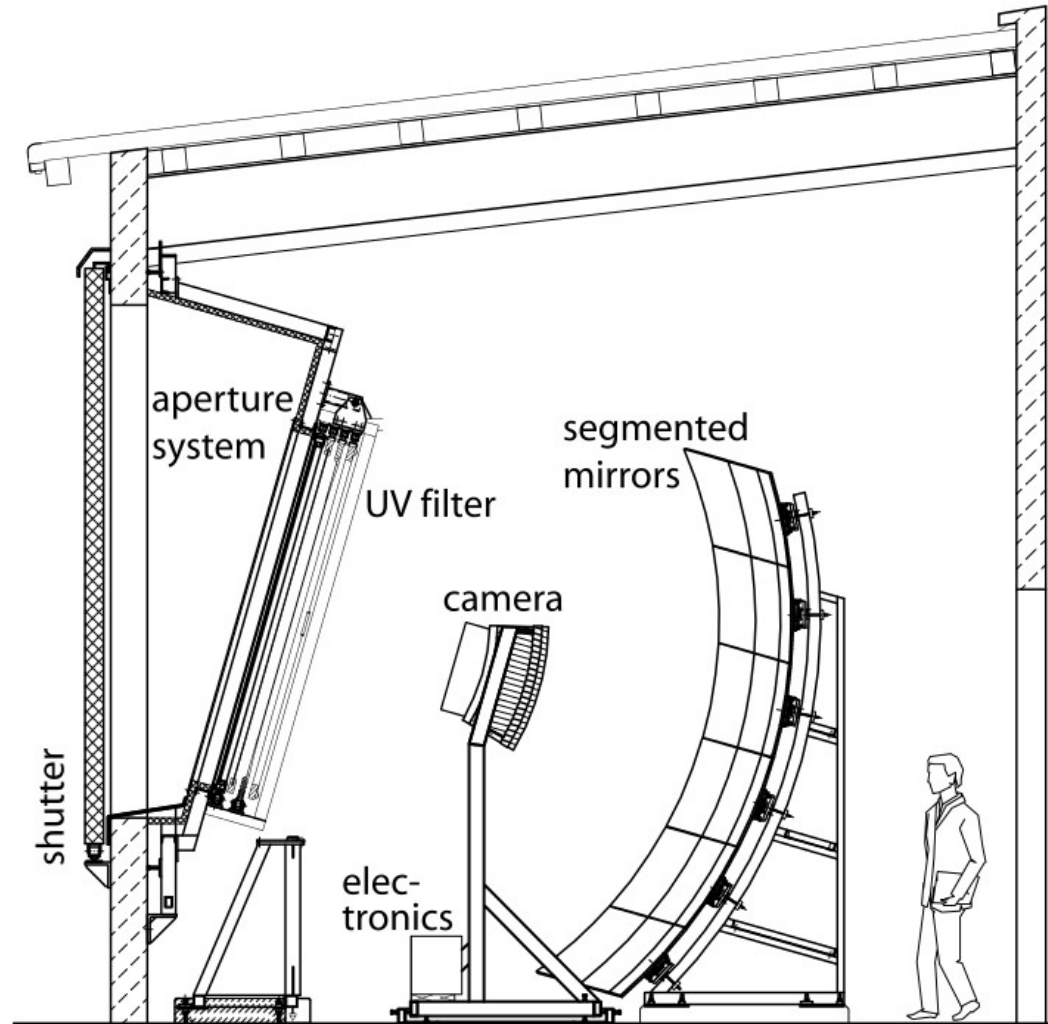
mirror 3 m²

440 pixel camera
10 MHz sampling

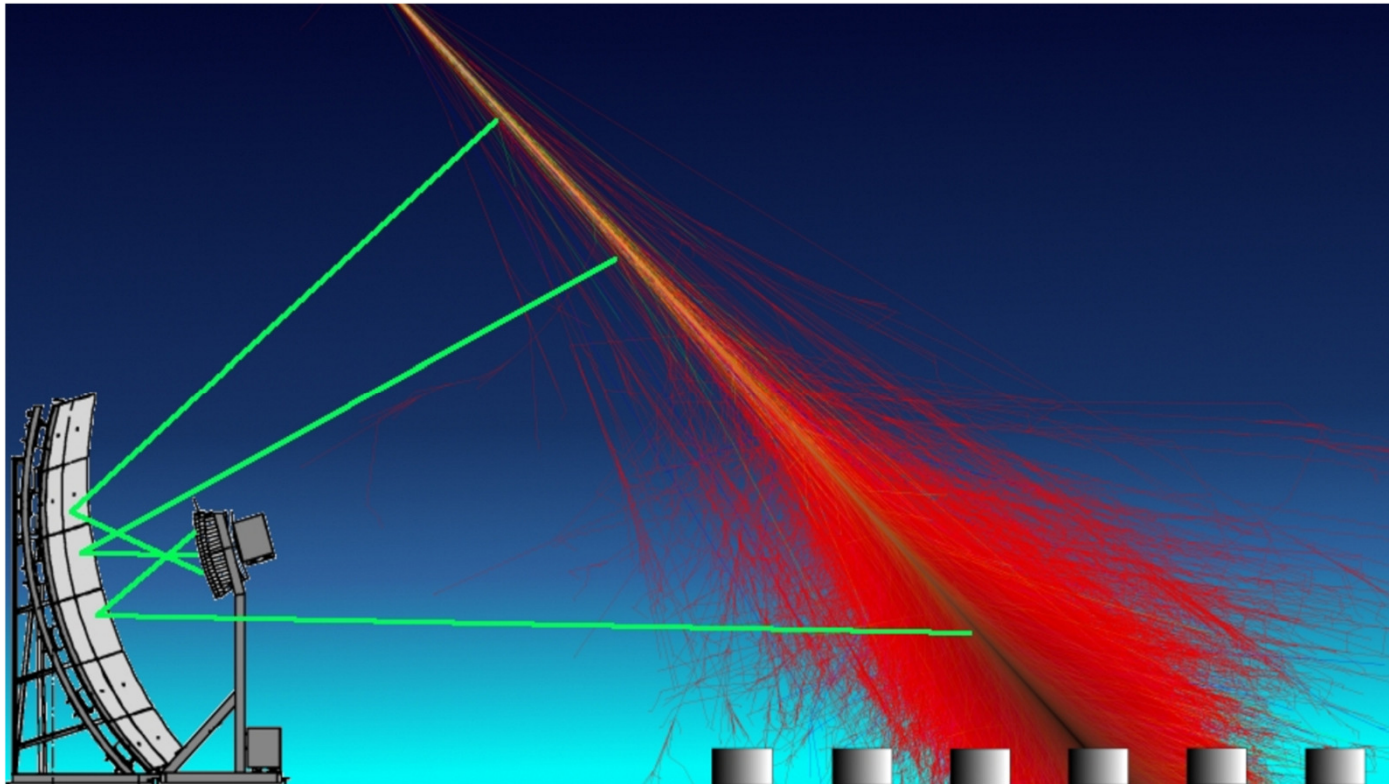


aperture, corrector ring and filter

A fluorescence telescope (FD)

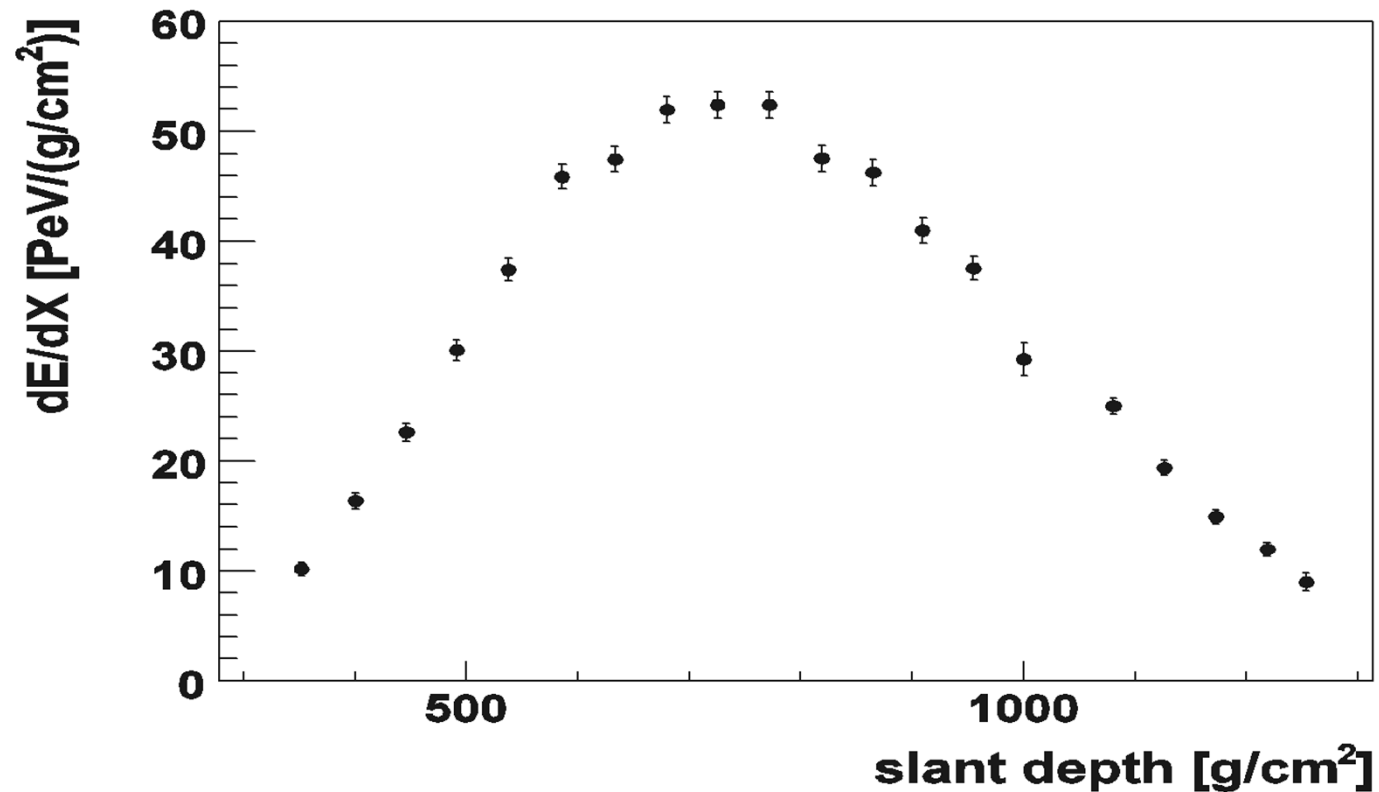


Telescope view



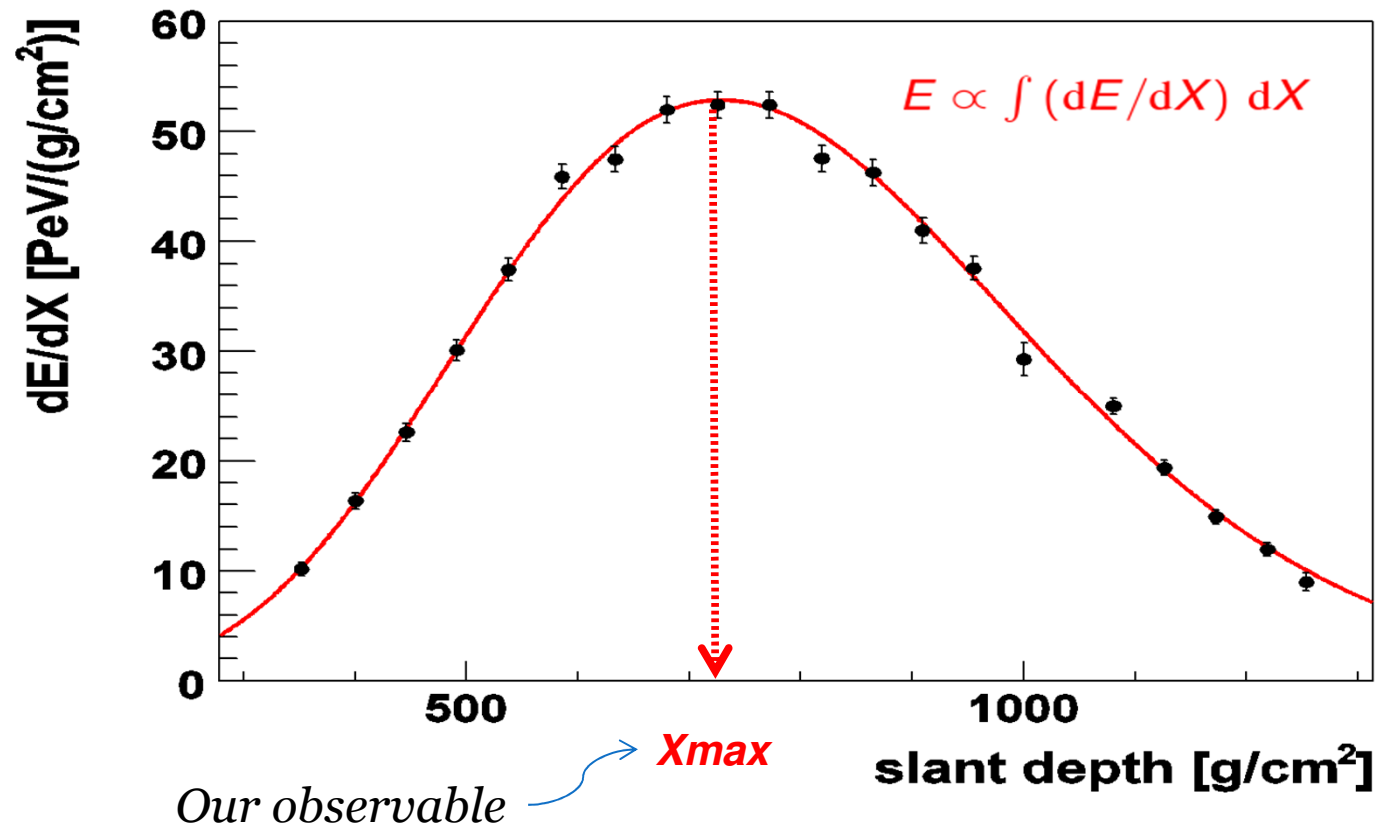
Telescopes measure the intensity and arrival time of the fluorescence light

Longitudinal Profile



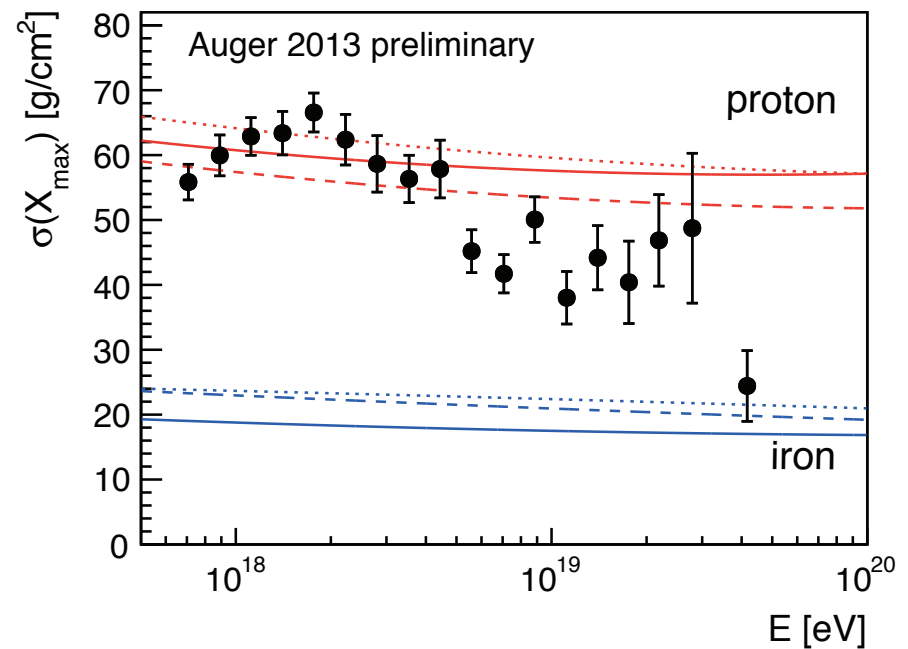
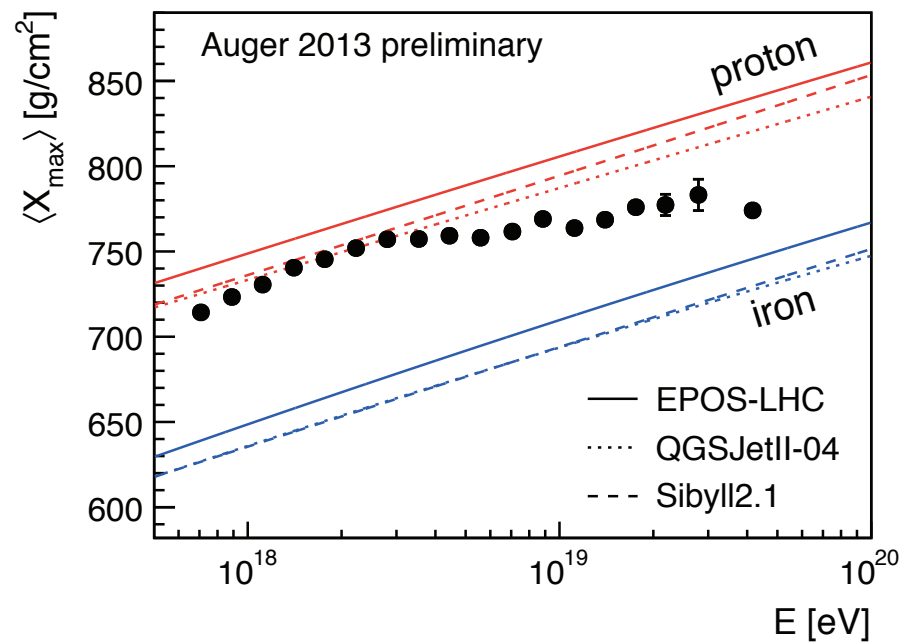
The intensity as a function of elevation can be transformed into the energy deposited in the atmosphere as a function of depth

Fitting the Longitudinal Profile

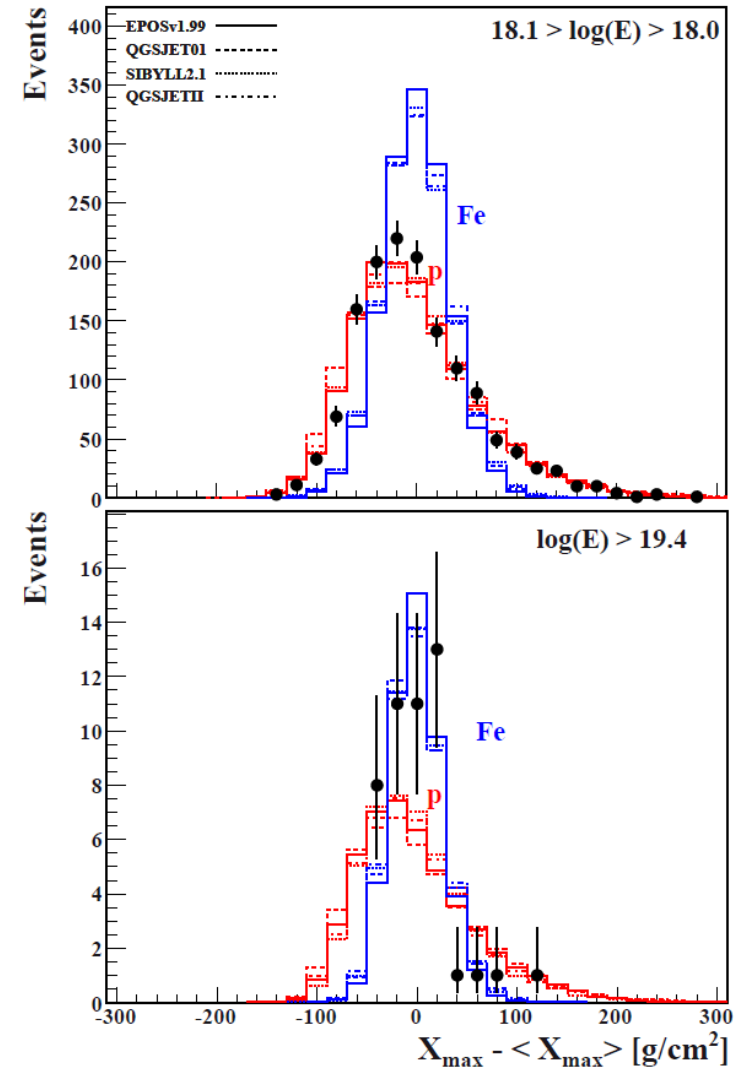
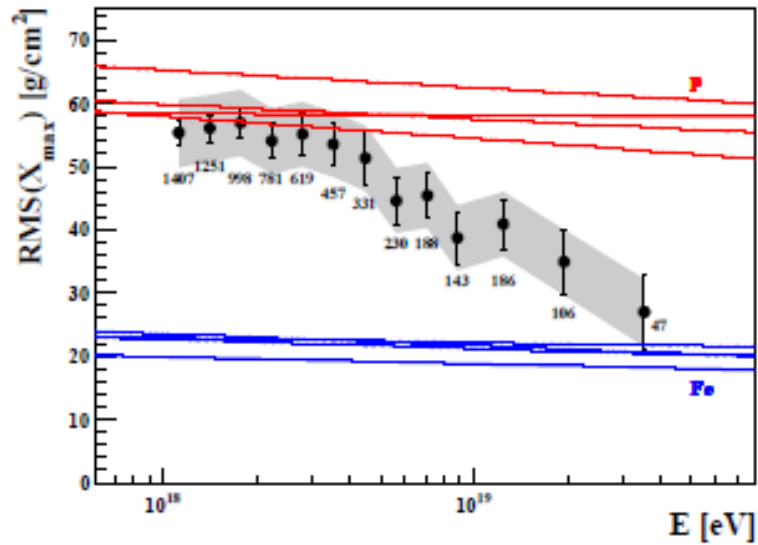
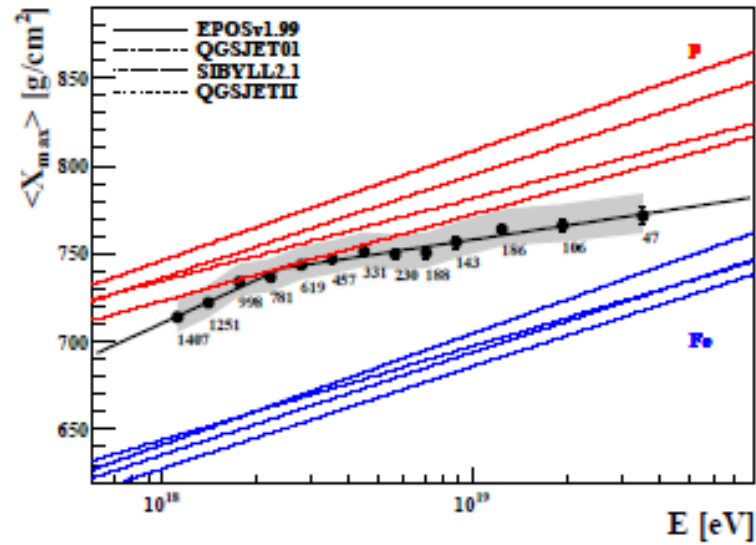


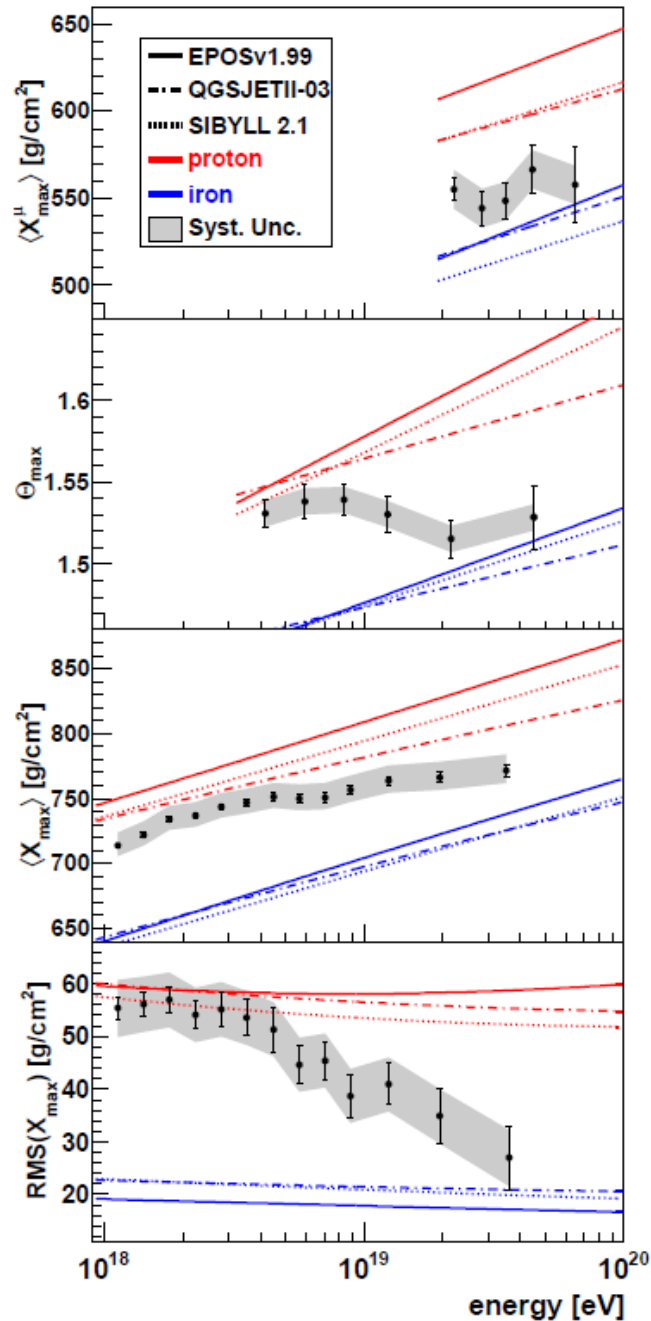
The total calorimetric energy of the shower is proportional to the integral of the energy deposited

Studying the primary mass composition



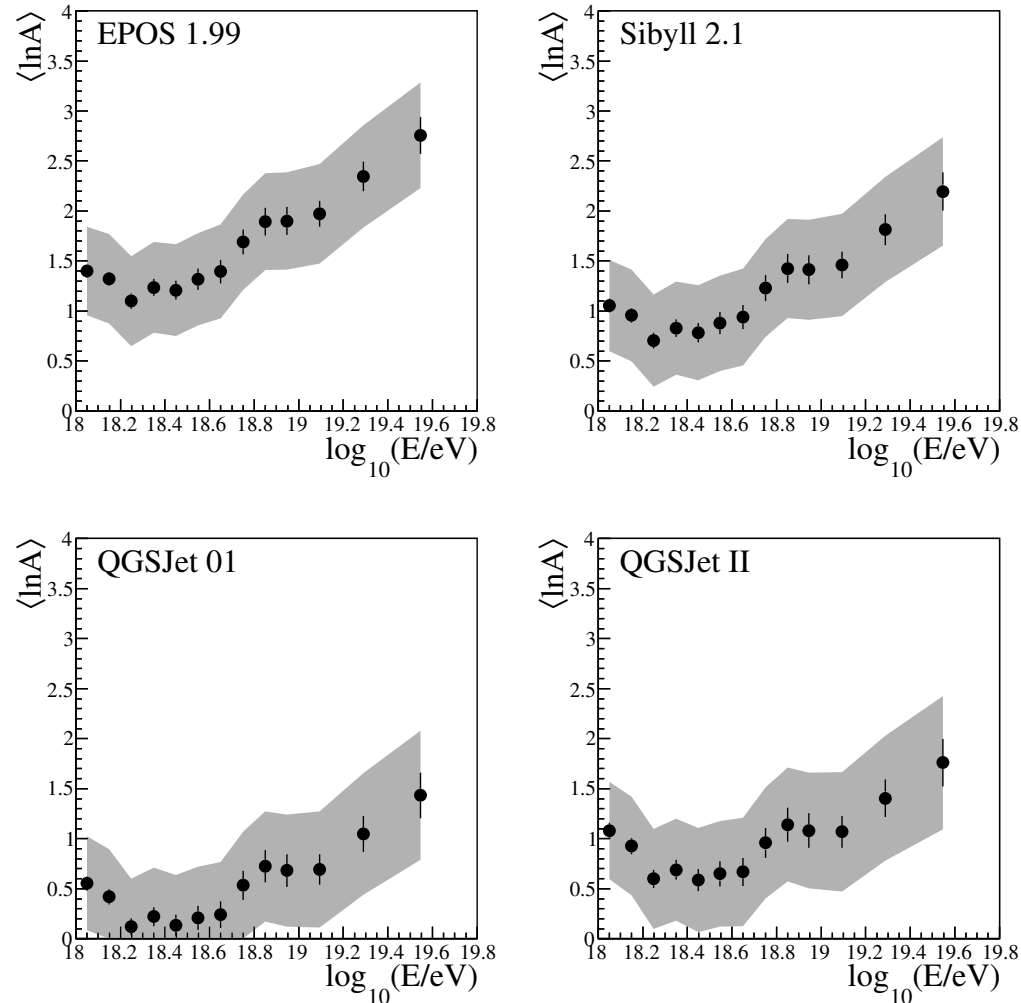
X_{max} and RMS (X_{max})





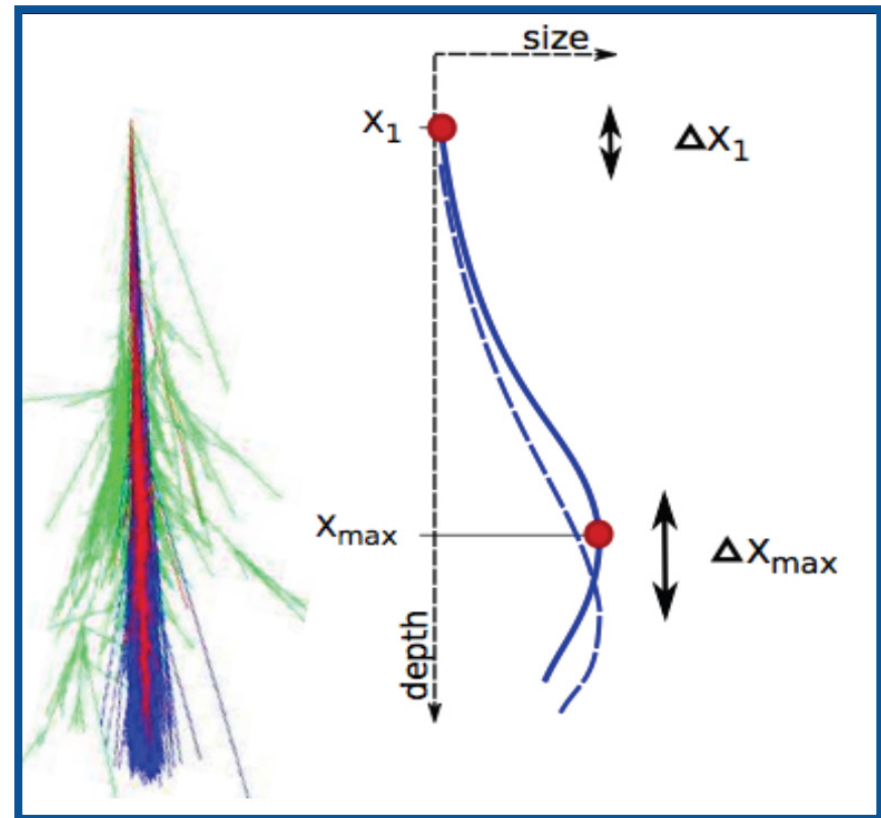
Results of the evolution of shower observables that are sensitive to X_{max} and comparison with predictions from various hadronic interaction models.

$\langle \ln A \rangle$ as function of $\log E$ from Auger data using different hadronic interaction models



Measuring cross section @ $\sqrt{s} = 57 \text{ TeV}$

- The path before the interaction, X_1 , is a function of the p -air cross section.
- The fluorescence telescopes measure the position of the maximum of the shower, X_{max} .
- Use MC models to related X_{max} to X_1 , and then σ_{p-air} .
- Using models, the value of σ_{inel} (p -air) is inferred, and then using a technique based on the Glauber method, σ_{inel} (pp) is evaluated.



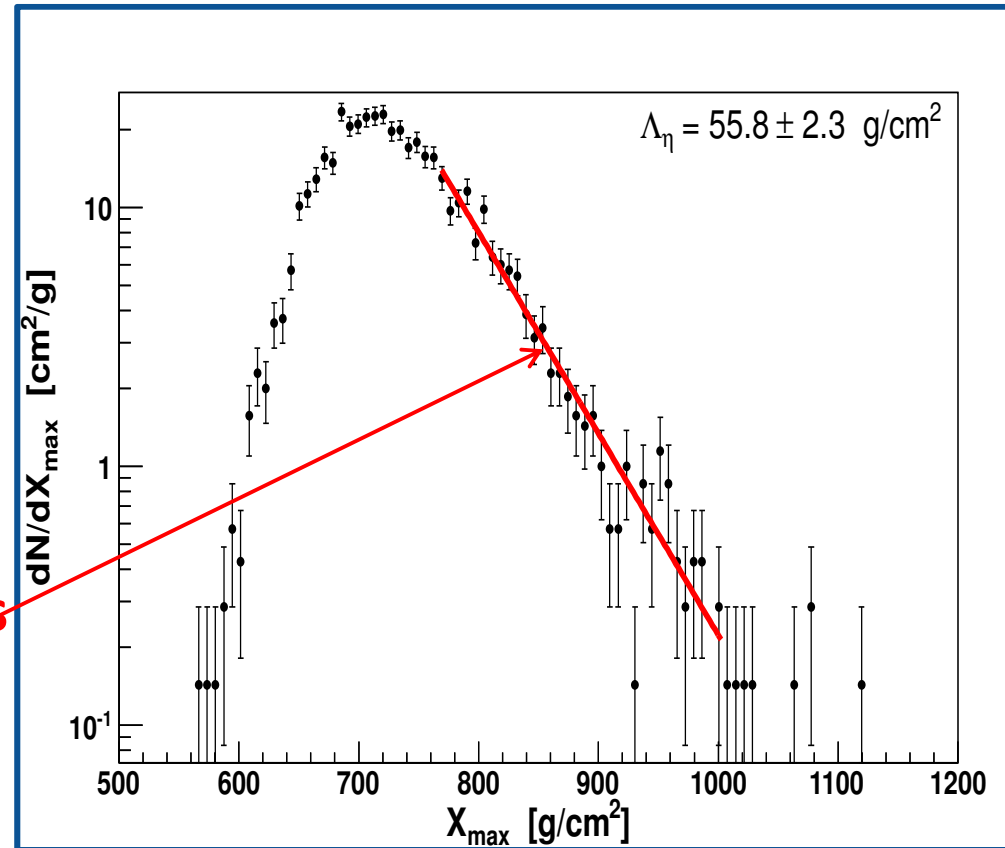
Selecting deep showers....

The position of the air shower maximum, at fixed energy, X_{max} , is sensitive to the cross section

$18 < \log_{10}(E/\text{eV}) < 18.5$
 \Rightarrow protons give a significant contribution to the overall flux

At these depths, protons showers dominate

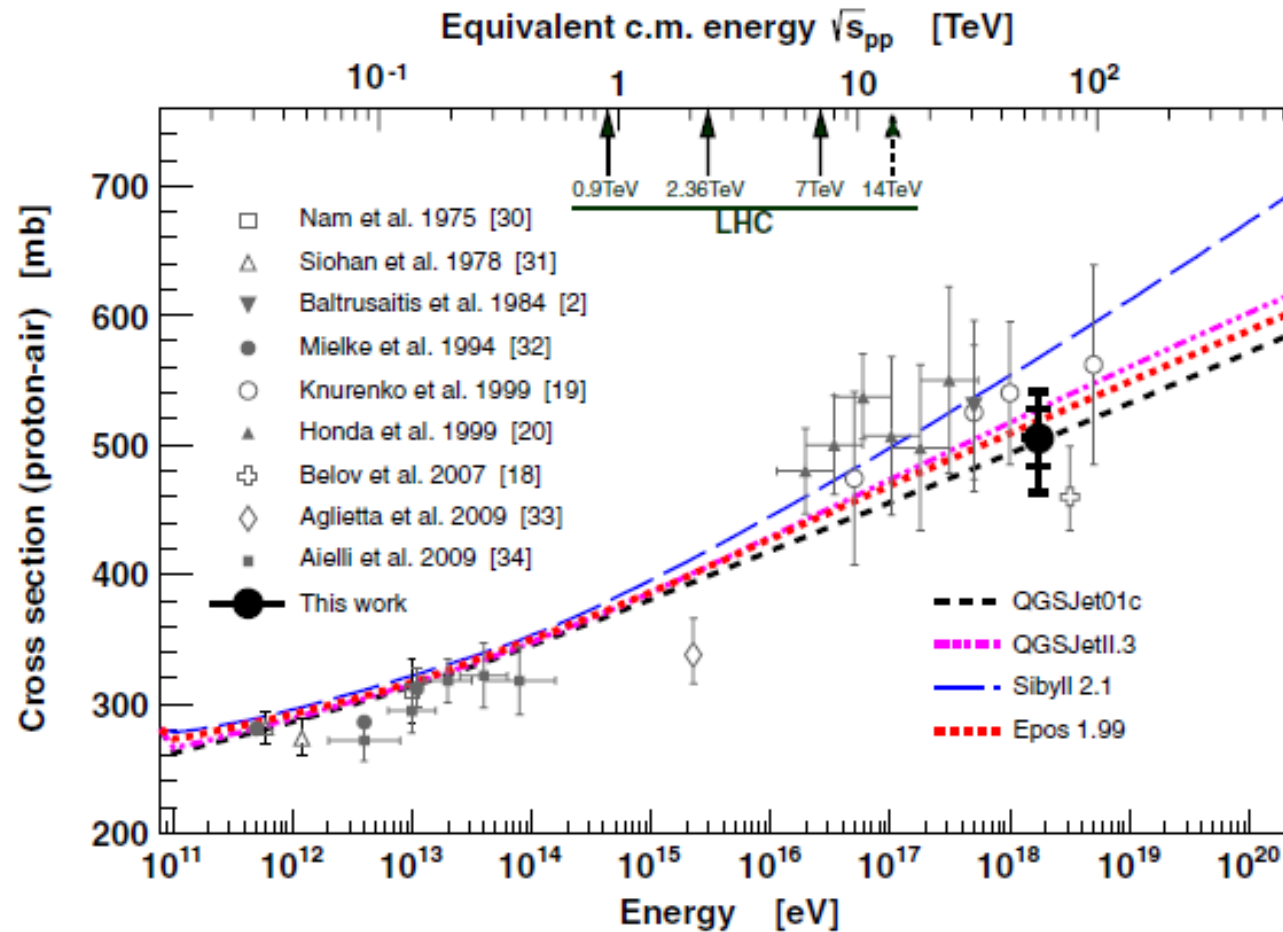
$$\frac{dN}{dX_{max}} = \exp\left(-\frac{X_{max}}{\Lambda_{\eta}}\right)$$



proton-air cross section for particle production

First:

$$\sigma_{p-air}^{prod} = \left(505 \pm 22 \Big|_{stat} \begin{array}{l} +28 \\ -36 \end{array} \Big|_{syst} \right) mb$$



The PA Collaboration, Phys. Rev. Lett. 109, 062002 (2012)

TABLE I. Summary of the systematic uncertainties.

Description	Impact on $\sigma_{p\text{-air}}^{\text{prod}}$
Λ_η systematics	± 15 mb
Hadronic interaction models	$-8 + 19$ mb
Energy scale	± 7 mb
Conversion of Λ_η to $\sigma_{p\text{-air}}^{\text{prod}}$	± 7 mb
Photons, $< 0.5\%$	$< + 10$ mb
Helium, 10%	-12 mb
Helium, 25%	-30 mb
Helium, 50%	-80 mb
Total (25% helium)	-36 mb, $+28$ mb

The basic concepts...

$$\Gamma_{hA}(\vec{b}) = 1 - \prod_{j=1}^A \left[1 - \int \Gamma_{hN}(\vec{b} - \vec{s}_j) \rho_j(\vec{r}_j) d^3 r_j \right]$$

$$\sigma_{hA}^{tot} = \frac{4\pi}{|\vec{k}|} \Im m \left\{ f_{\ddot{u}}^{hA}(S, \vec{q}_2 \rightarrow 0) \right\} = 2\Re e \int \Gamma_{hA}(\vec{b}) d^2 b$$

$$\sigma_{hA}^{ela} = \int \frac{1}{|\vec{k}|^2} \left| f_{\ddot{u}}^{hA}(S, \vec{q}_2) \right|^2 d^2 q = \int |\Gamma_{hA}(\vec{b})|^2 d^2 b$$

$$\Gamma_{pp \rightarrow pX}(s, \vec{b}) = \lambda(s) \Gamma_{pp \rightarrow pp}(s, \vec{b})$$

Elastic scattering and inelastic diffraction

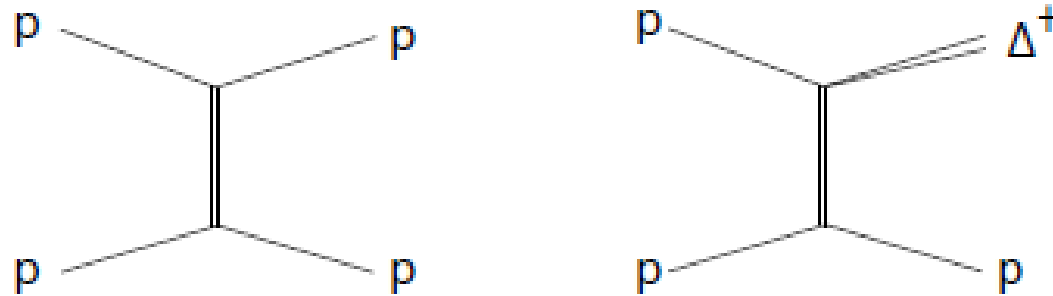
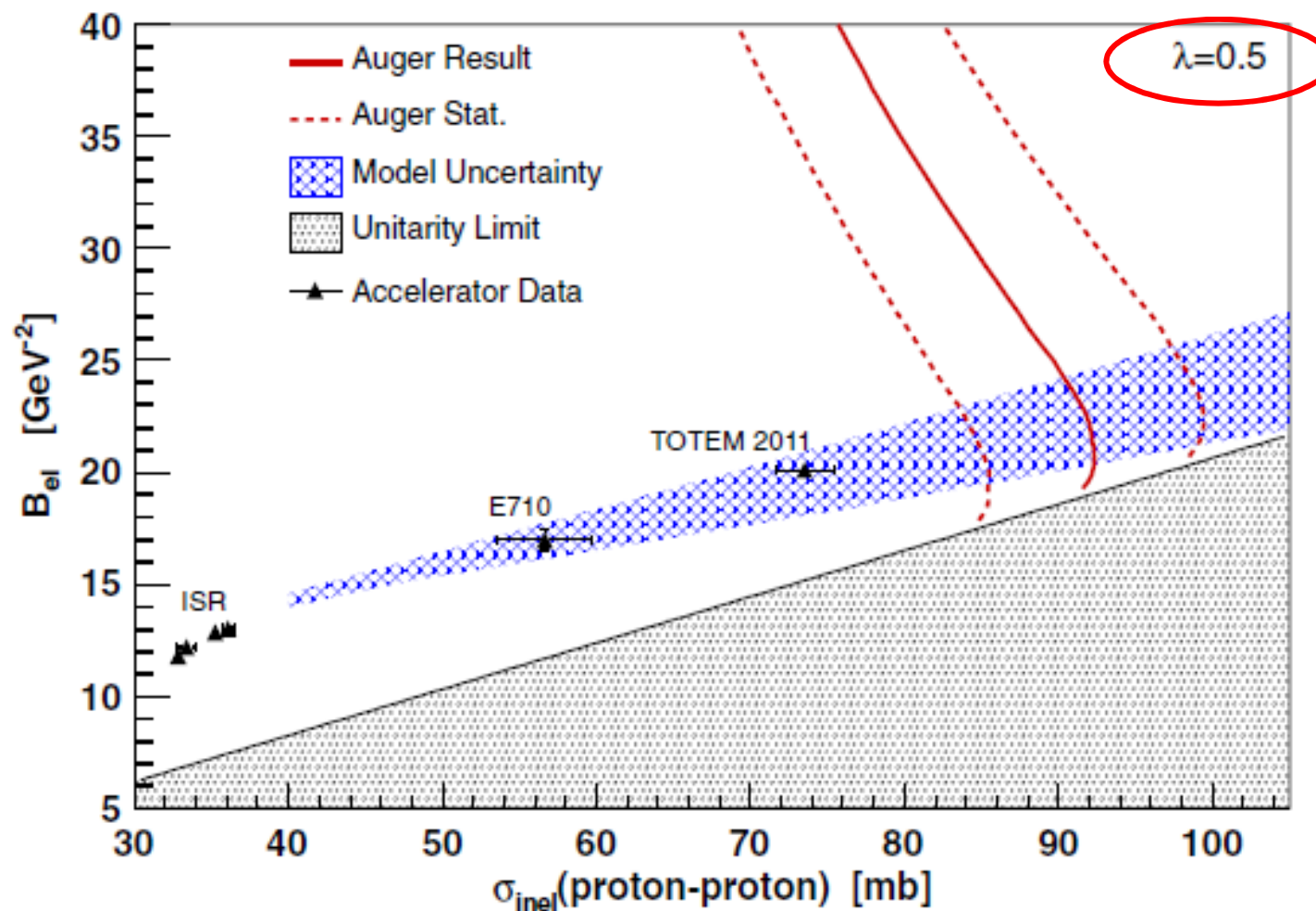


Figure 1: Left panel: Elastic scattering, Right panel: inelastic diffraction.

Step by step....

- Parametrization of the p-p and p-n cross sections and of the ratio real/imaginary part of the scattering amplitude
- Parametrization of the energy dependence of the data on single diffraction dissociation
- Determination of the ratio of the cross sections for diffraction dissociation and elastic scattering in dependence on the maxima mass of the diffractive system as parameter $M_{D,max}$
- Calibration of the parameter $M_{D,max}$ with proton-carbon data
- Application of this value for the mass limit in the Glauber calculation to convert the measured proton-air cross section to the pp inelastic cross section

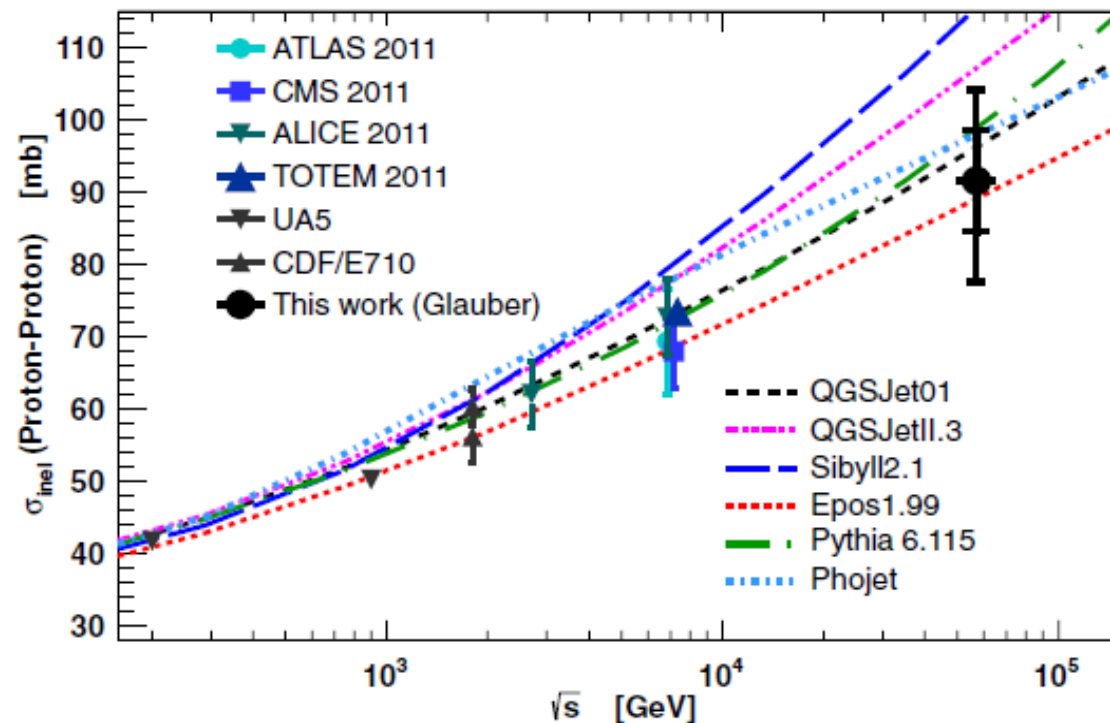
Correlation of B_{el} and $\sigma^{inel}(pp)$ in Glauber



proton-proton cross section

$$\sigma_{pp}^{inel} = \left(92 \pm 7 \Big|_{stat} \begin{array}{c} +9 \\ -11 \end{array} \Big|_{syst} \pm 7 \Big|_{Glauber} \right) mb$$

$$\sigma_{pp}^{tot} = \left(133 \pm 13 \Big|_{stat} \begin{array}{c} +17 \\ -20 \end{array} \Big|_{syst} \pm 16 \Big|_{Glauber} \right) mb$$



σ_{pp} is inferred using an extended Glauber approach: calculation with parameters from accelerator measurements @ lower energies that are extrapolated to CR energies

Final results:

$$\sigma_{pp}^{inel} = [92 \pm 7 (stat) + 9 / -11 (syst) \pm 7 (Glauber)] \text{ mb}$$

$$\sigma_{pp}^{tot} = [133 \pm 13 (stat) + 17 / -20 (syst) \pm 16 (Glauber)] \text{ mb}$$

ACKNOWLEDGEMENTS

