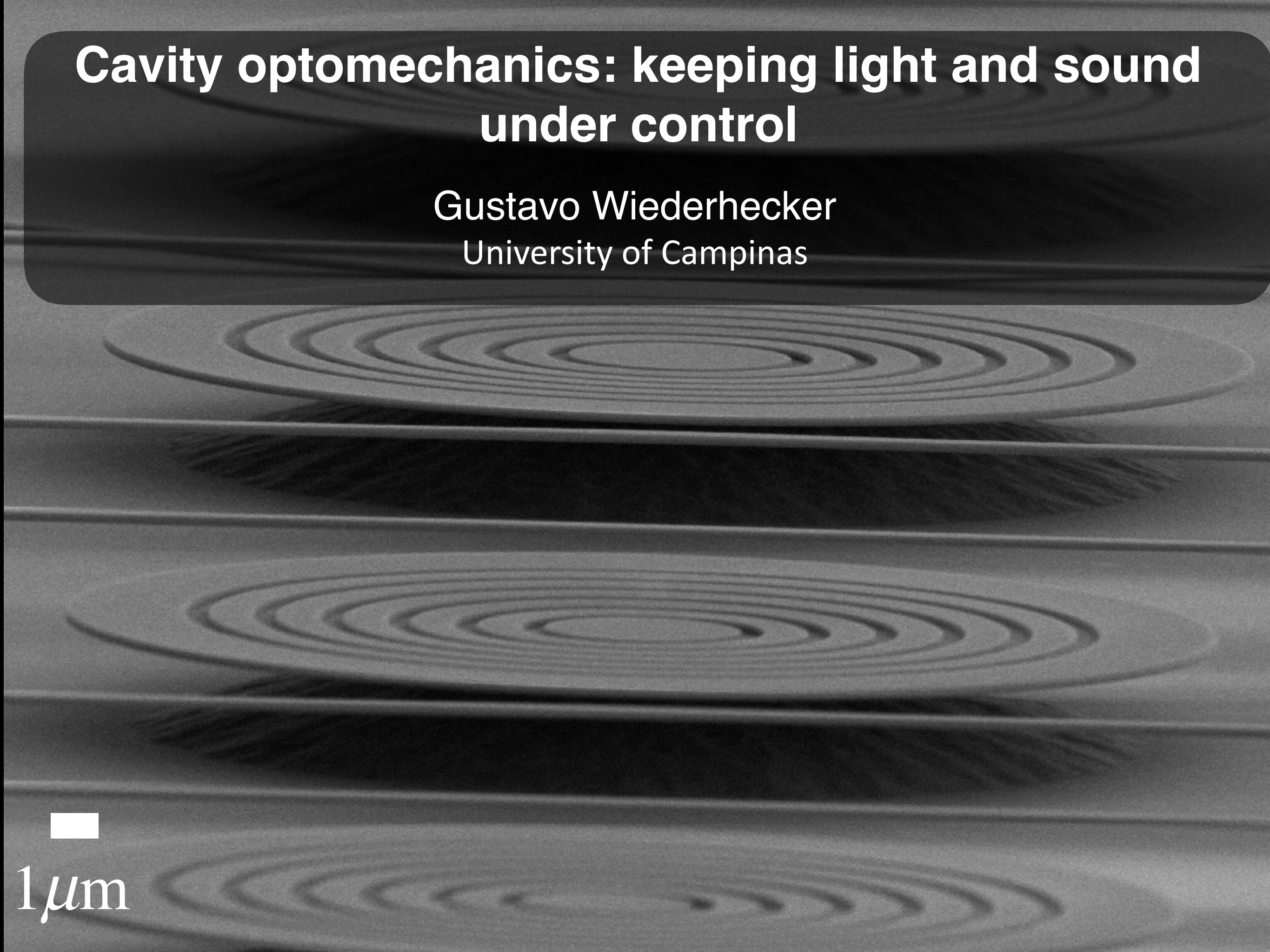


Cavity optomechanics: keeping light and sound under control

Gustavo Wiederhecker
University of Campinas



1 μm

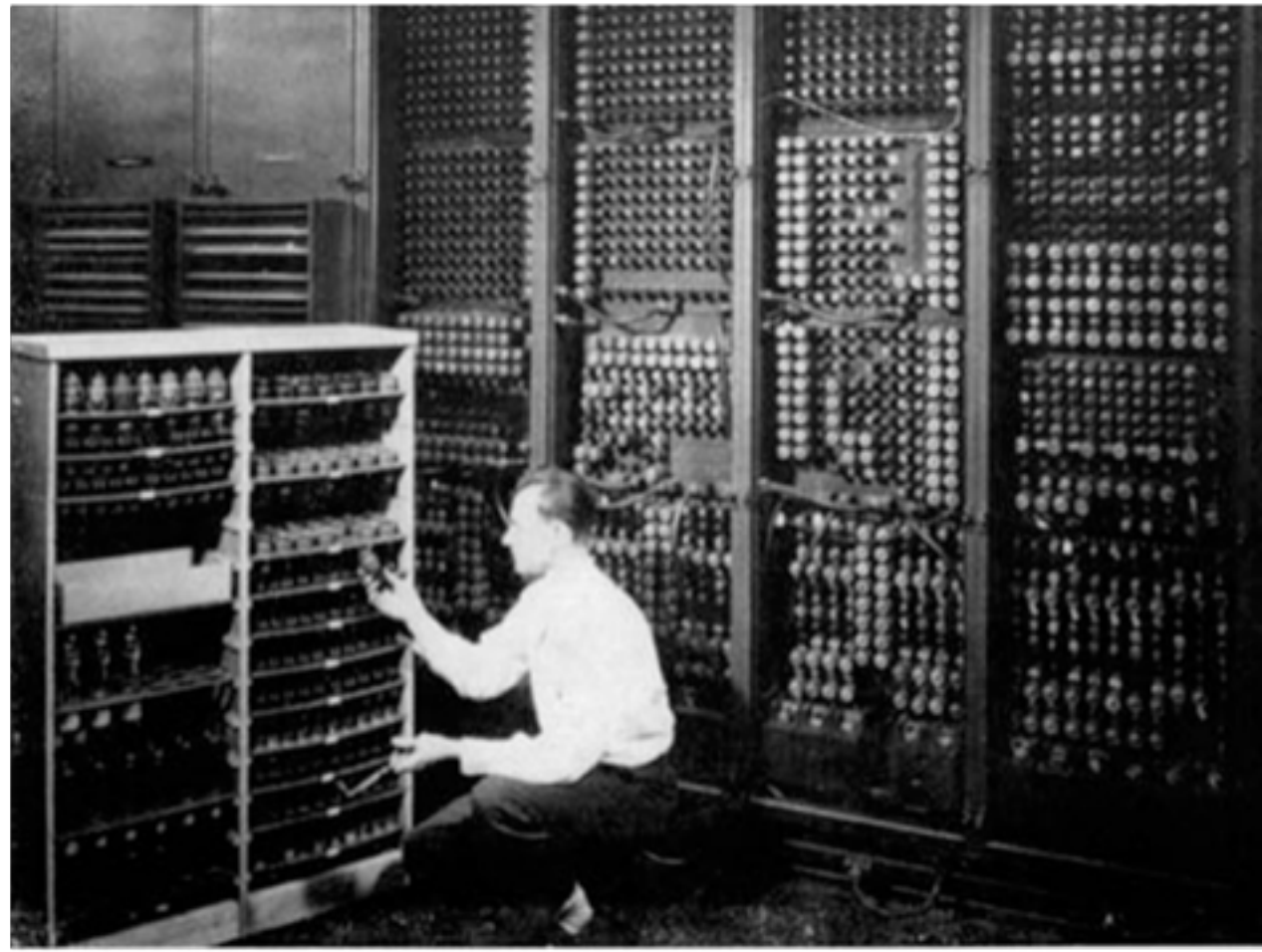


Silicon Nanophotonics



Z2 (Germany, 1939)

- **1.2 flops**
- 300 kg
- 1 kW

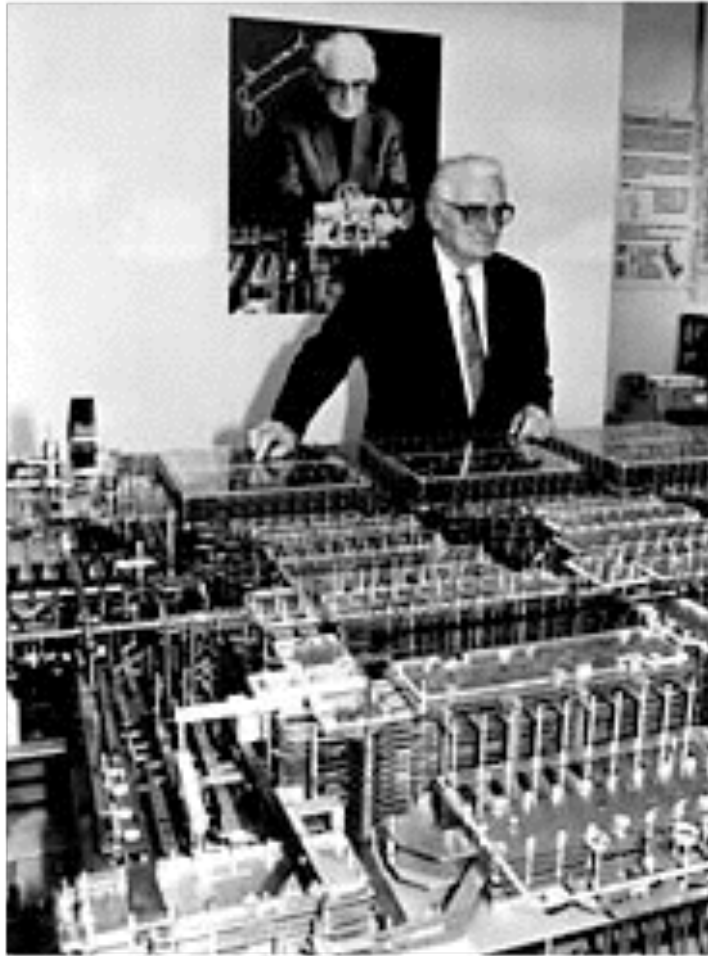


ENIAC (USA, 1946)

- **300 flops**
- 27,000 kg
- 150 kW

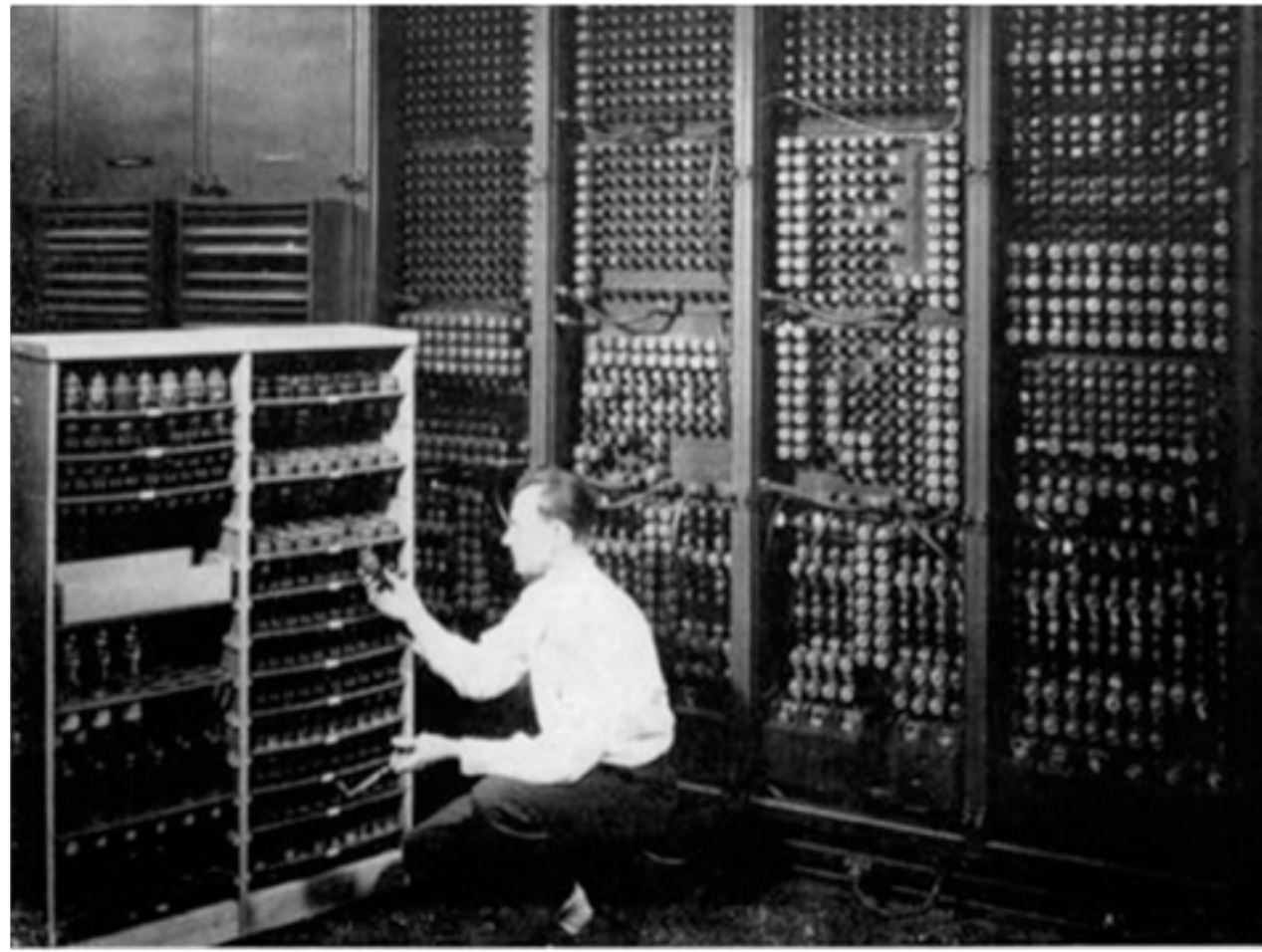


Silicon Nanophotonics



Z2 (Germany, 1939)

- **1.2 flops**
- 300 kg
- 1 kW



ENIAC (USA, 1946)

- **300 flops**
- 27,000 kg
- 150 kW



Iphone 6 (2014)

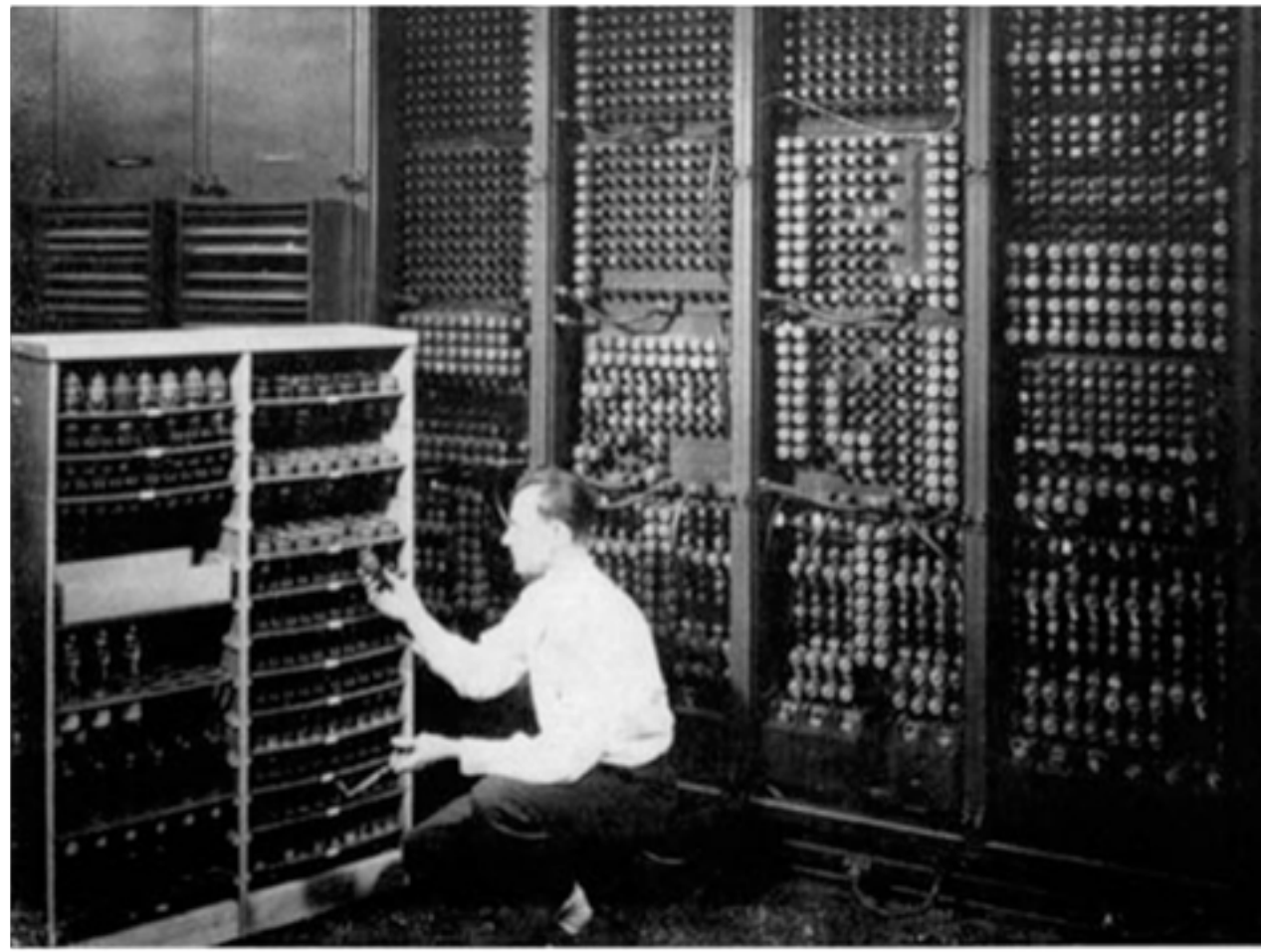


Silicon Nanophotonics



Z2 (Germany, 1939)

- **1.2 flops**
- 300 kg
- 1 kW



ENIAC (USA, 1946)

- **300 flops**
- 27,000 kg
- 150 kW



Iphone 6 (2014)

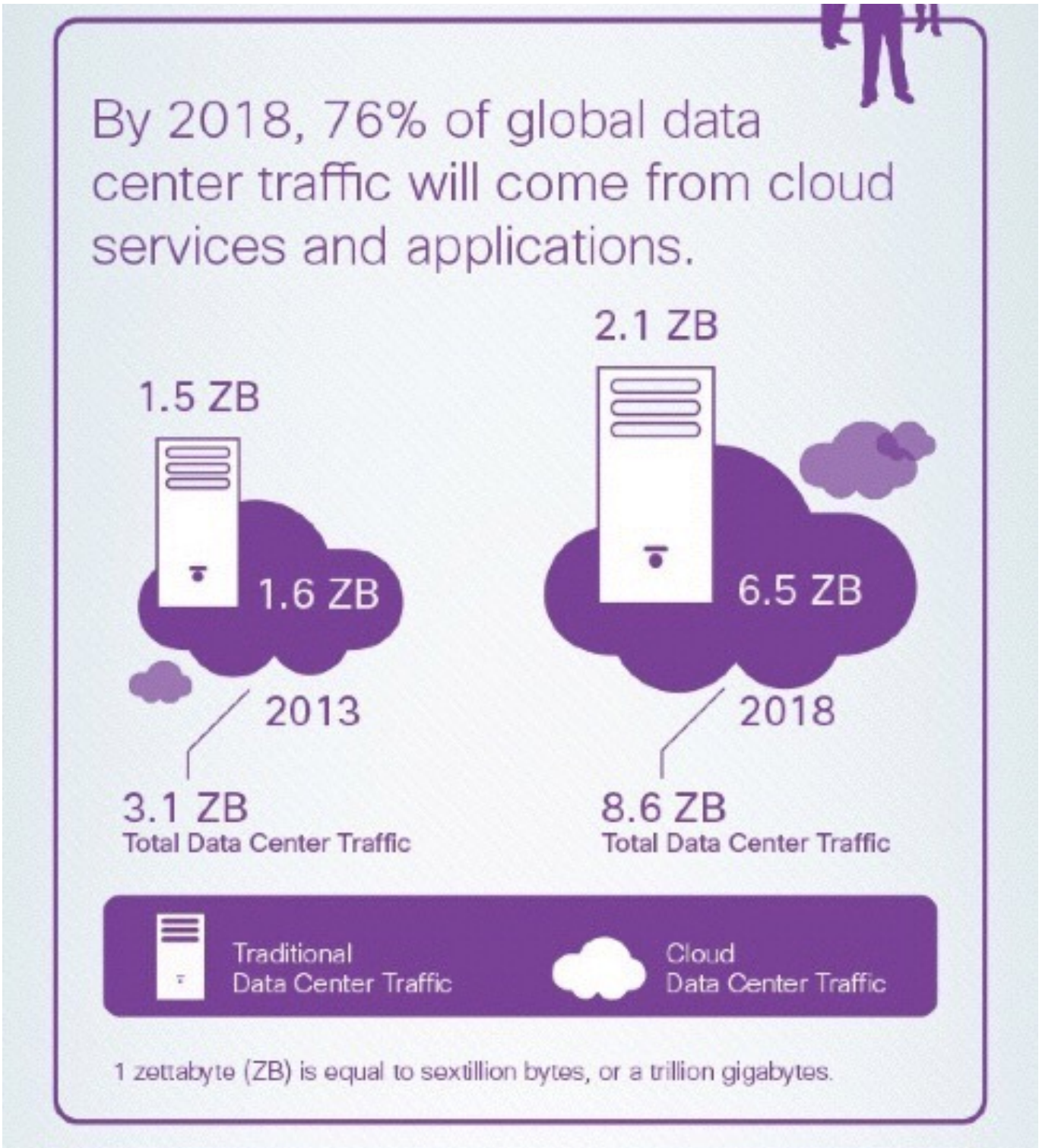
- **150 Gflops**
- 130 g
- battery

Silicon Nanophotonics



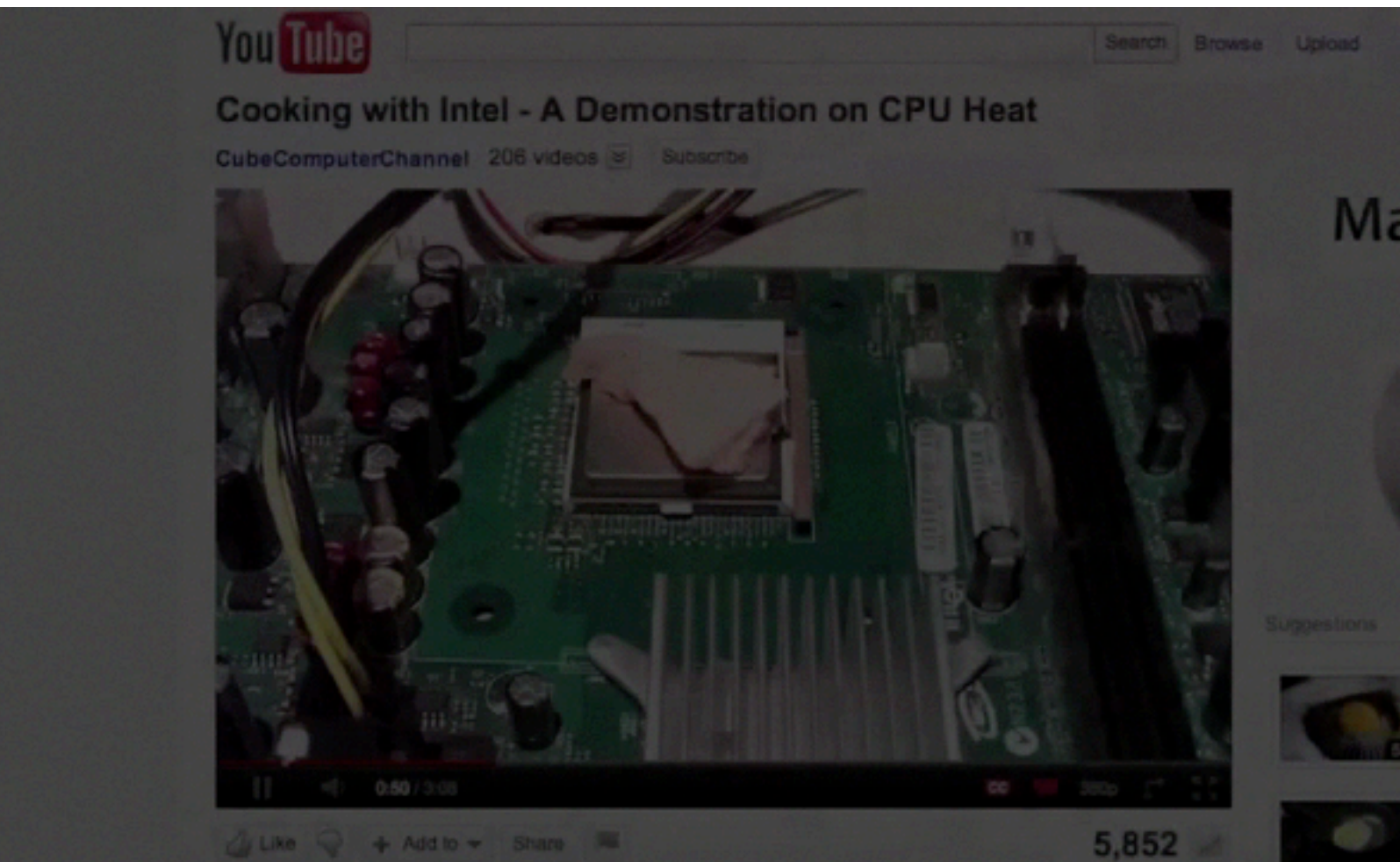


Datacenter bottleneck

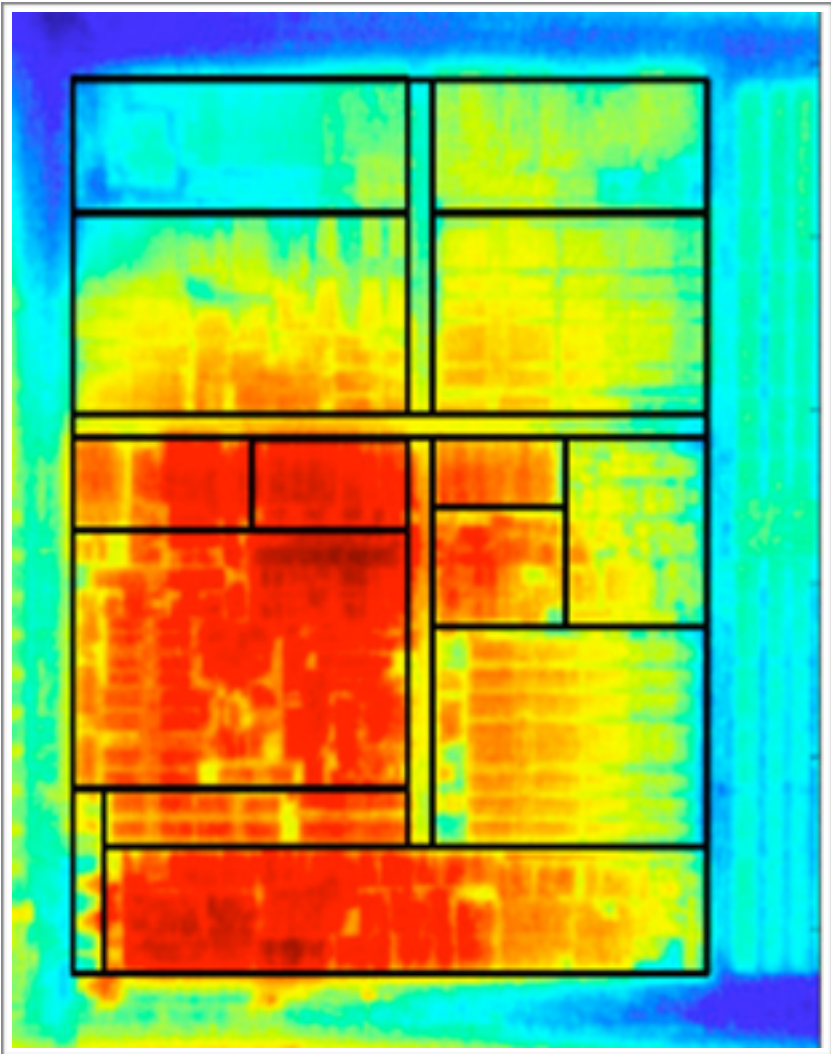




Temperature Issues



Bacon “on-chip”

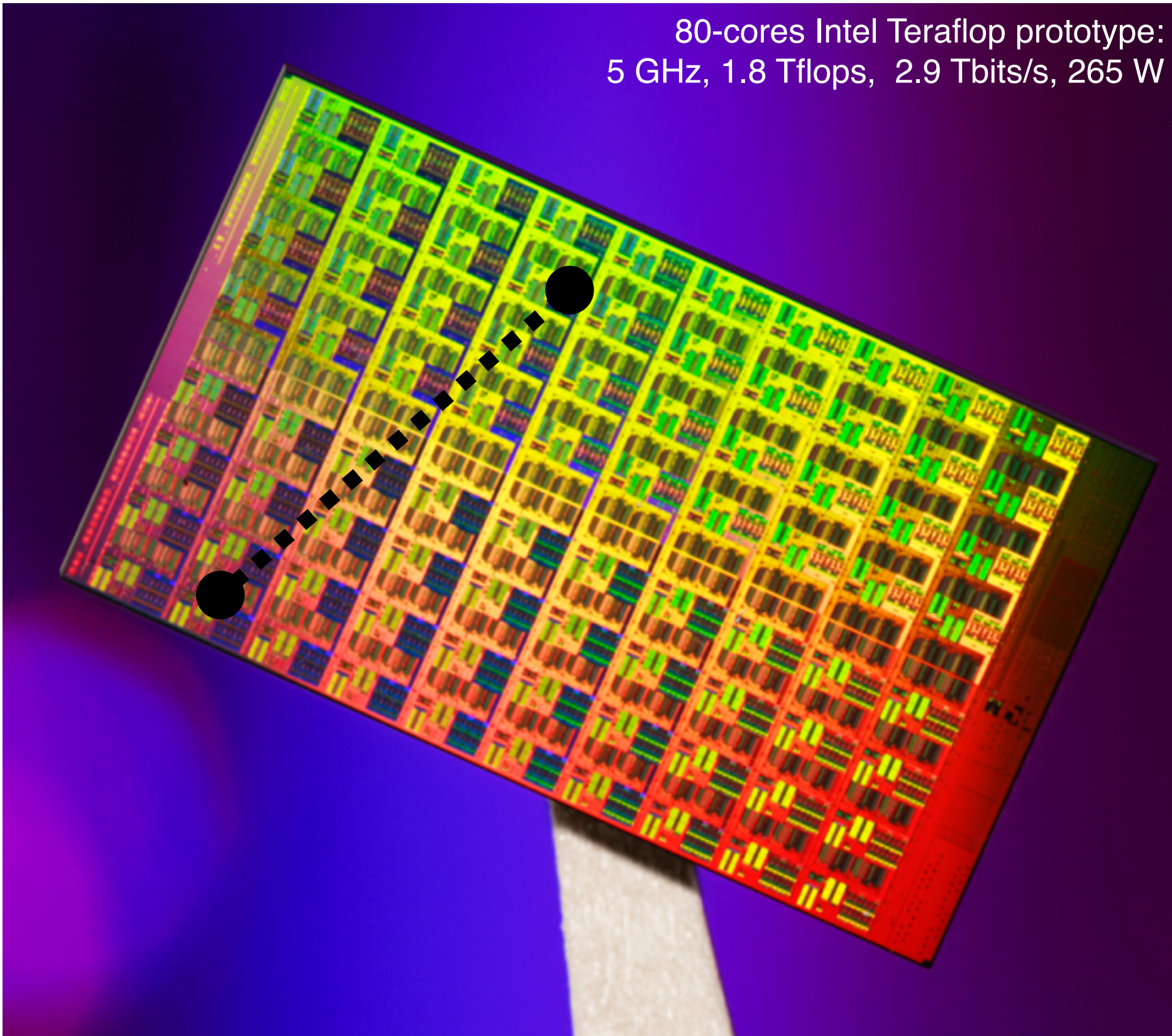


Temperature map
 $\Delta T \approx 80\text{ K}$



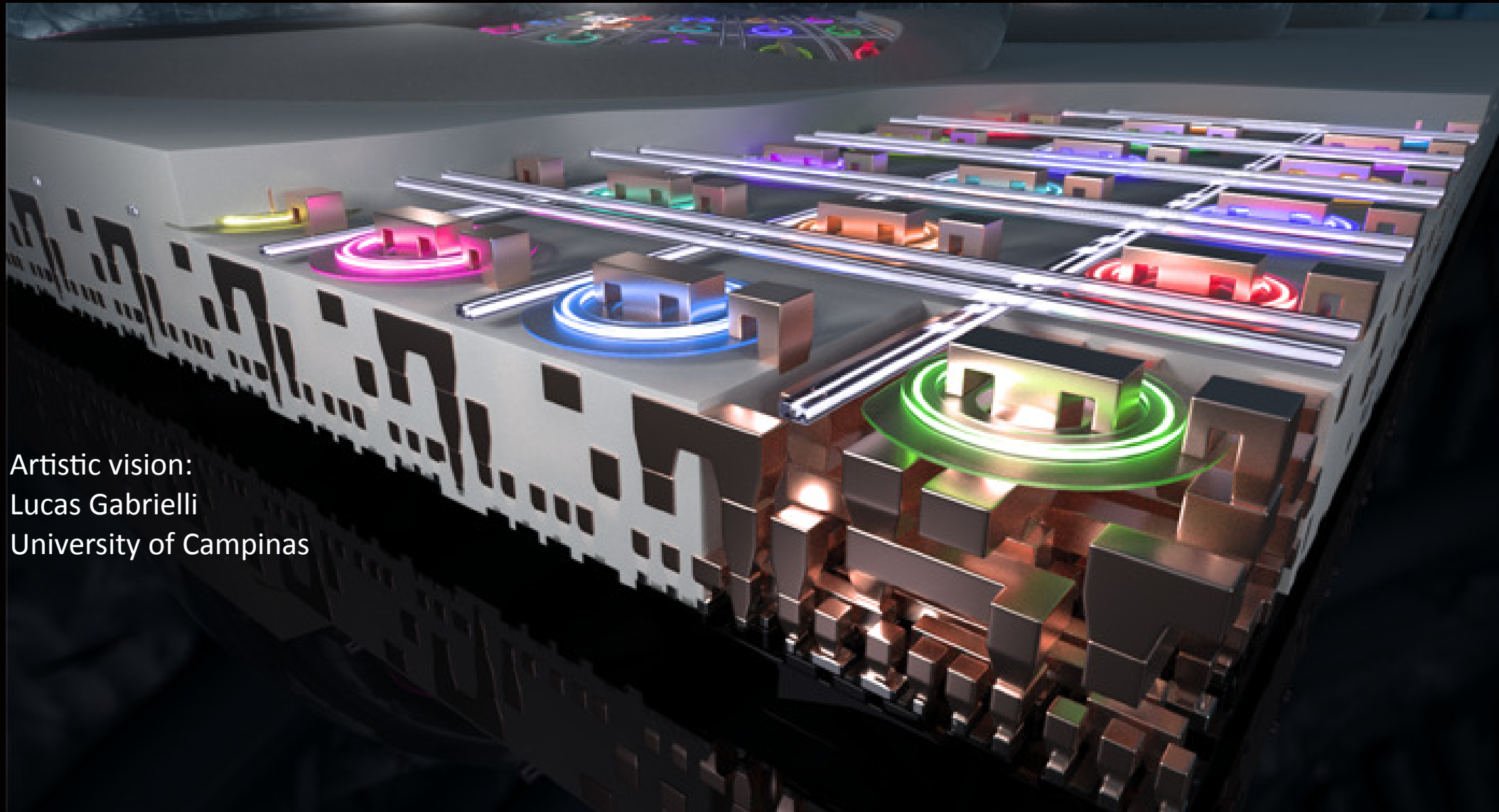
The multi-core solution

80-cores Intel Teraflop prototype:
5 GHz, 1.8 Tflops, 2.9 Tbits/s, 265 W





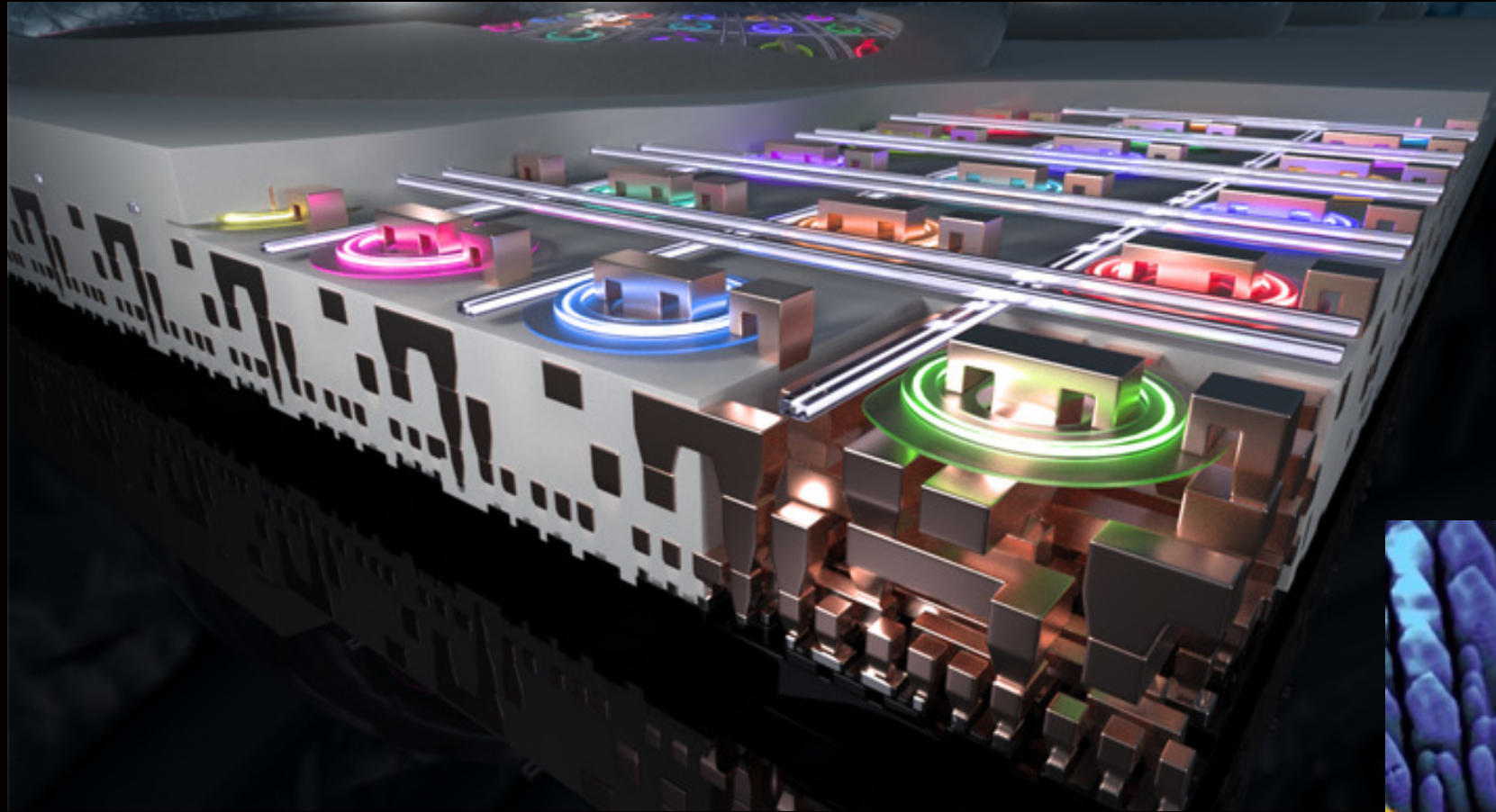
Si Nanophotonics



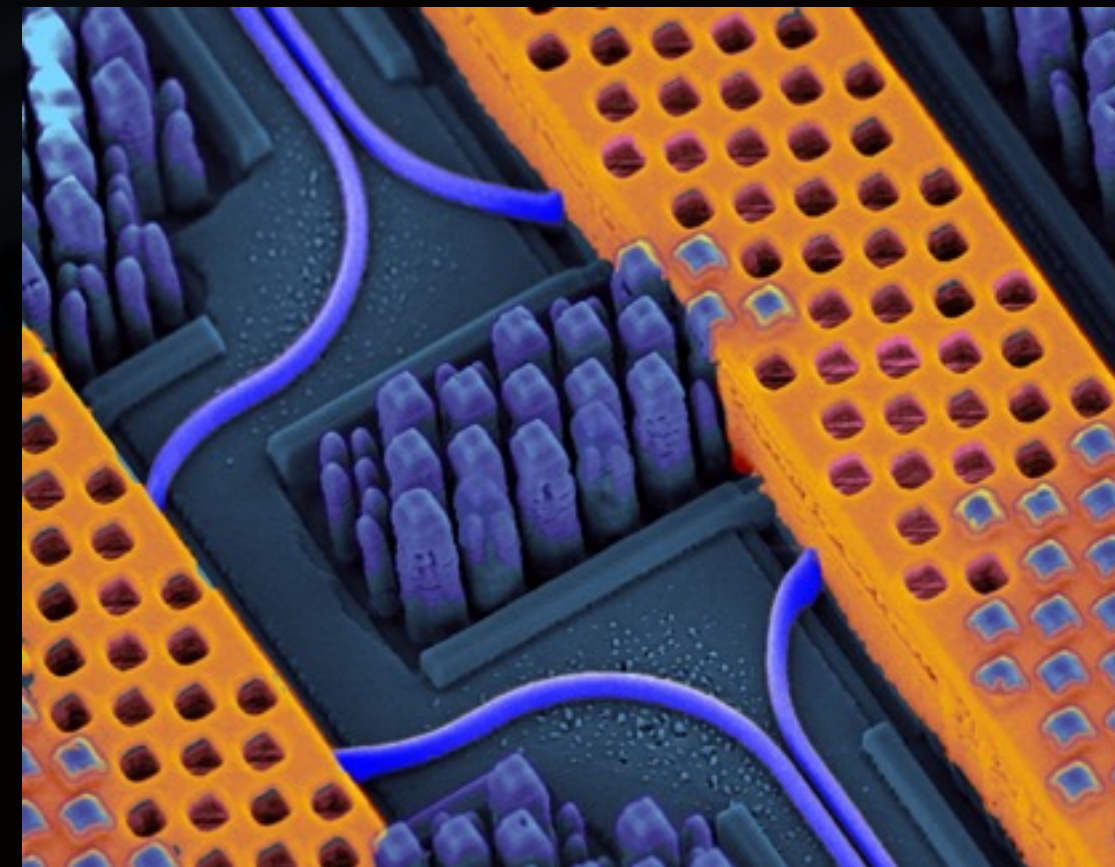
Artistic vision:
Lucas Gabrielli
University of Campinas



Si Nanophotonics



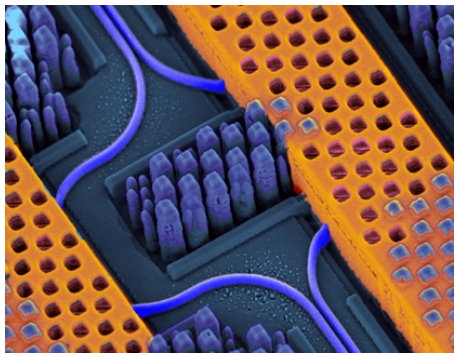
**What kind of
nanophotonic building
blocks we need here?**



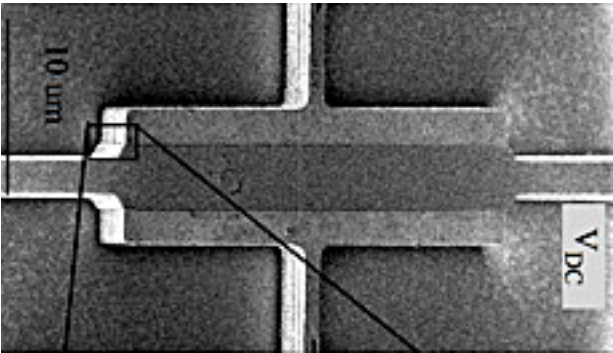
IBM 90-nm BEOL integration (2013)



Technological viewpoint: Si Nanophotonics Building Blocks

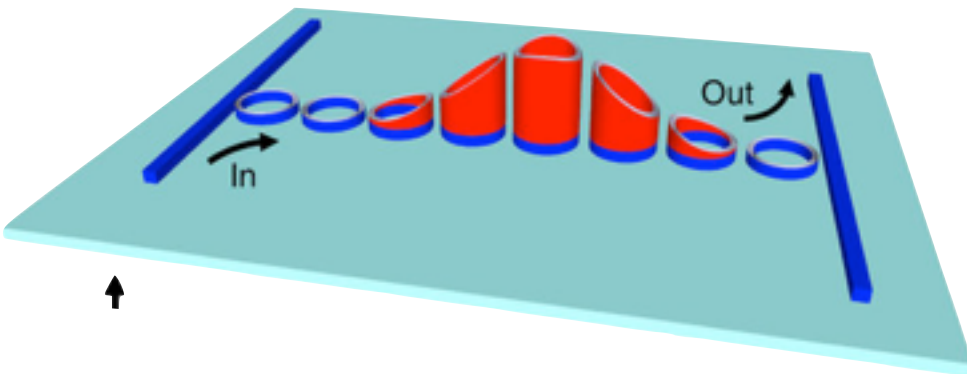


Silicon MEMS oscillators



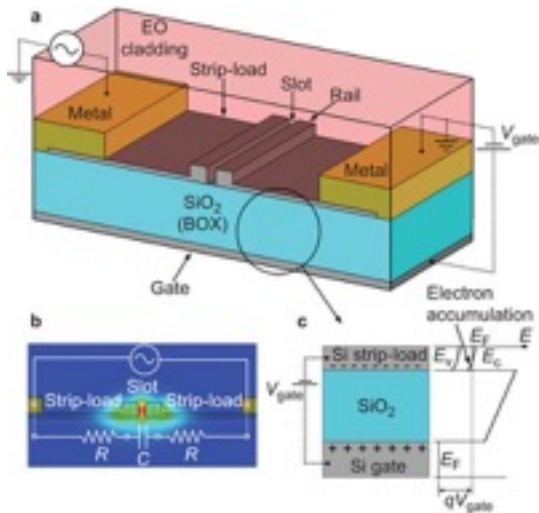
Bhave's group (Cornell)
Journal of MEMS (2009)

Wavelength Conversion



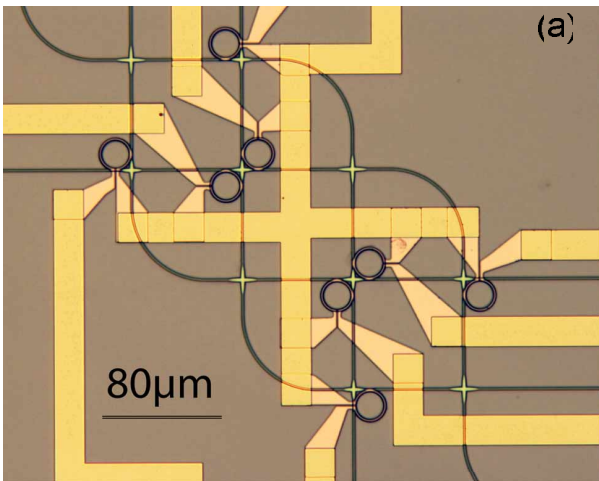
Melloni Group (Milano)
Nature Comm. 2011

Modulators



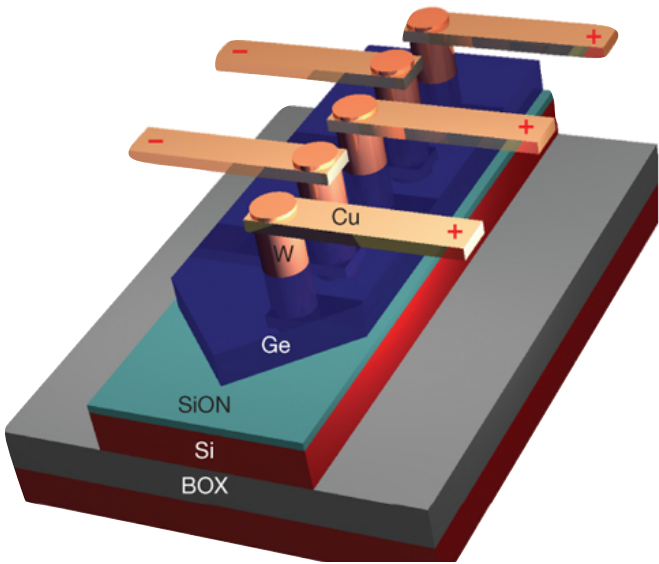
Leuthold group (KIT)
Nature 2005

Routers



Lipson's group (Cornell)
Opt. Express 2008

Photodetectors



Vlasov's group (IBM)
Nature 2010



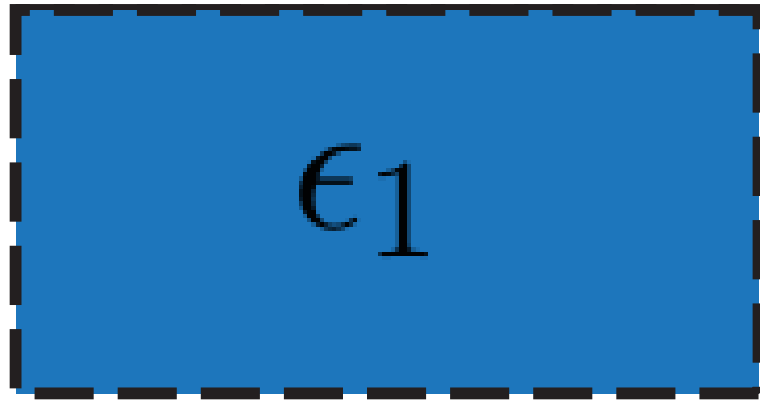
Outline

- ★ **Optical and acoustic mode interaction**
- ★ Optical force actuation
- ★ Dynamical back-action
- ★ Optomechanical clocks
- ★ Bullseye - a case study
- ★ Outlook



Light-sound interaction

ϵ_2



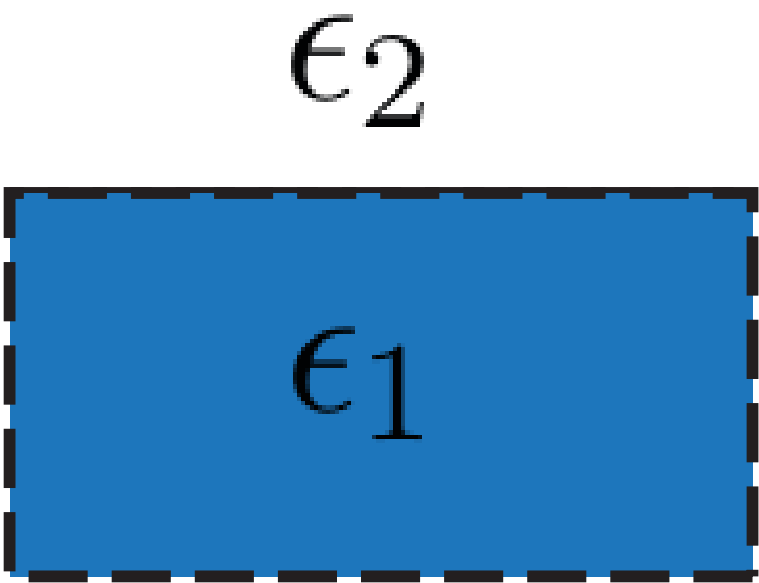
Dielectric
waveguide

E. P. Ippen and R. H. Stolen, “Stimulated Brillouin scattering in optical fibers,” Appl. Phys. Lett., vol. 21, pp. 539–541, Dec. 1972.

R. H. Stolen - “The Early Years of Fiber Nonlinear Optics”. JLT, VOL. 26, NO. 9, MAY 1, 2008



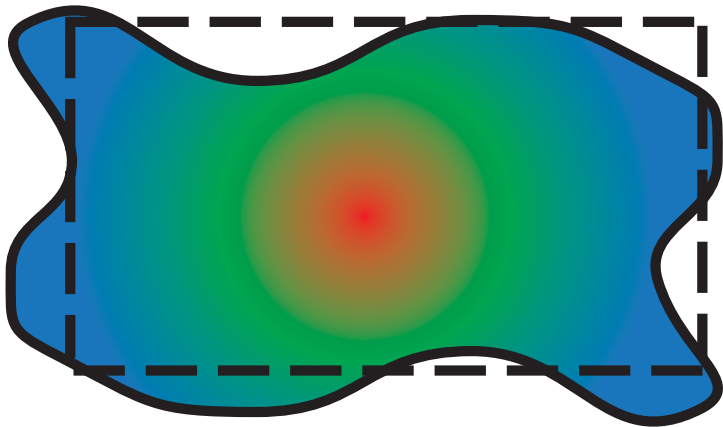
Light-sound interaction



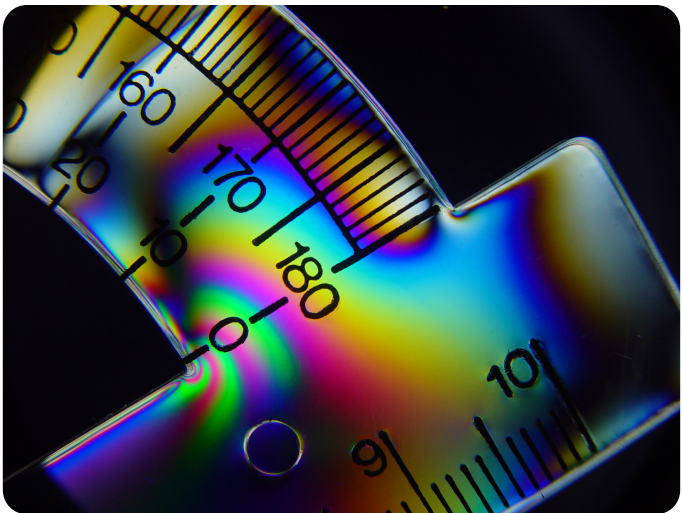
Dielectric waveguide



Mechanical mode
Deformation



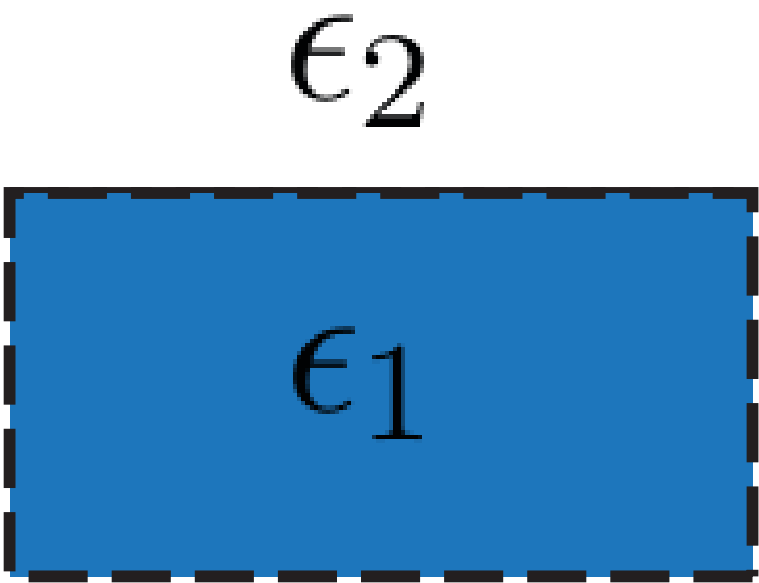
☆ boundary distortion
☆ strain



strain-optic effect



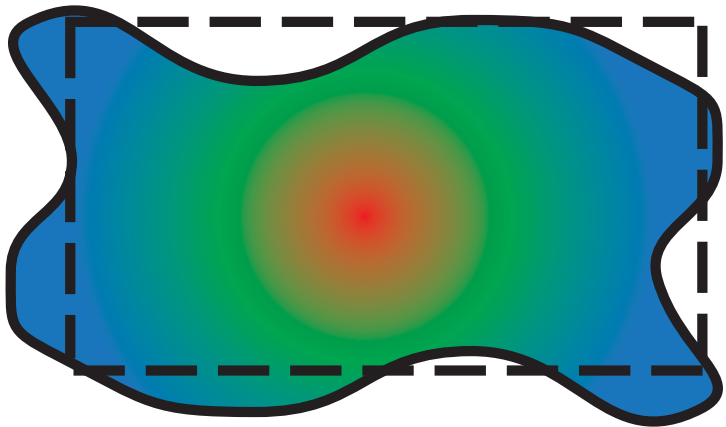
Light-sound interaction



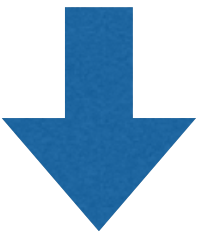
Dielectric waveguide



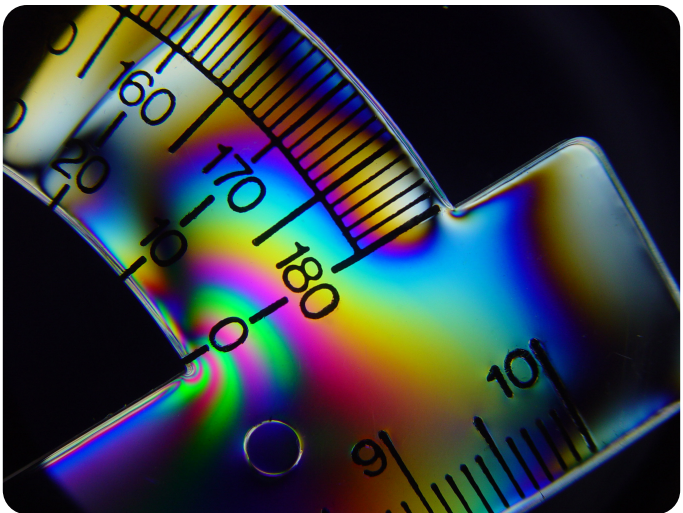
Mechanical mode Deformation



- ☆ boundary distortion
- ☆ strain



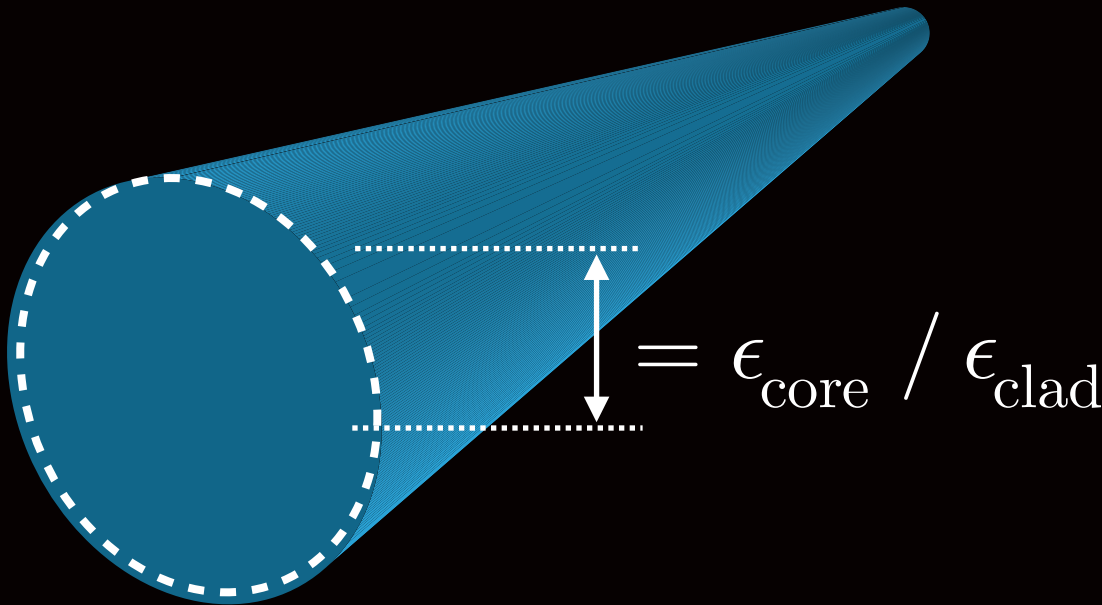
refractive index modulation



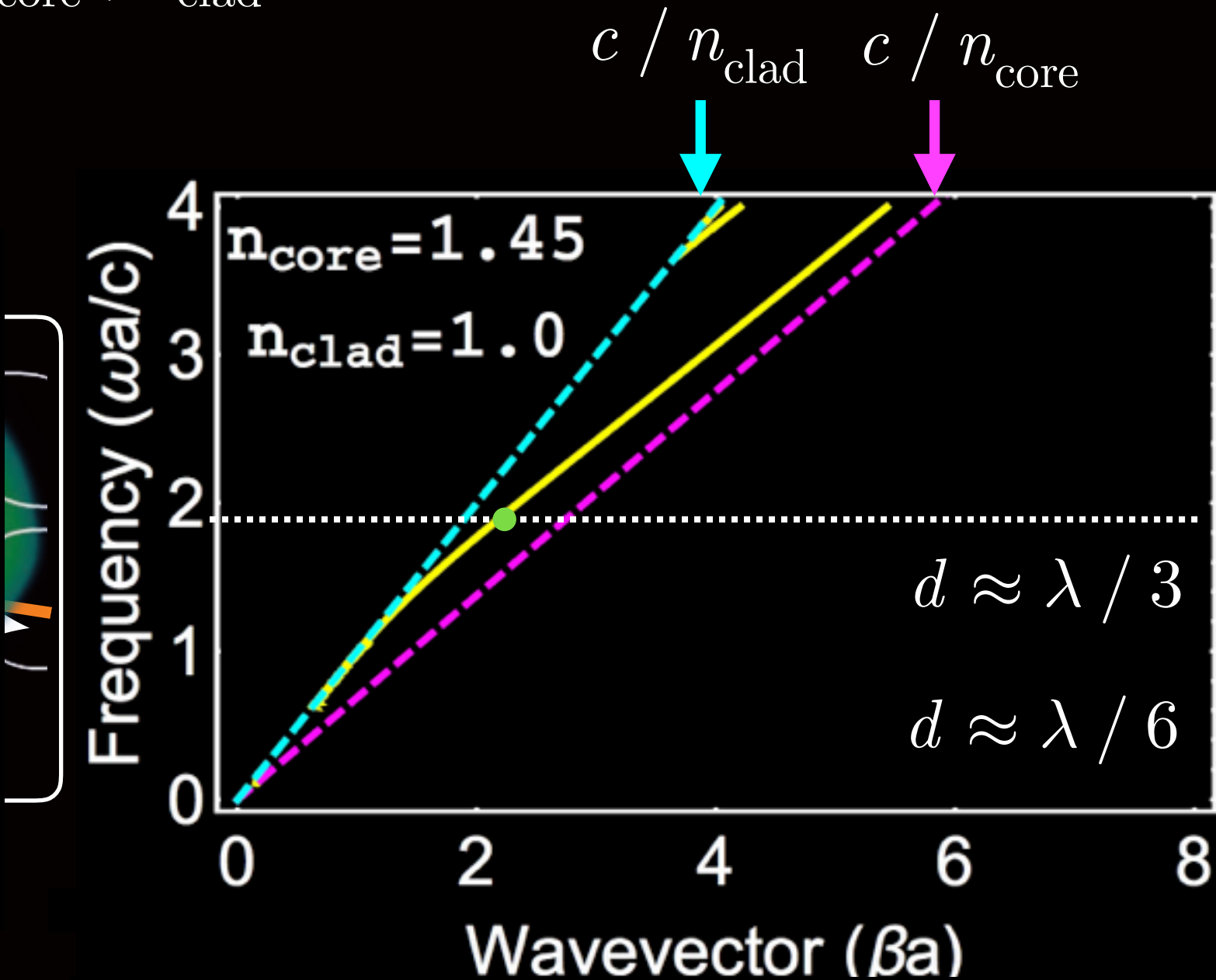
strain-optic effect



Sub-wavelength confinement



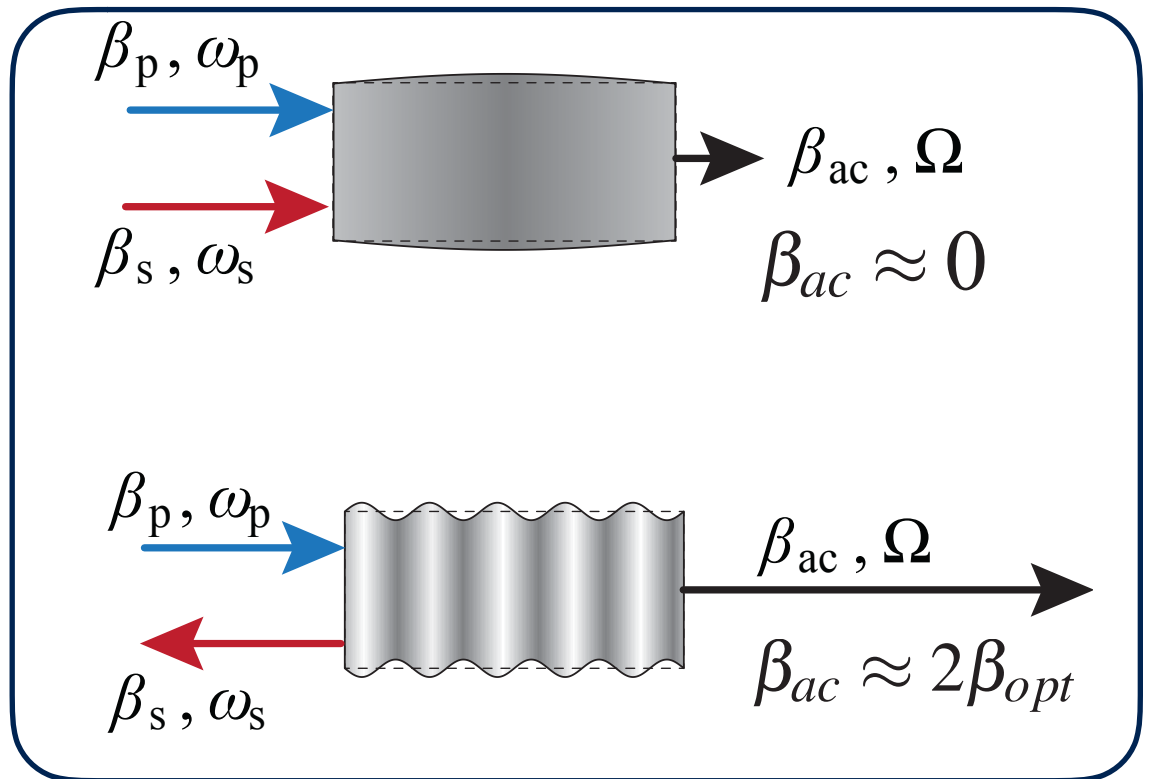
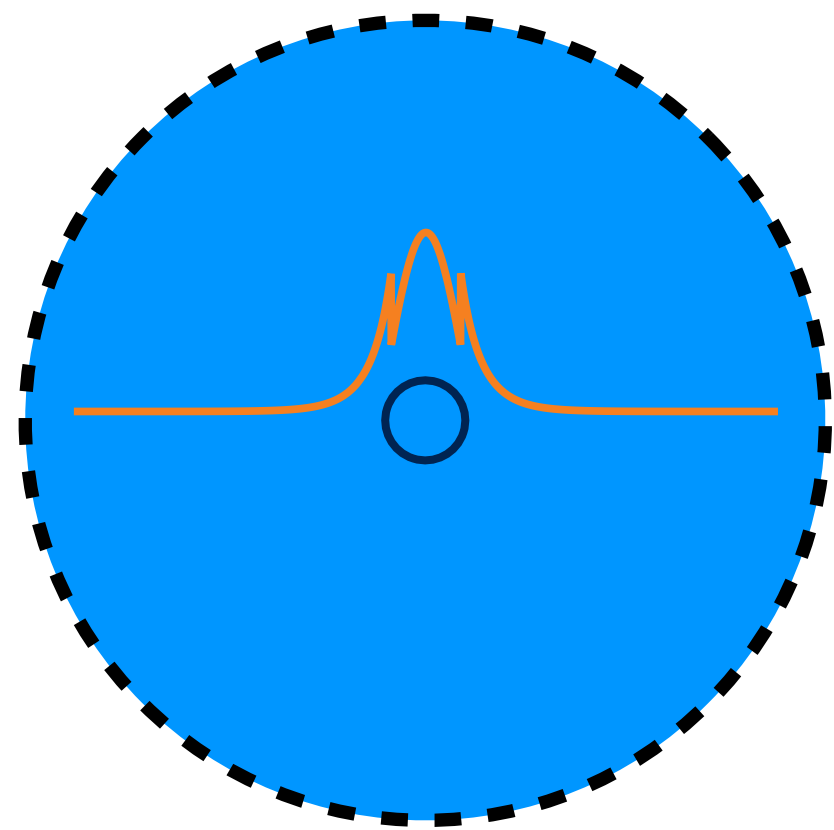
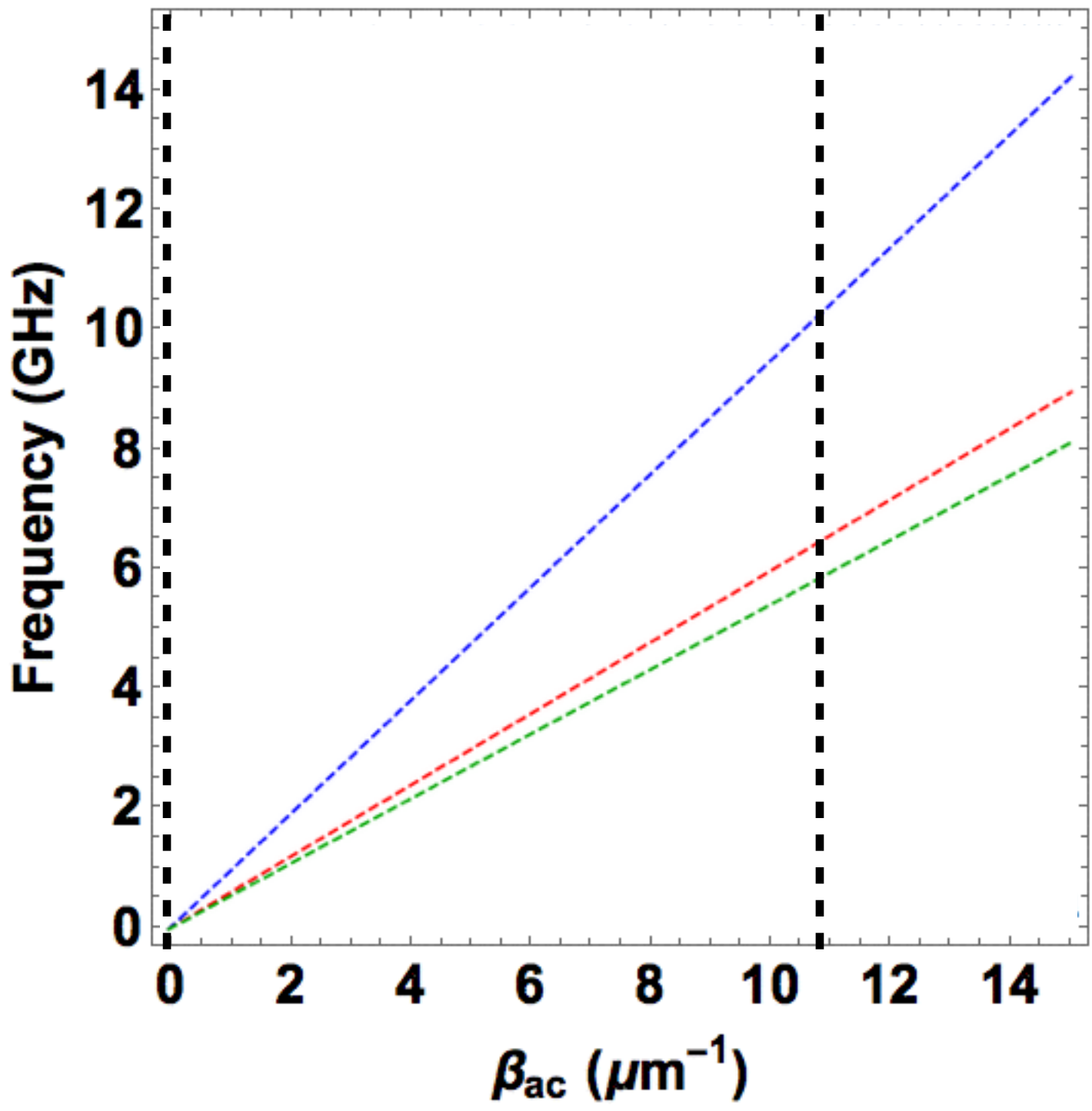
$$n_{\text{Si}} \approx 3.5$$
$$n_{\text{SiO}_2} \approx 1.45$$





Light-sound interaction: Brillouin scattering

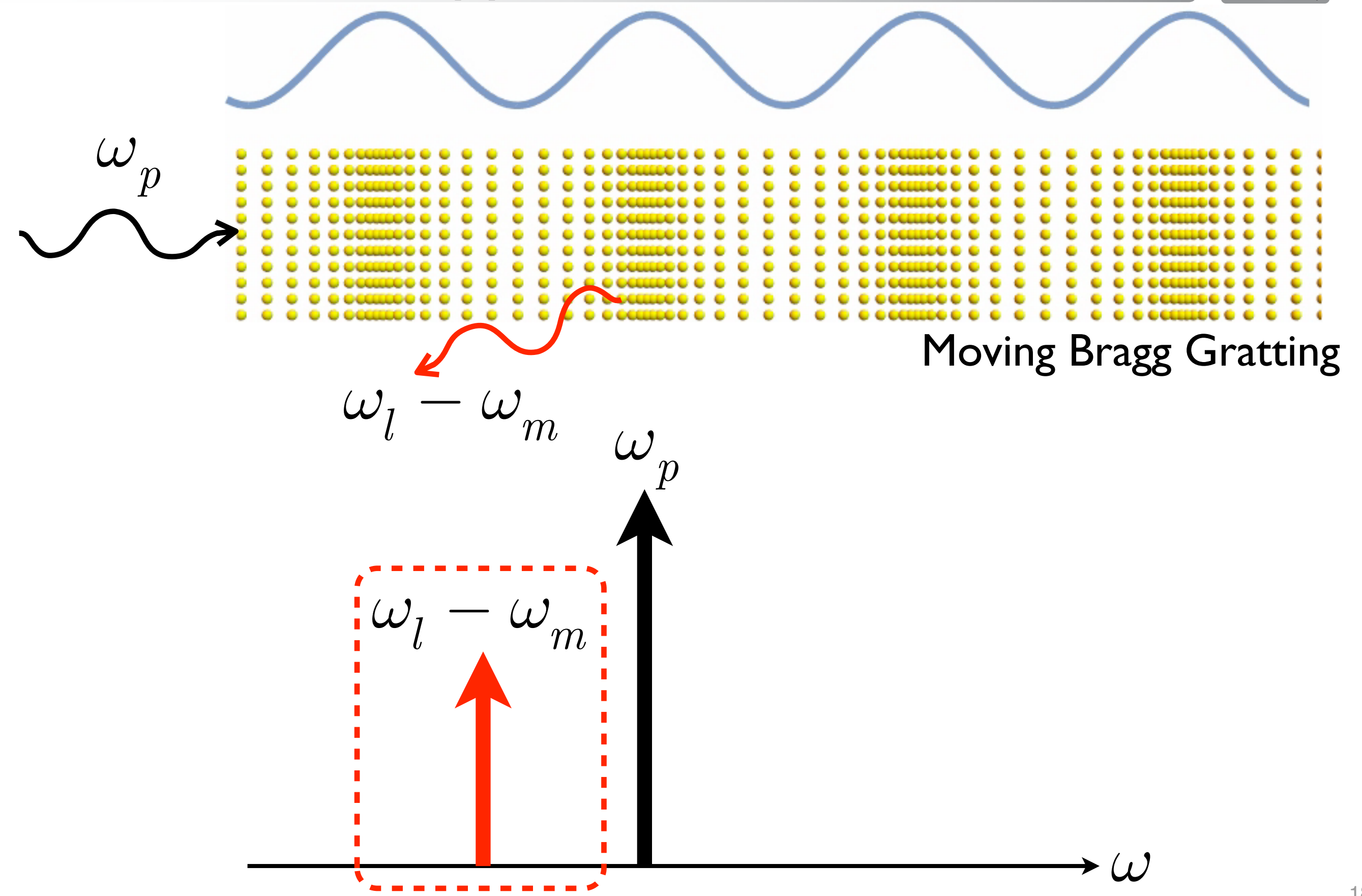
Dispersion relation: R_{0m} modes



Phase matching for Stokes scattering

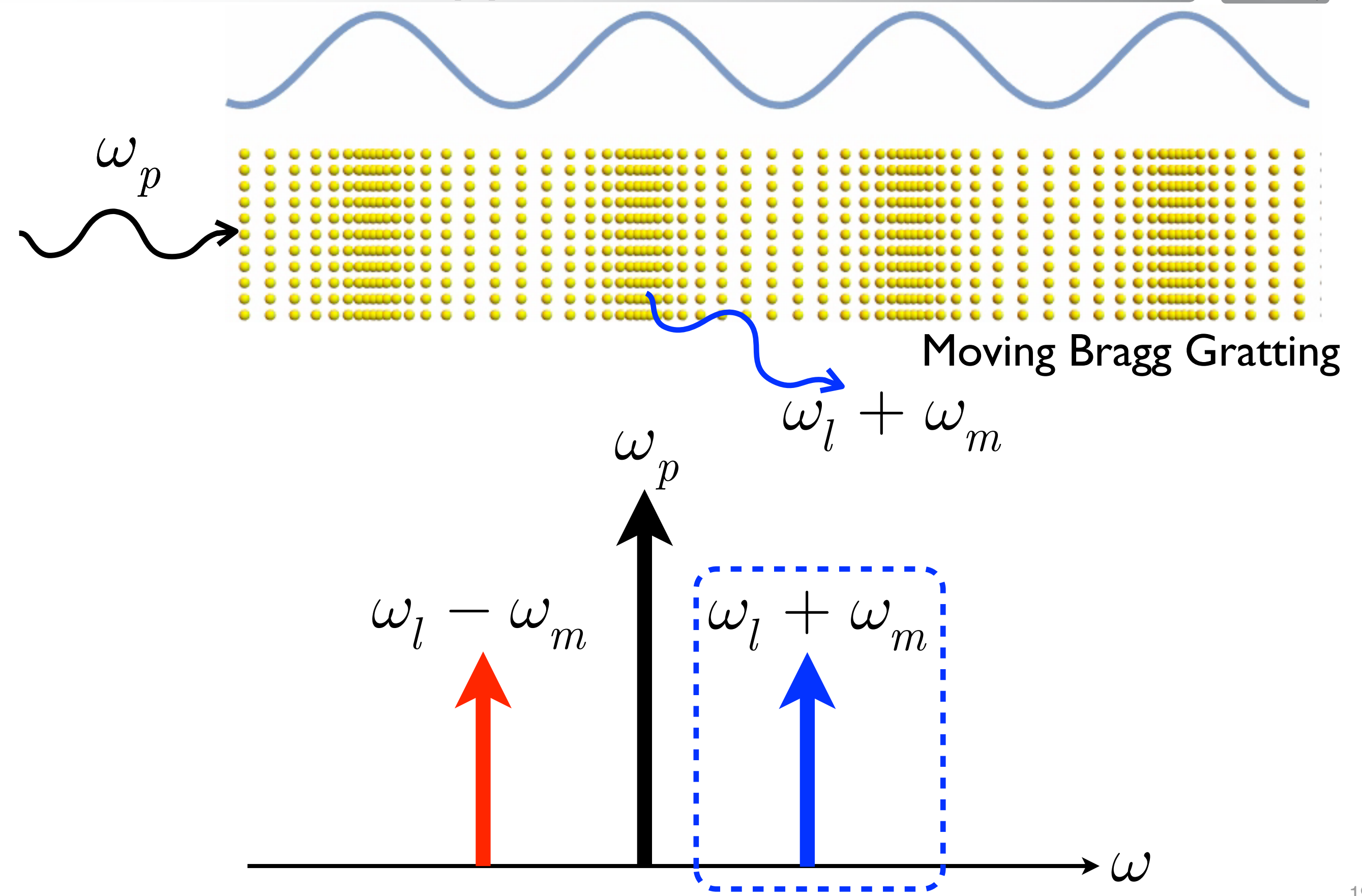


Photoelastic: Doppler Shift



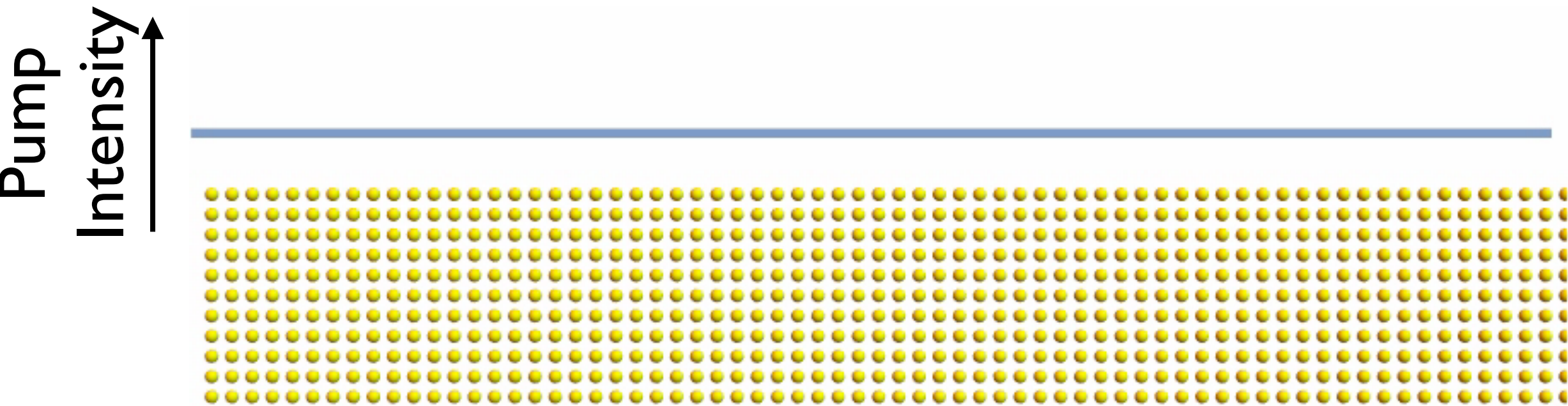


Photoelastic: Doppler Shift





Light-sound interaction: Brillouin scattering

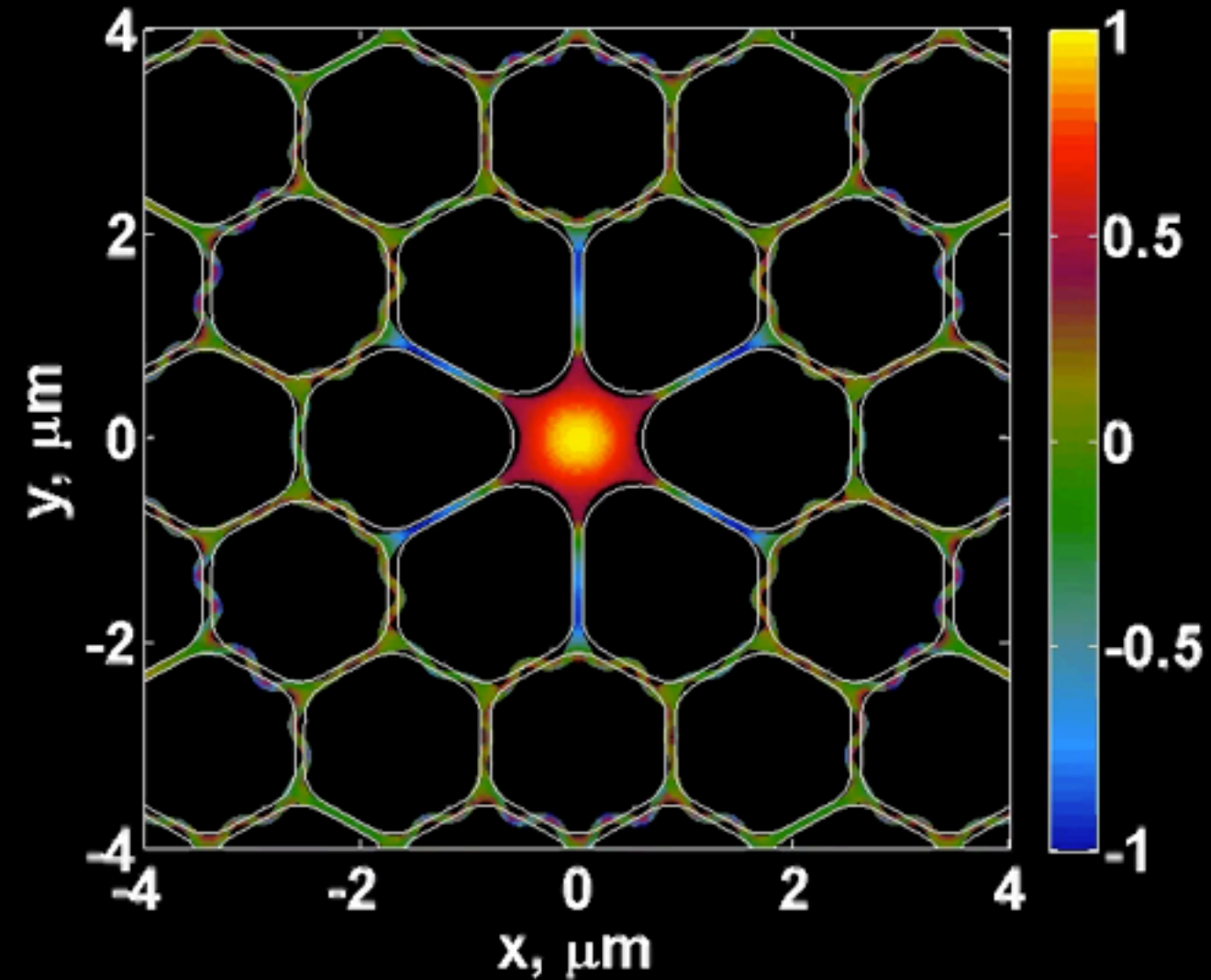
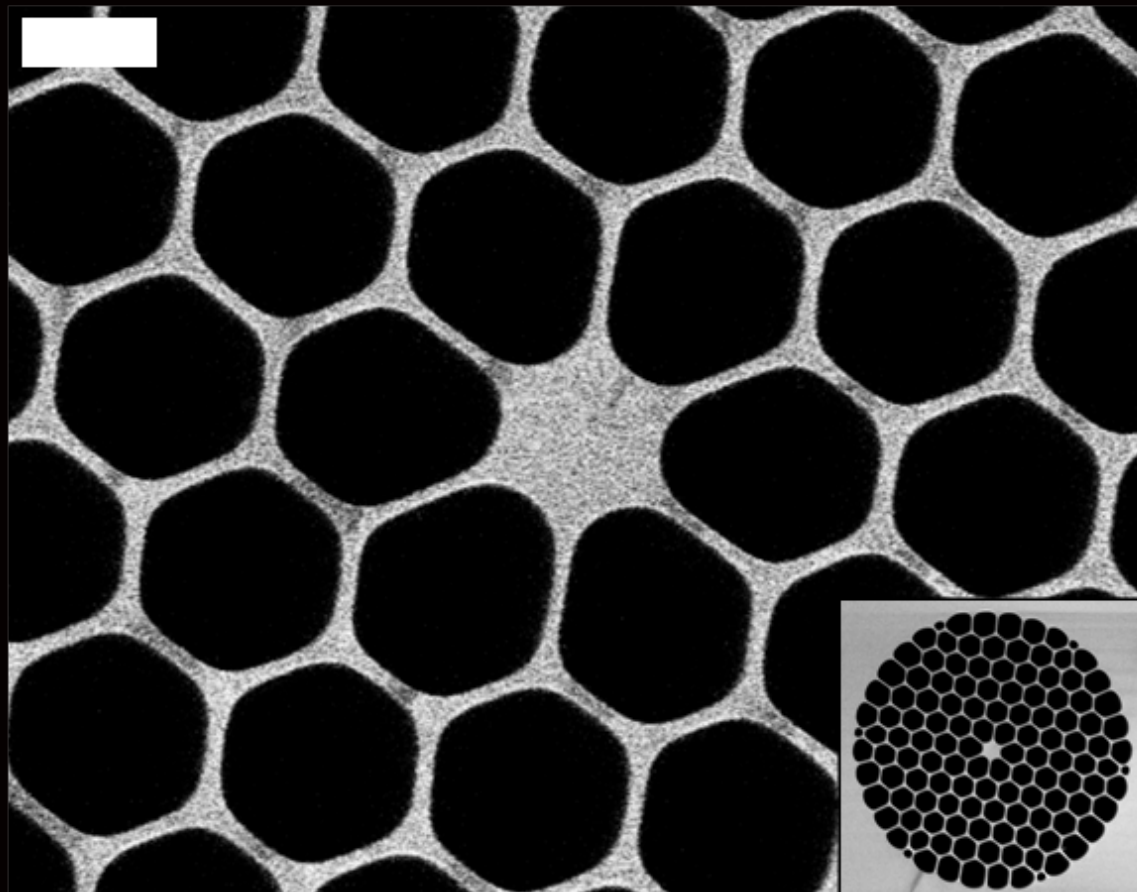


Steady Bragg Grating

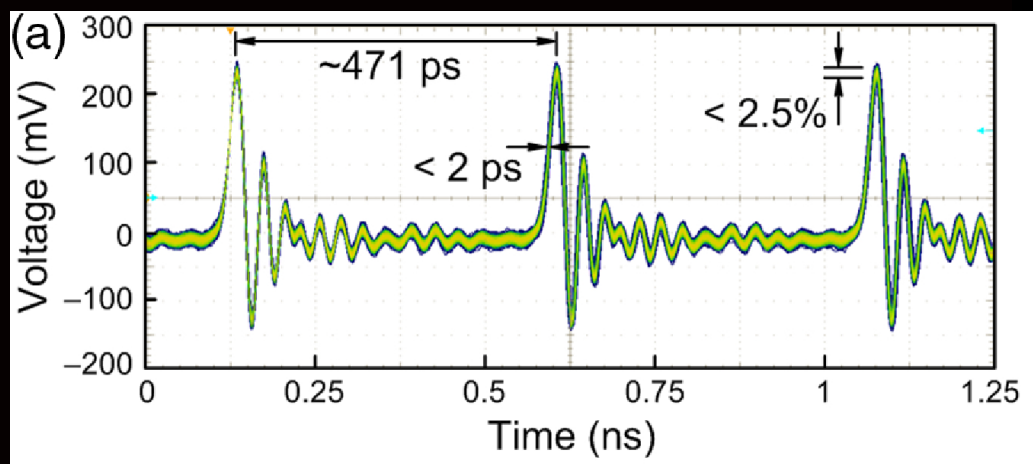
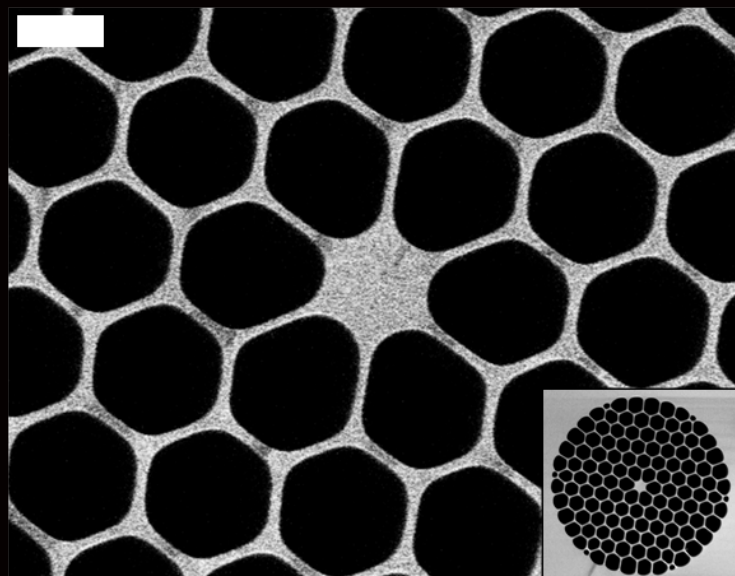


Brillouin Scattering in microwaves

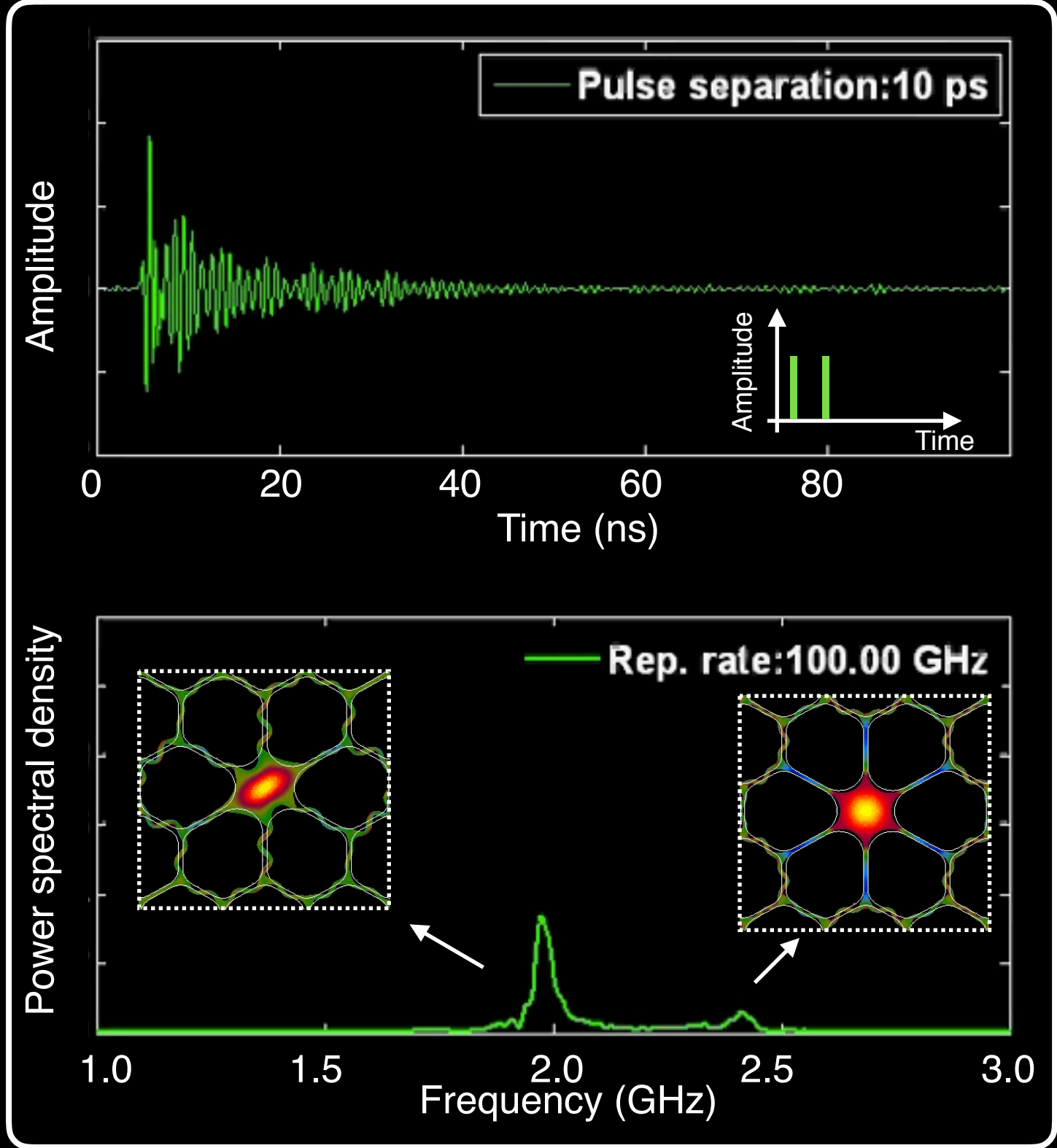
a.



- Dainese, P., et al. (2006). Nature Physics, 2(6), 388.
 Dainese, P., et al (2006). Optics Express, 14(9), 4141–4150
 Wiederhecker, G. S., et al. (2008). PRL, 100(20), 203903.
 Kang, M., et al (2008). Applied Physics Letters, 93, 131110.
 Brenn, A., et al (2009). Josa B, 26(8), 1641–1648.

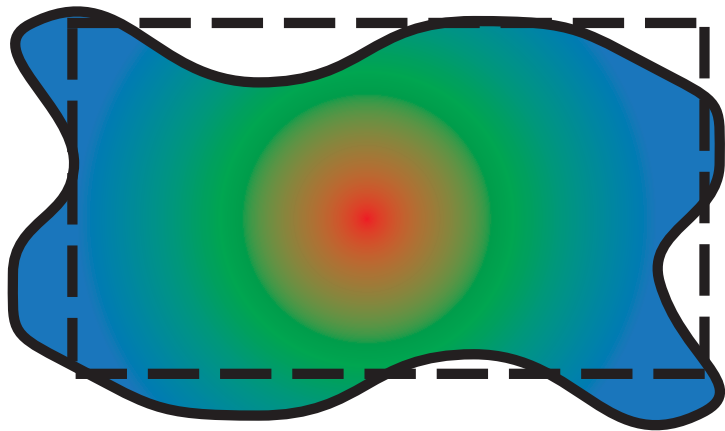


W. He, et al, Optics Express, 23(19), 24945-24954 (2015)
M. Pang, Optica, Vol. 2, Issue 4, pp. 339-342 (2015)

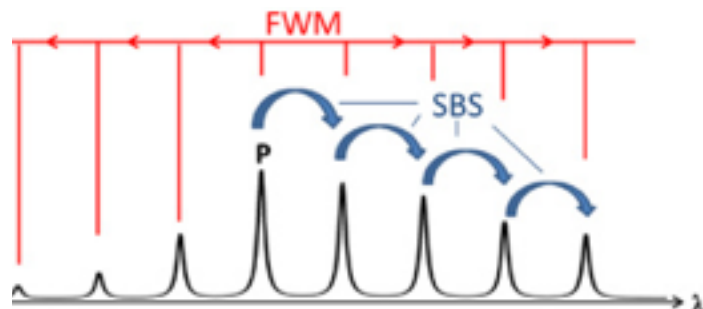




Light-sound interaction: Brillouin scattering

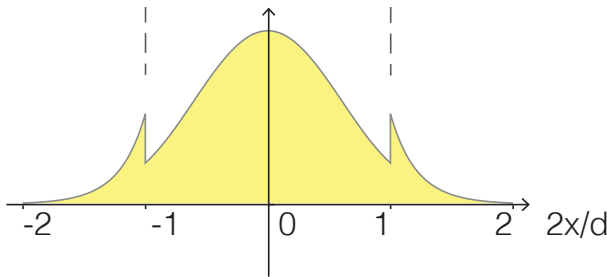
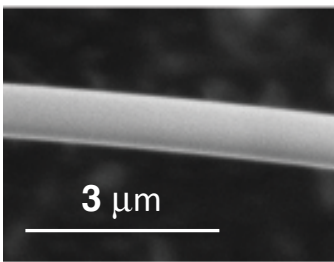


Cascaded Brillouin Scattering



Büttner T.F.S. et al., *Optica* 1, 311-314 (2014)

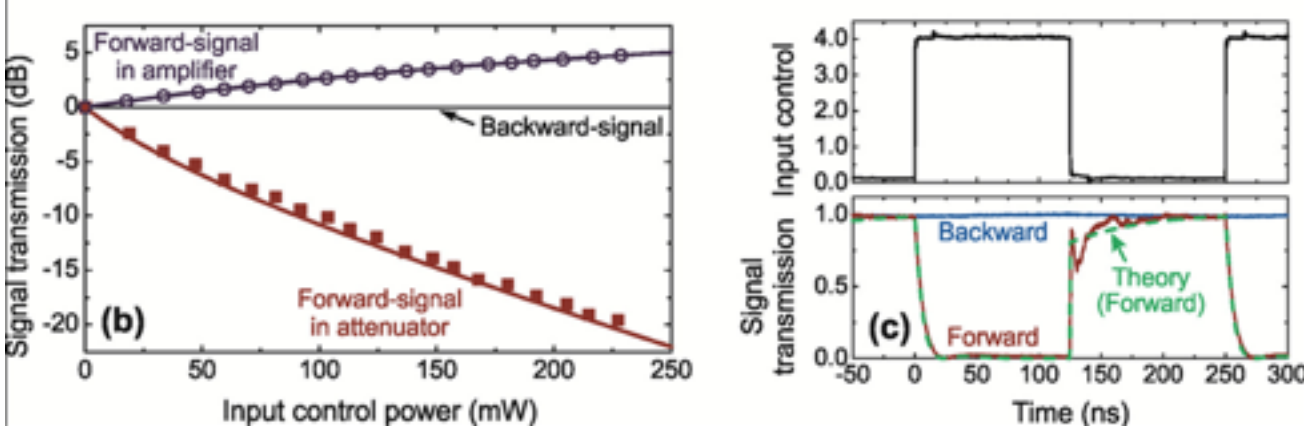
Brillouin Self-Cancellation



Florez, O., et al. *Nature Communications*, 7, 11759.

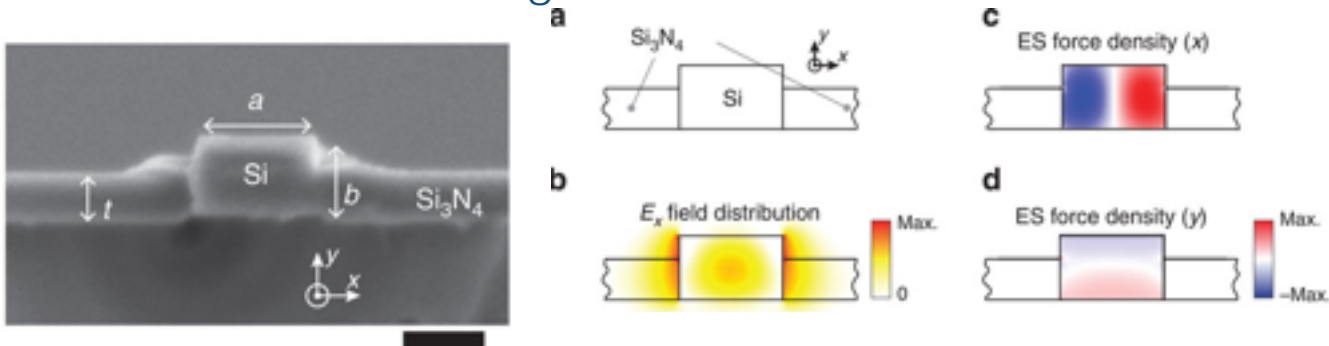
S. G. Johnson et al., *Physical Review E* **65**, (2002)

Optical isolator and switching



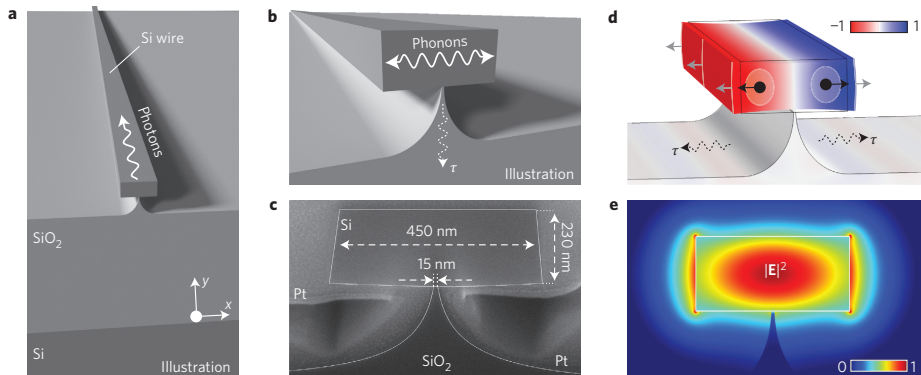
M. Kang et al. *Nature Photonics* **5**, 549–553 (2011).

Brillouin in silicon waveguides



Shin, H., et al. *Nature Communications*, 4. (2013)

Brillouin amplification

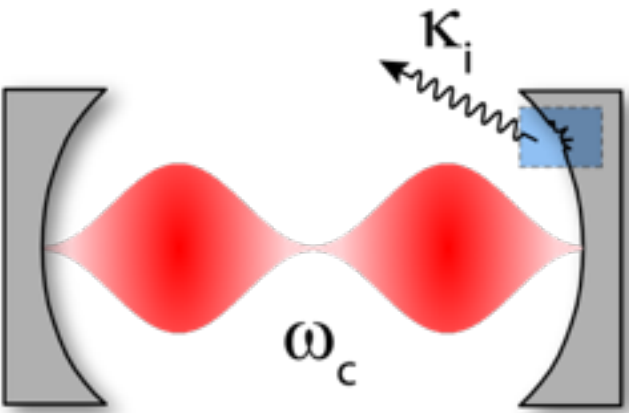


R. Van Laer et al, *Nature Photonics* **9**, 199–203 (2015)

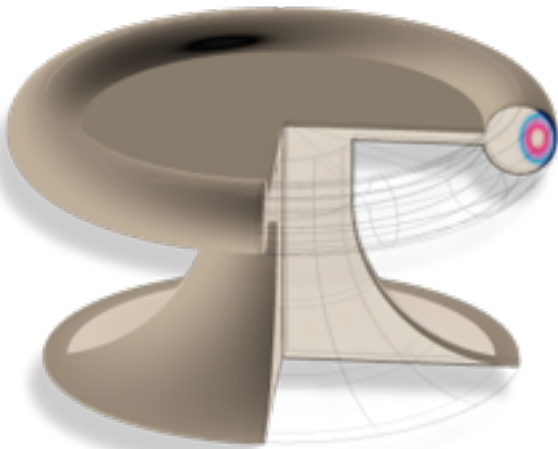


Examples of optical cavities

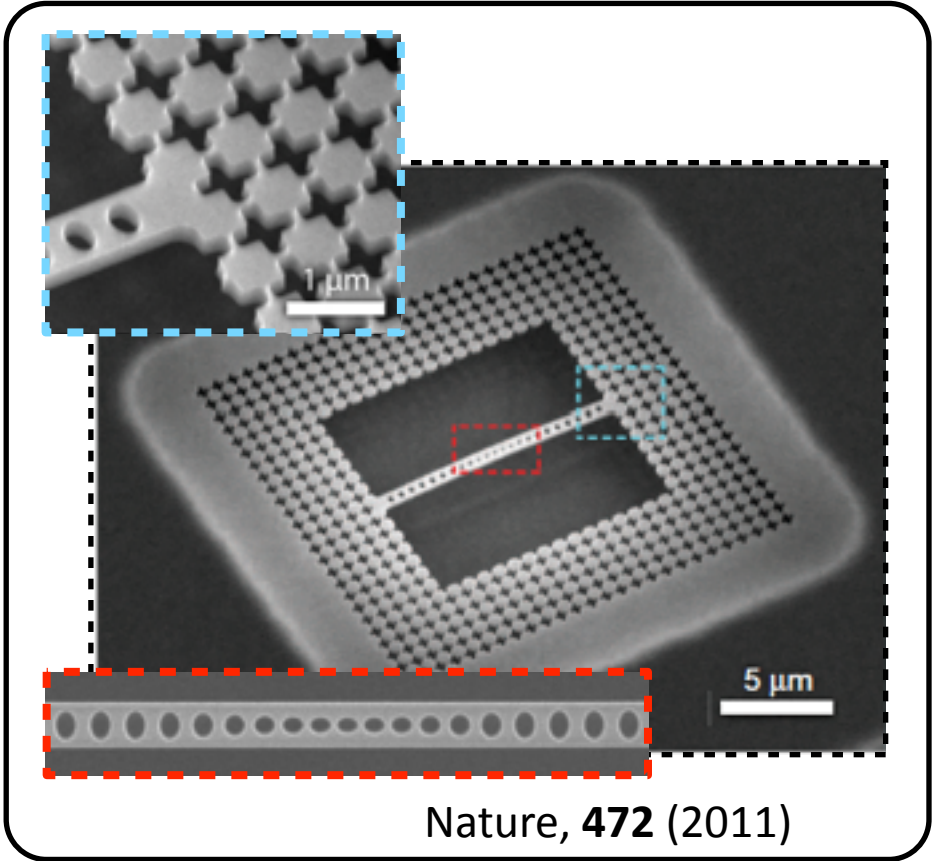
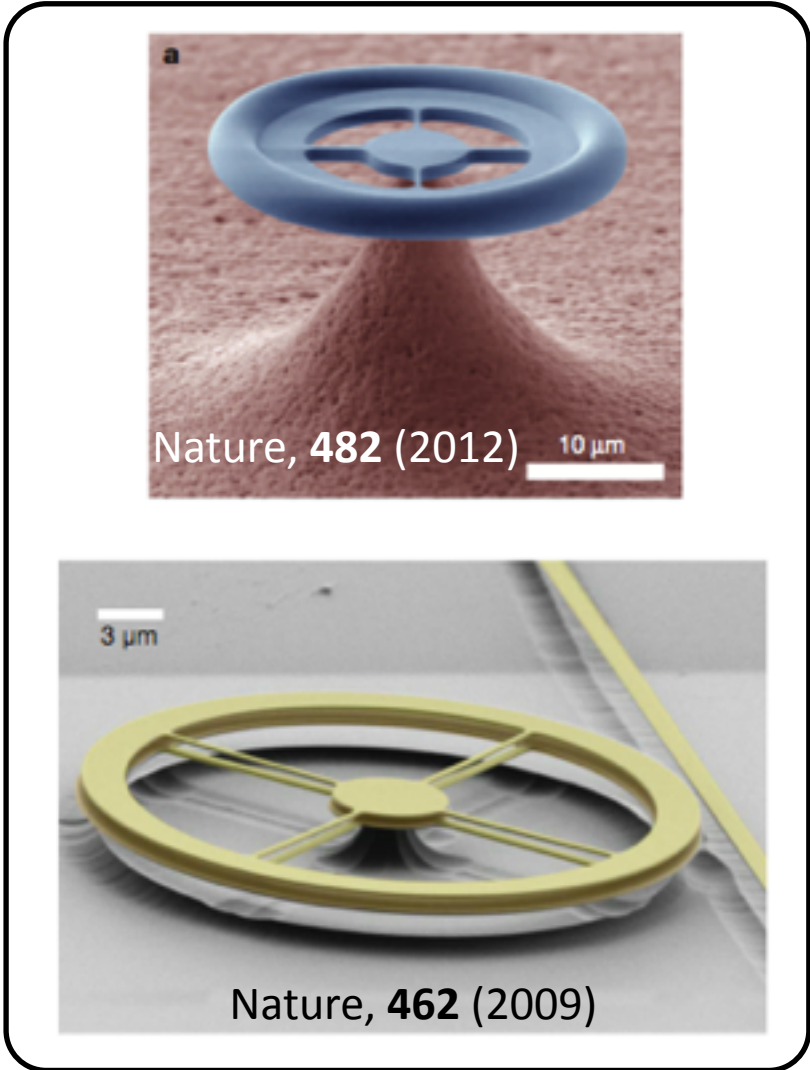
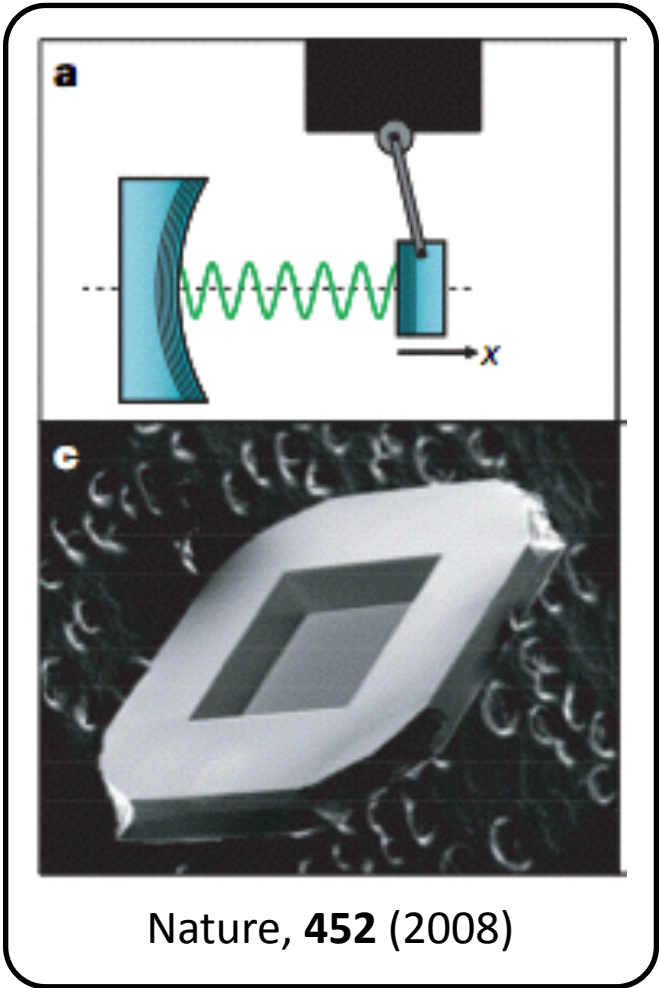
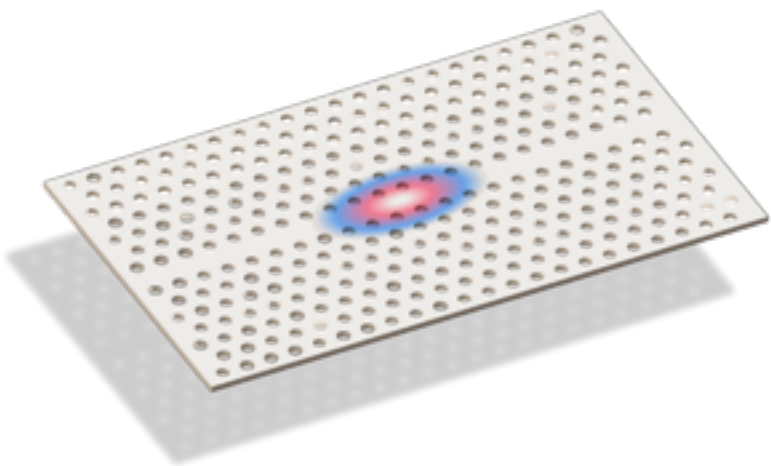
Fabry-Perot



Ring and disk resonator

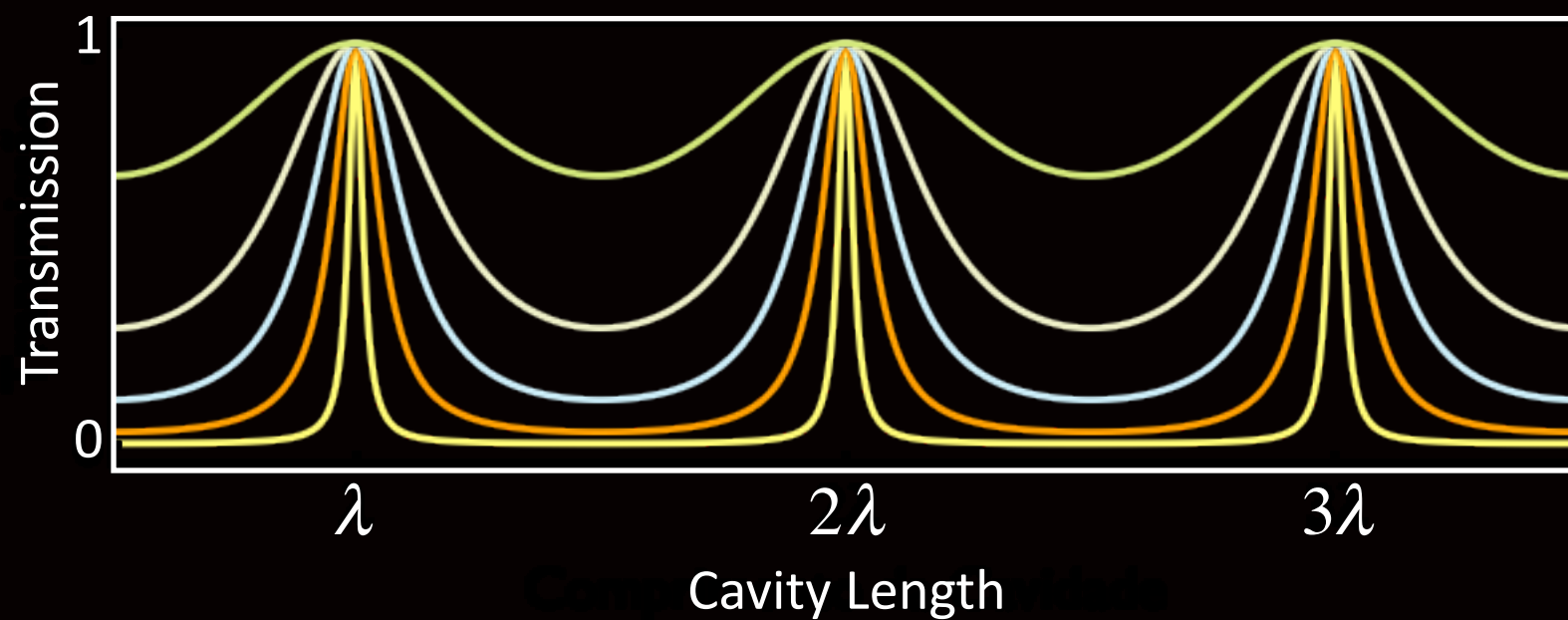
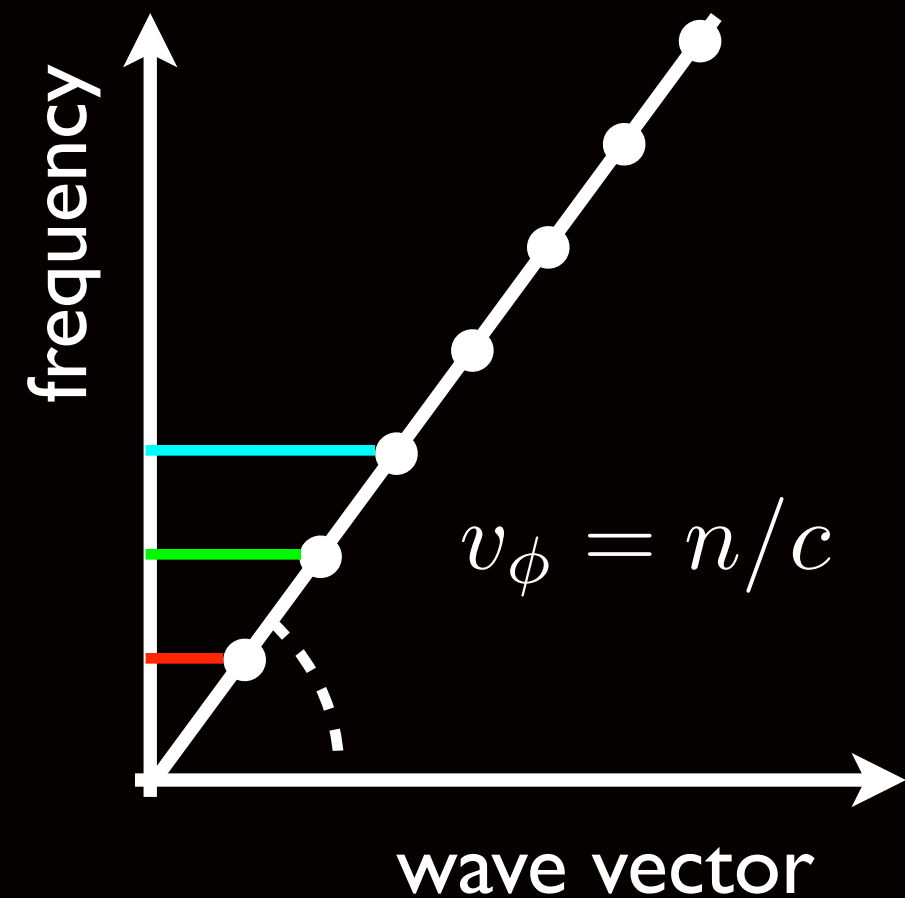
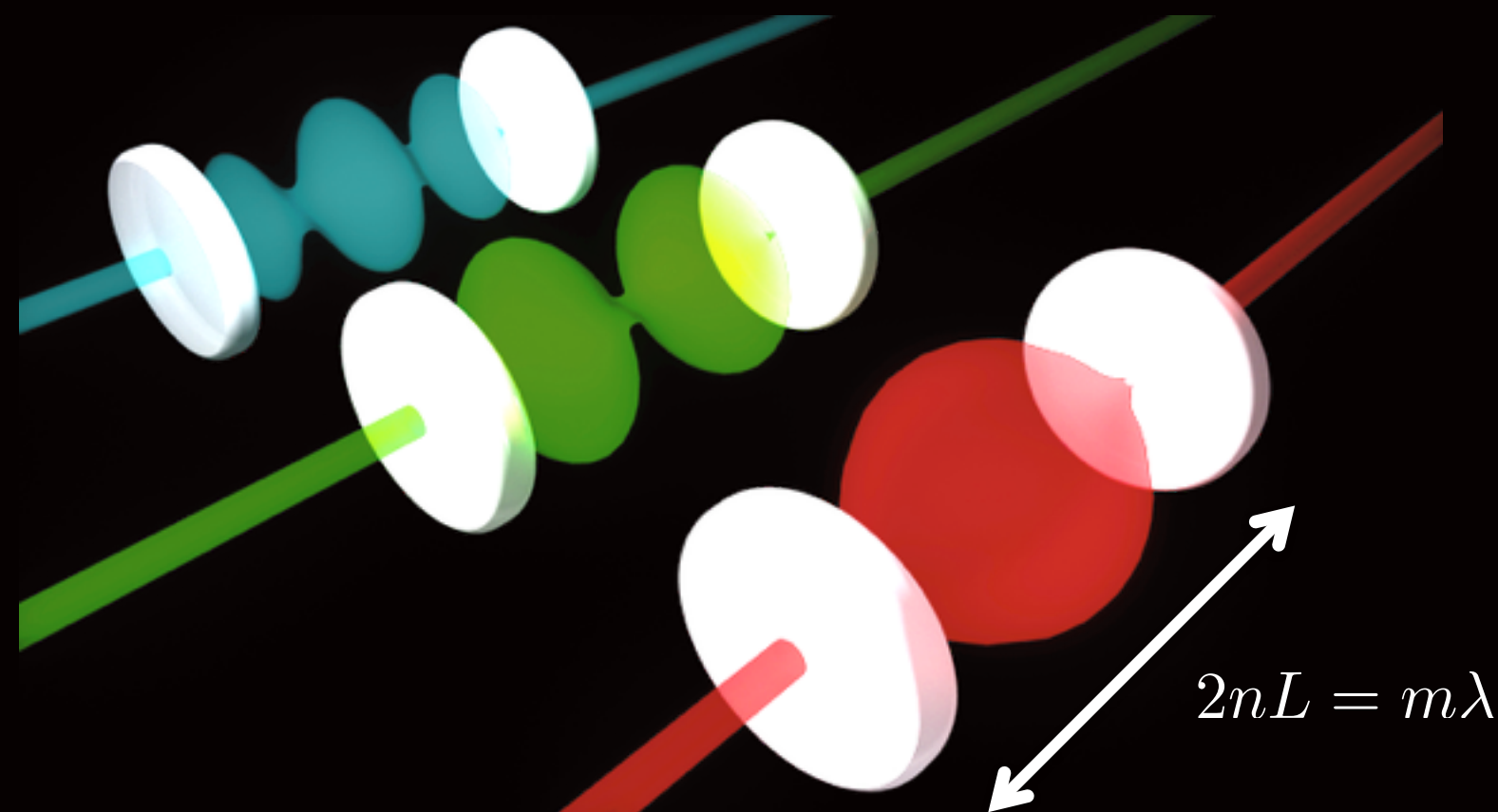


PhC





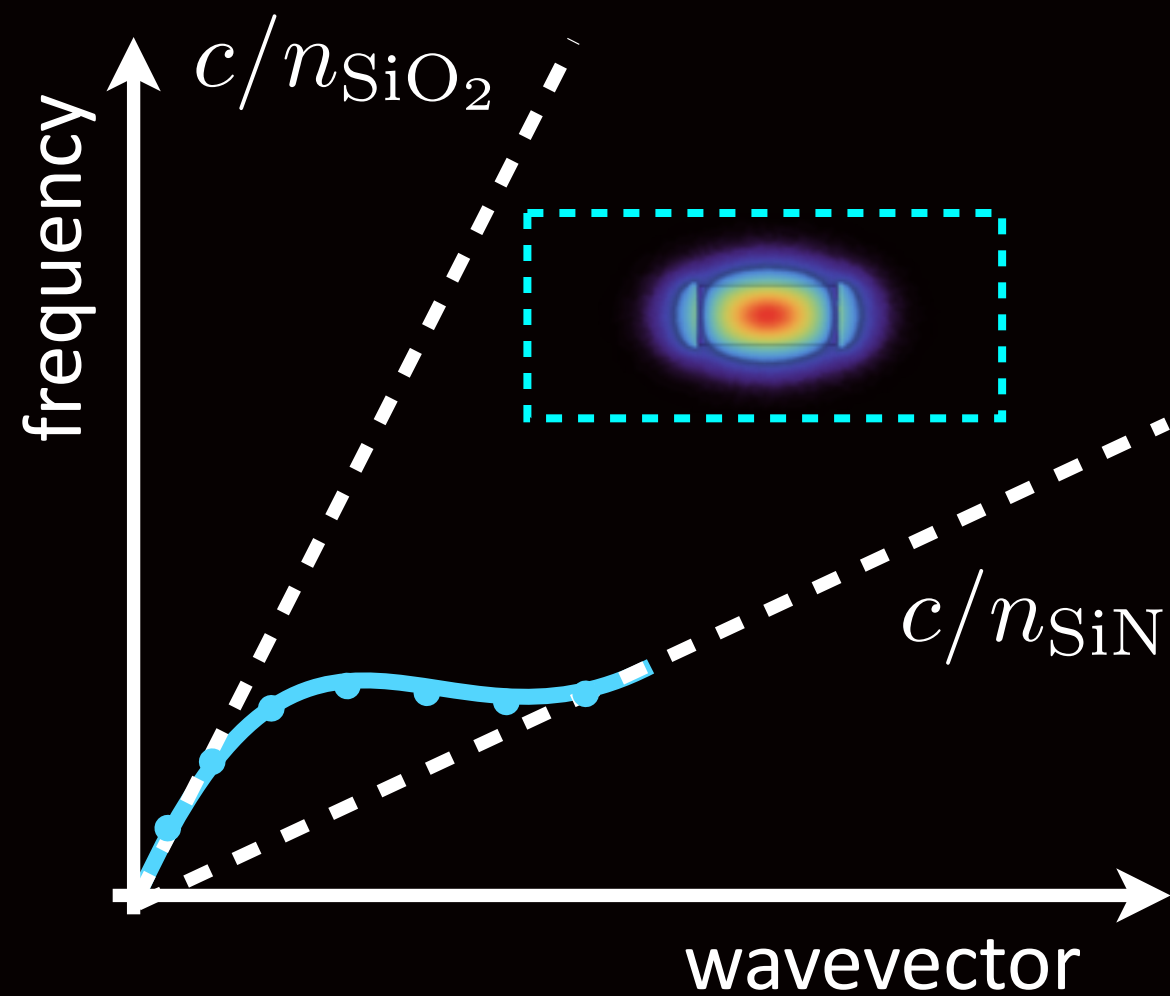
Optical Cavities



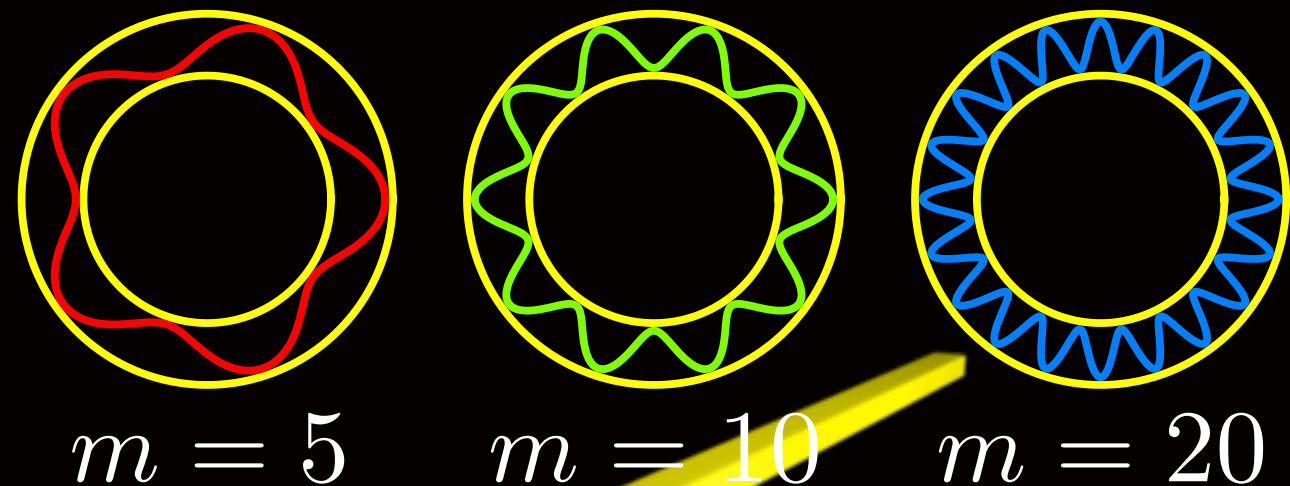
$$T = \left| \sum_{m=0}^{\infty} R^m e^{im\delta} \right|^2$$



Ring resonators



$$\omega(k) = k_z \frac{c}{n}$$



Power enhancement

$$P_{\text{circ}} \propto \mathcal{F} P_0$$

$$R = 15 \text{ } \mu\text{m}$$

$$Q = 10^6 \text{ } (\tau \approx 1 \text{ ns})$$

$$\mathcal{F} \approx 50 \times 10^3$$



A brief story of micro cavities

Volume 137, number 7,8

PHYSICS LETTERS A

QUALITY-FACTOR AND NONLINEAR PROPERTIES OF OPTICAL WHISPERING-GALLERY MODES

V.B. BRAGINSKY, M.L. GORODETSKY and V.S. ILCHENKO

Department of Physics, Moscow University, 119899 Moscow, USSR

Received 10 March 1989; accepted for publication 21 March 1989

Communicated by V.M. Agranovich

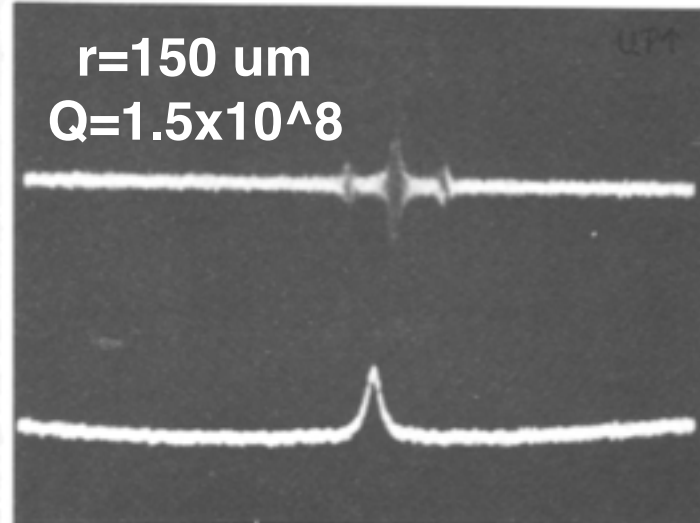


Fig. 2. Detailed resonant curve of the mode; second trace is the precise marks ± 7 MHz. Bandwidth of the mode ≈ 3 MHz, quality-factor 1.5×10^8 (the whispering-gallery microresonator $150 \text{ } \mu\text{m}$ in diameter).

Braginsky, V. B., et al. Physics Letters A, 137(7-8), 393–397. (1989)

Collot, L., Europhys. Lett. 23, 327-334 (1993).

M. L. Gorodetsky et al, Opt. Commun.113, 133 (1994).

J. C. Knight, et al. Opt. Lett. 20, 1515-1517 (1995)

D. W. Vernooy and H. J. Kimble, Phys. Rev. A 55, 1239 (1997).



A brief story of micro cavities

Volume 137, number 7,8

PHYSICS LETTERS A

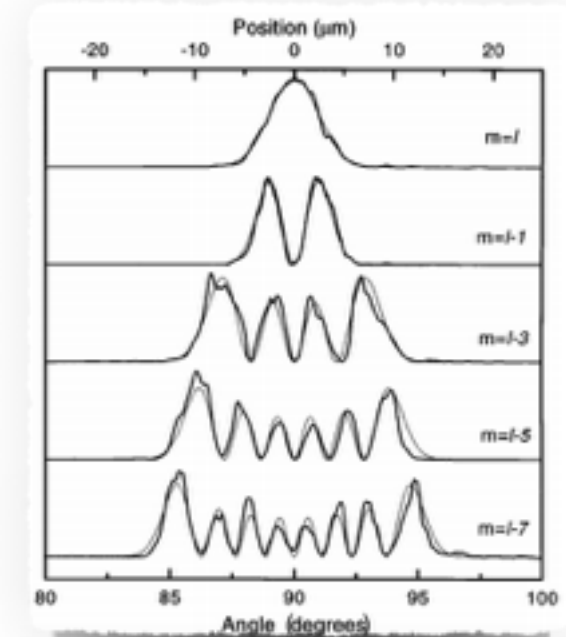
QUALITY-FACTOR AND NONLINEAR PROPERTIES OF OPTICAL WHISPERING-GALLERY MODES

V.B. BRAGINSKY, M.L. GORODETSKY and V.S. ILCHENKO

Department of Physics, Moscow University, 119899 Moscow, USSR

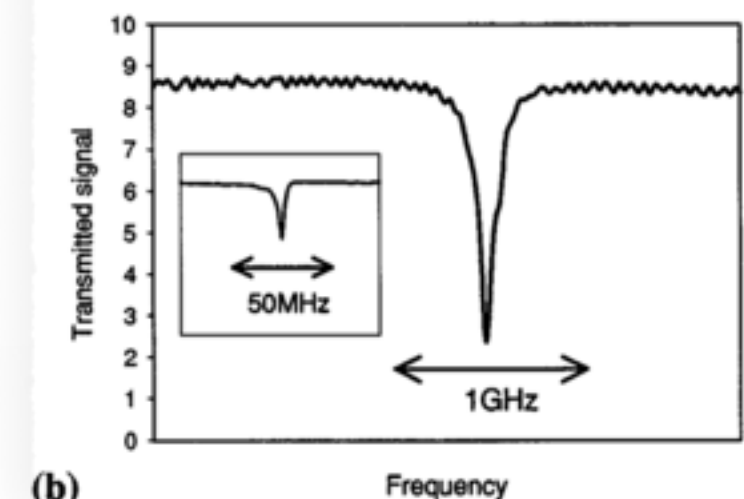
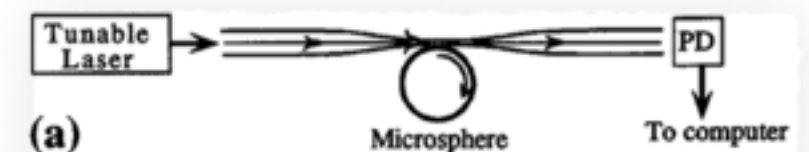
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J. C. Knight, et al.

Opt. Lett. 20, 1515-1517 (1995)



Knight, J. C., et al.

Optics Letters, 22(15), 1129–1131 (1997).

Braginsky, V. B., et al. Physics Letters A, 137(7-8), 393–397. (1989)

Collot, L., Europhys. Lett. 23, 327-334 (1993).

M. L. Gorodetsky et al, Opt. Commun.113, 133 (1994).

J. C. Knight, et al. Opt. Lett. 20, 1515-1517 (1995)

D. W. Vernooy and H. J. Kimble, Phys. Rev. A 55, 1239 (1997).



A brief story of micro cavities



Nature 415, 621-623 (2002).
Silica Raman Laser



Nature 421, 925-928 (2003)
Ultra-high Q thoroids

Spillane, S. M., et al. Nature 415, 621-623 (2002)
D. K. Armani, et al., Nature 421, 925-928 (2003)



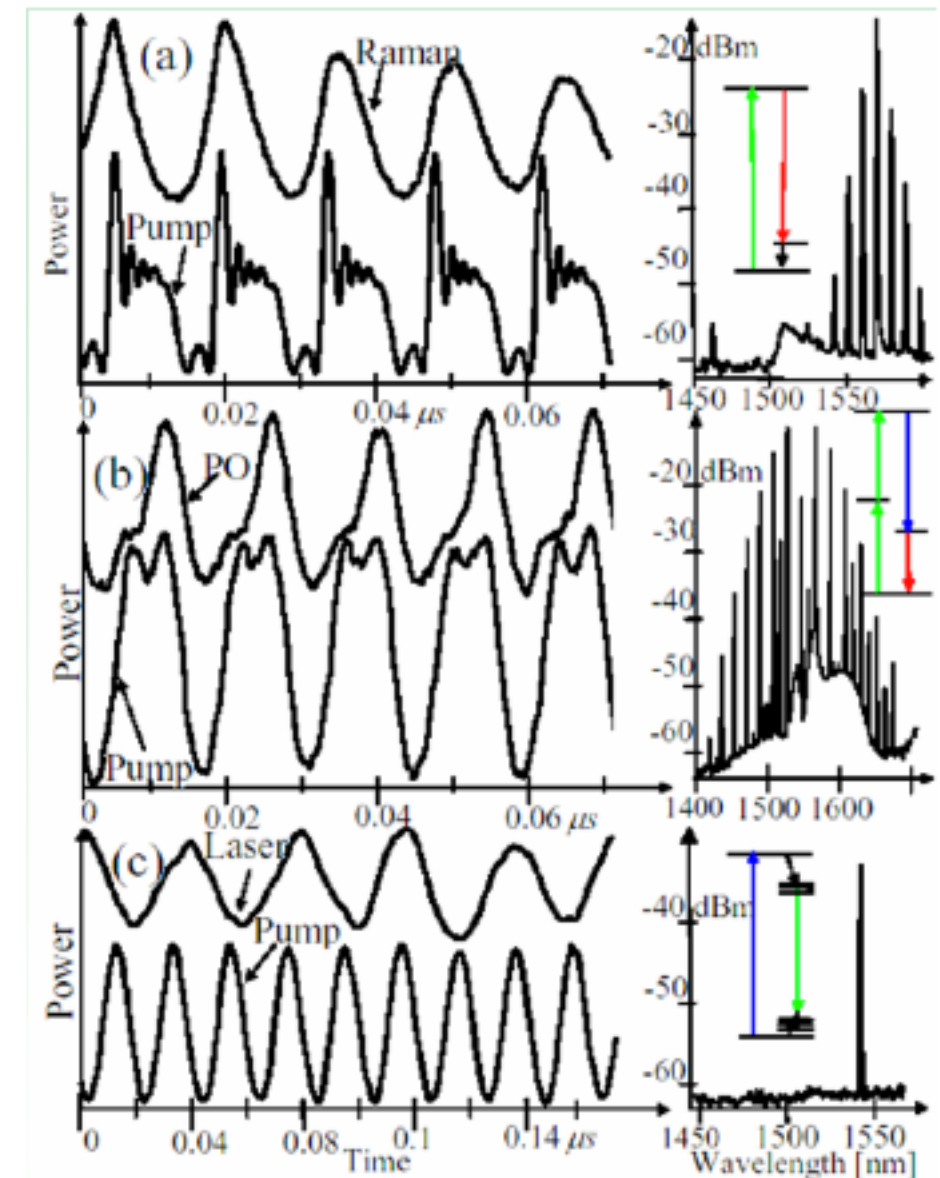
A brief story of micro cavities



Nature 415, 621-623 (2002).
Silica Raman Laser



Nature 421, 925-928 (2003)
Ultra-high Q thoroids



Carmon, CLEO 2005

Carmon, T., et al.
Physical Review Letters, 94(22), 223902. (2005)

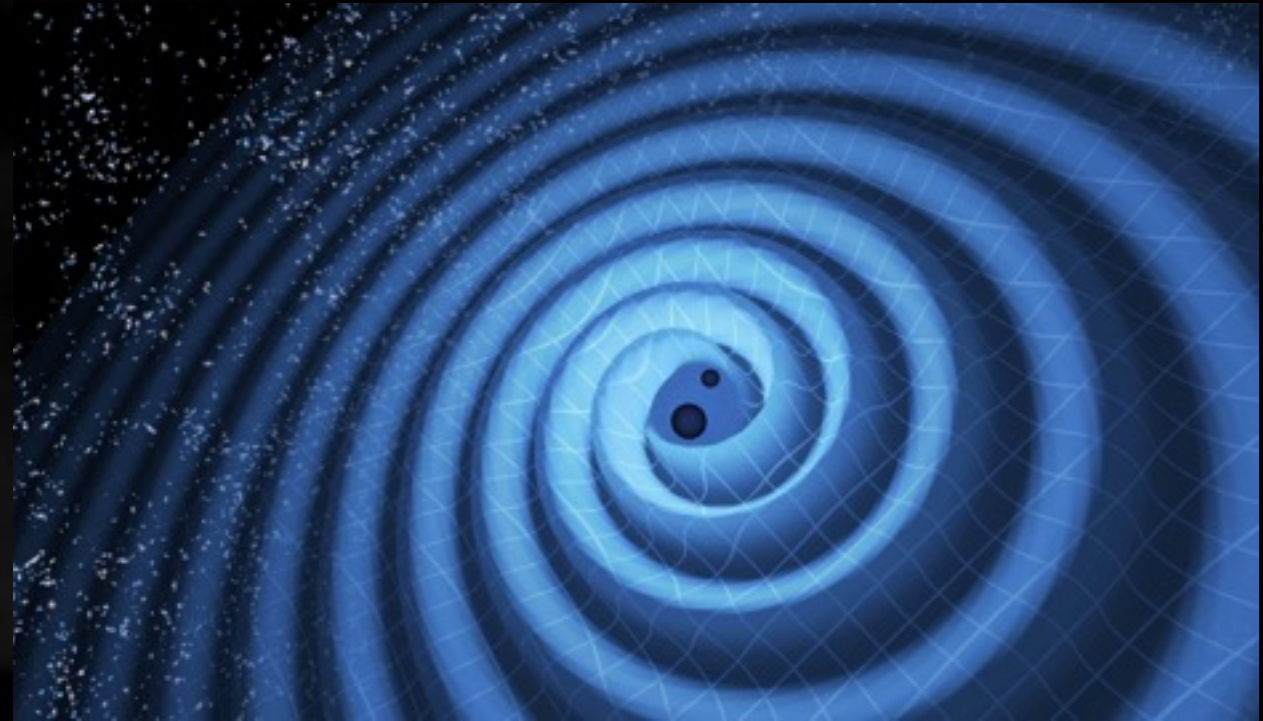
Spillane, S. M., et al. Nature 415, 621-623 (2002)
D. K. Armani, et al., Nature 421, 925-928 (2003)



Standing on the shoulder of giants



Vladimir Braginsky, 1931-2016



Credit: LIGO website

V.B. Braginsky, Y.I. Vorontsov, K.S. Thorne: *Science* **209**, 547 (1980)

V.B. Braginsky, S.E. Strigin, S.P. Vyatchanin, Parametric oscillatory instability in Fabry–Perot interferometer, *Physics Letters A*, Volume 287, Issues 5–6, 3 September 2001, Pages 331-338,



Mechanical effects of light

Photons' linear momentum results in *radiation pressure*

$$\vec{p} = \hbar \vec{k}$$



tail



Mechanical effects of light

Photons' linear momentum results in *radiation pressure*

$$\vec{p} = \hbar \vec{k}$$



tail

Radiation pressure driven



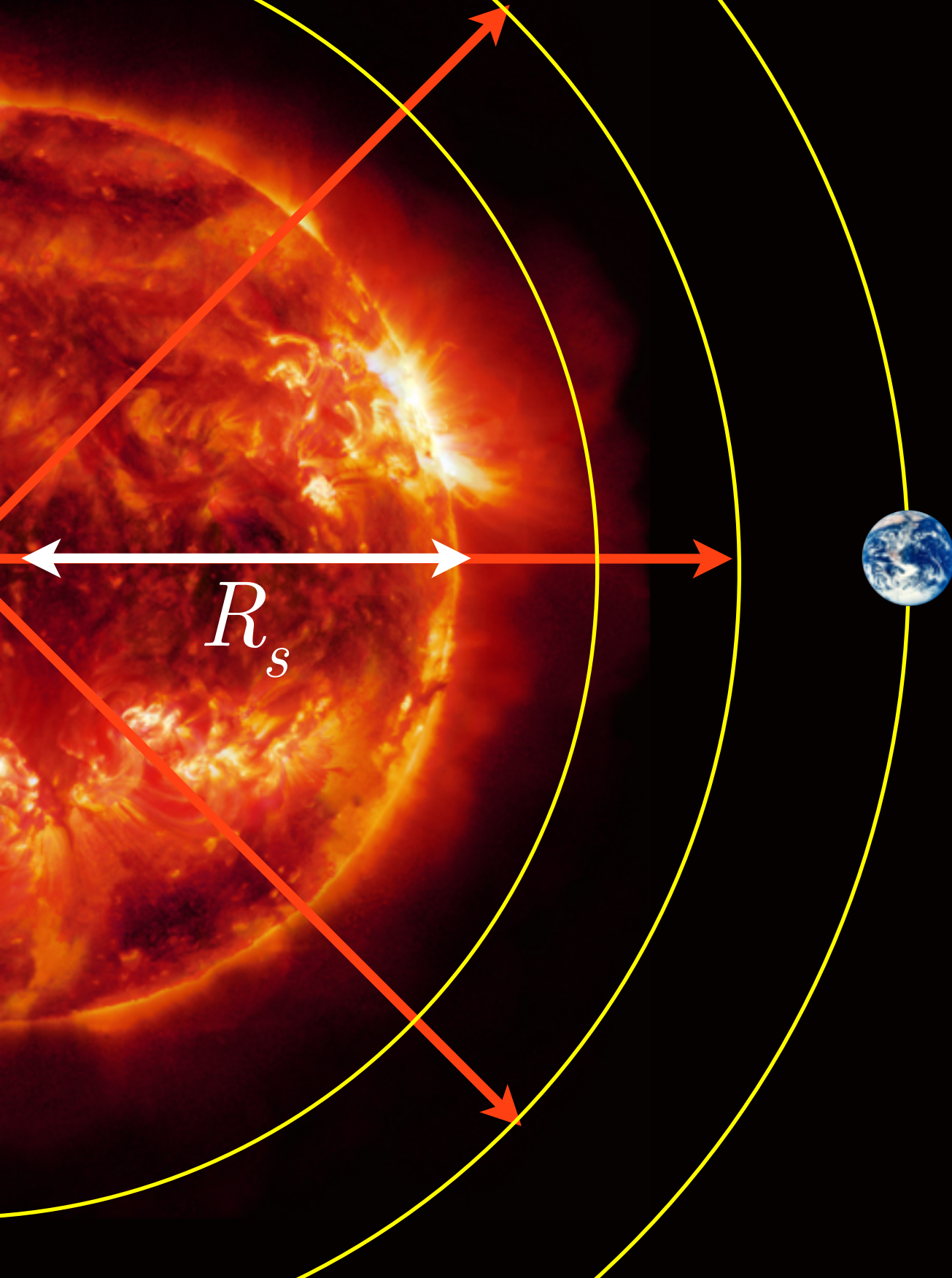
Solar wind driven





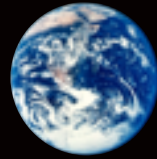
Radiation pressure force

$$P_{rad} = (4\pi R_s^2)\sigma T^4$$





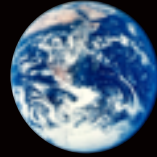
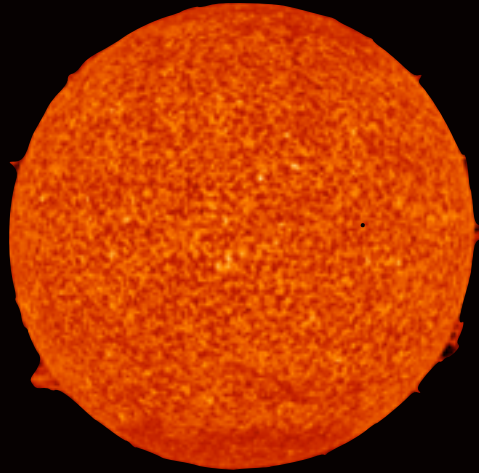
Sun vs laser pointer



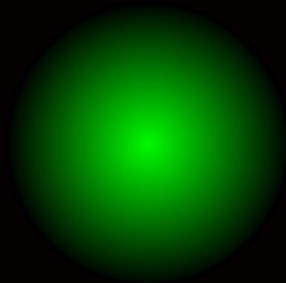
$$F \approx \frac{P}{c} \approx 0.6 \times 10^9 \text{ N}$$



Sun vs laser pointer



$$F \approx \frac{P}{c} \approx 0.6 \times 10^9 \text{ N}$$

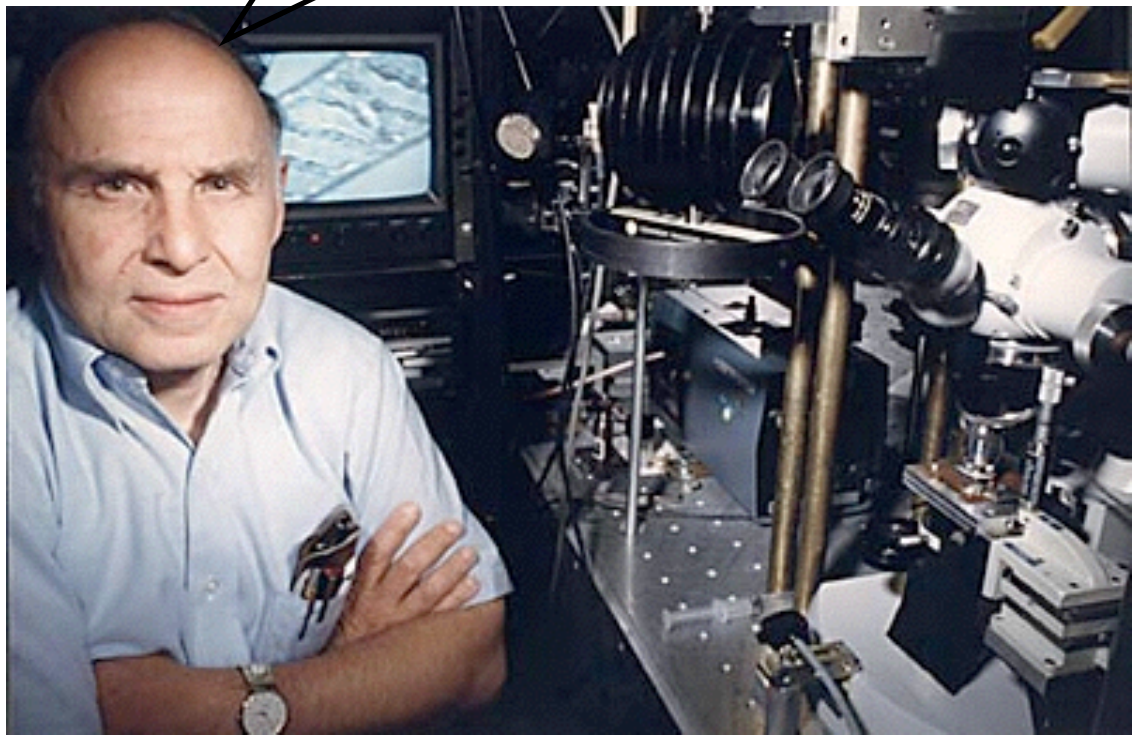
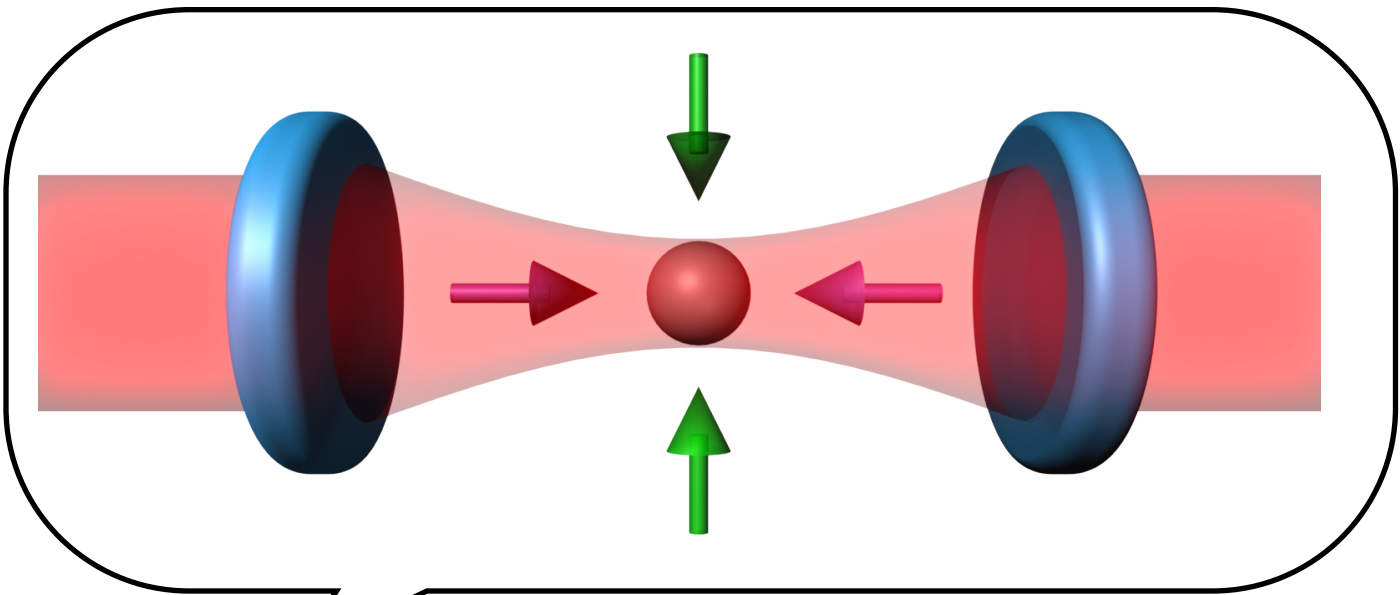


$$F \approx 15 \times 10^{-12} \text{ N}$$

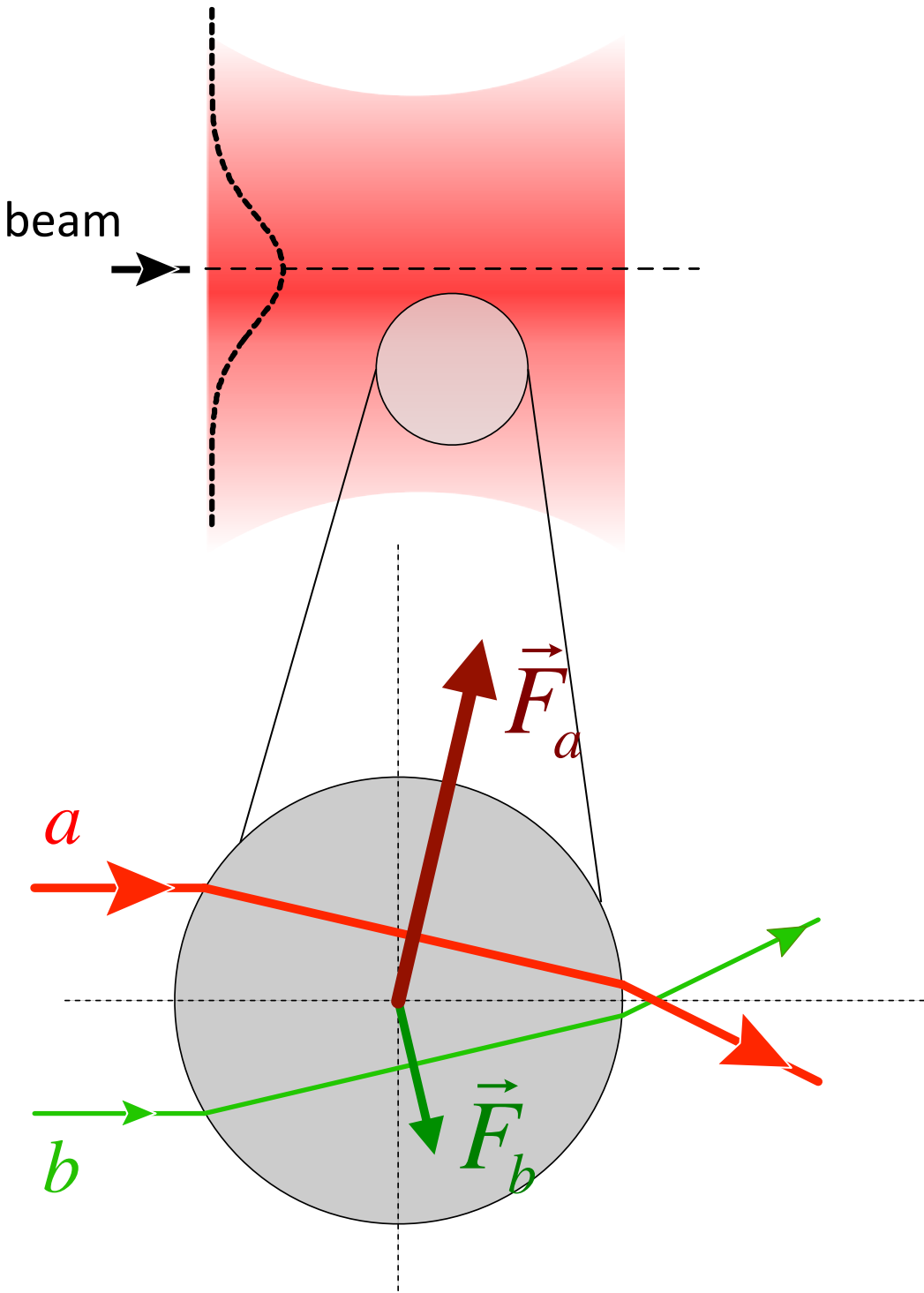
What can we do with this force?



Optical Forces in the Microscopic World



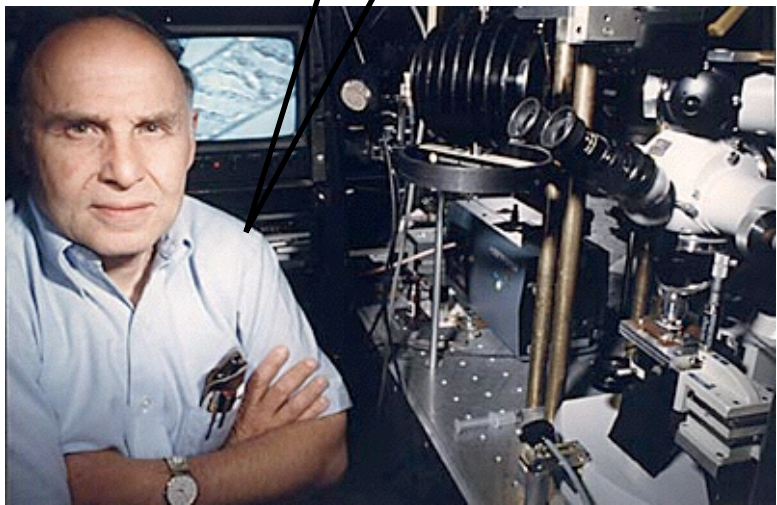
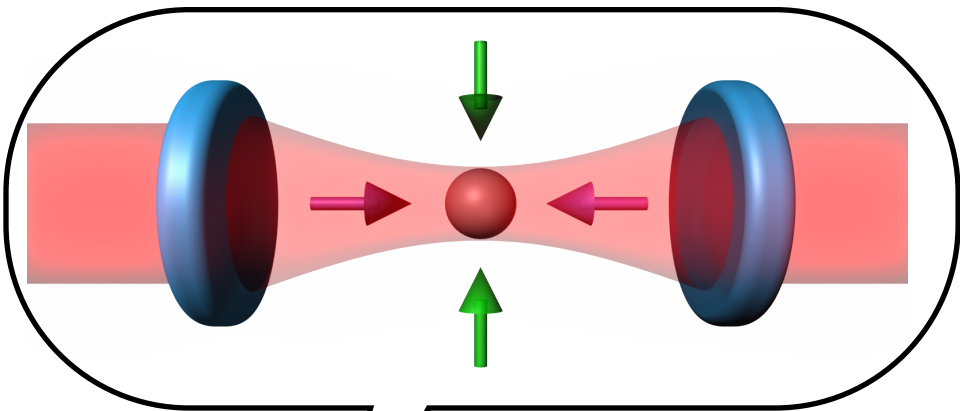
A. Ashkin et al. Science (1987)



A. Ashkin. IEEE JSTE, 6(6):841–856, 2000.
S. Chu, et al. PRL, 57(3):314–317, 1986.



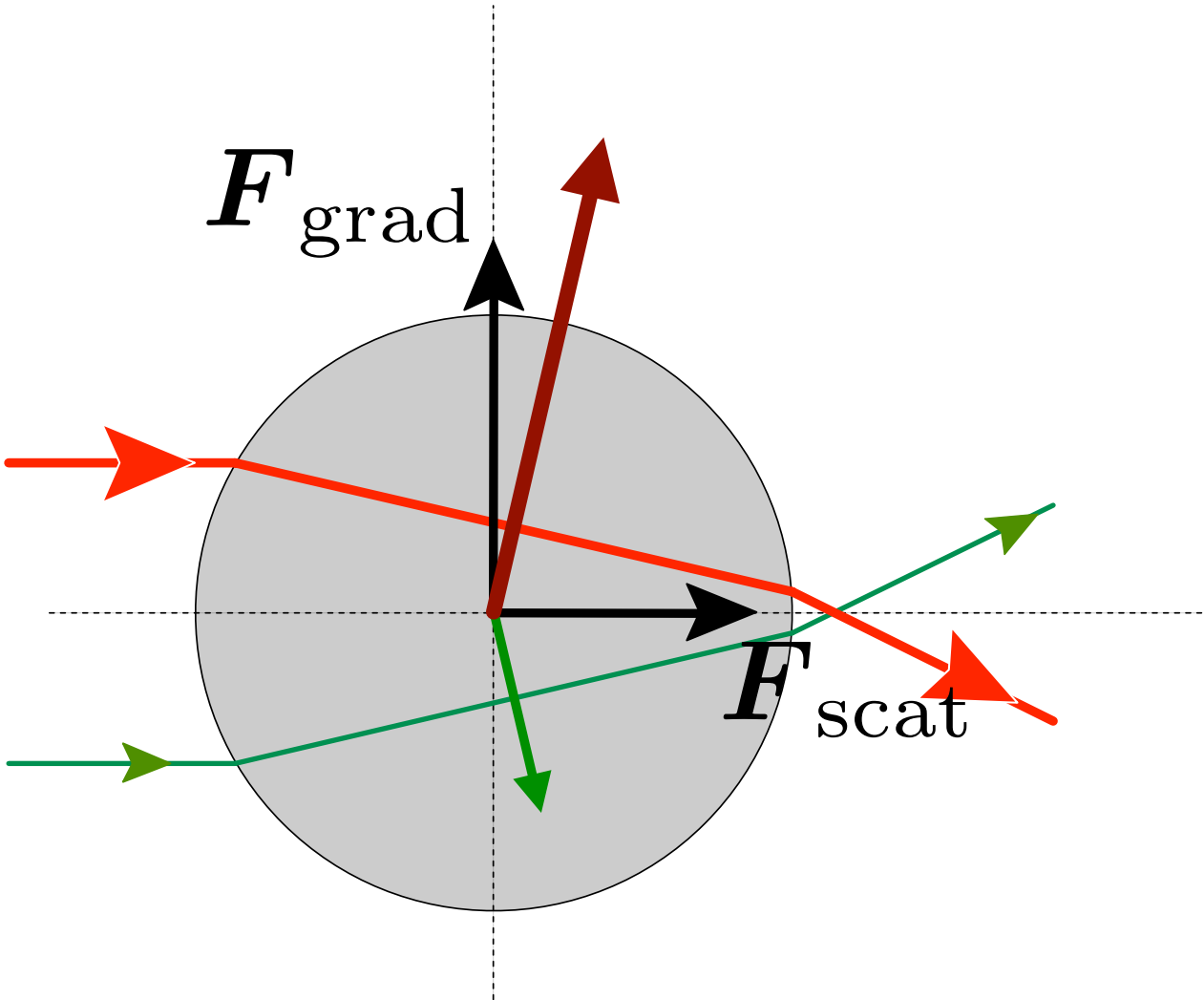
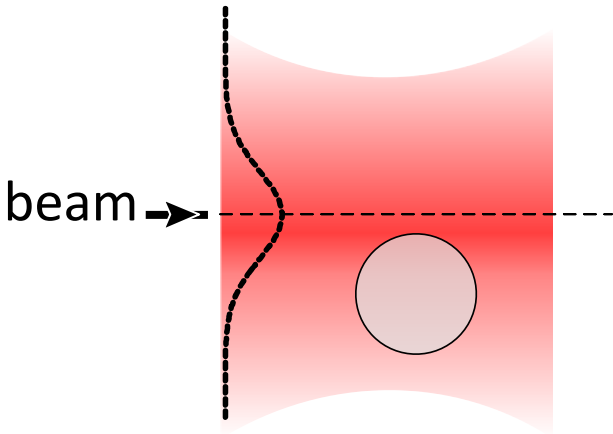
Optical Forces in the Microscopic World



A. Ashkin et al. Science (1987)

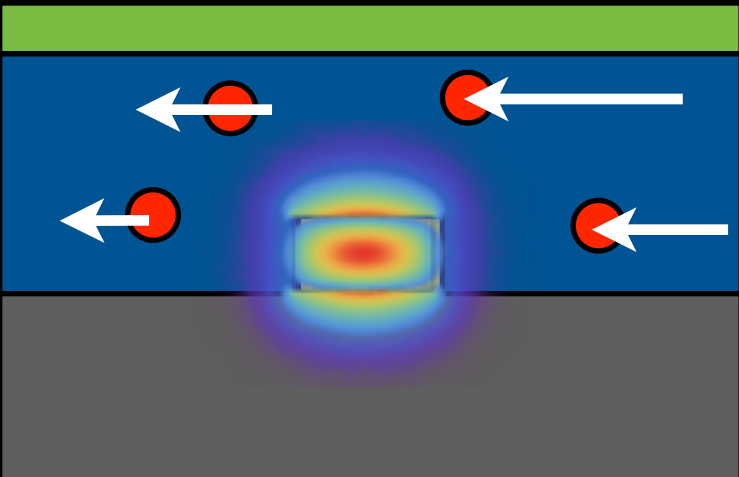
$$\vec{F} = \hbar n_{\text{photons}} \frac{\Delta \vec{k}}{\Delta t}$$

A. Ashkin. IEEE JSTE, 6(6):841–856, 2000.
S. Chu, et al. PRL, 57(3):314–317, 1986.

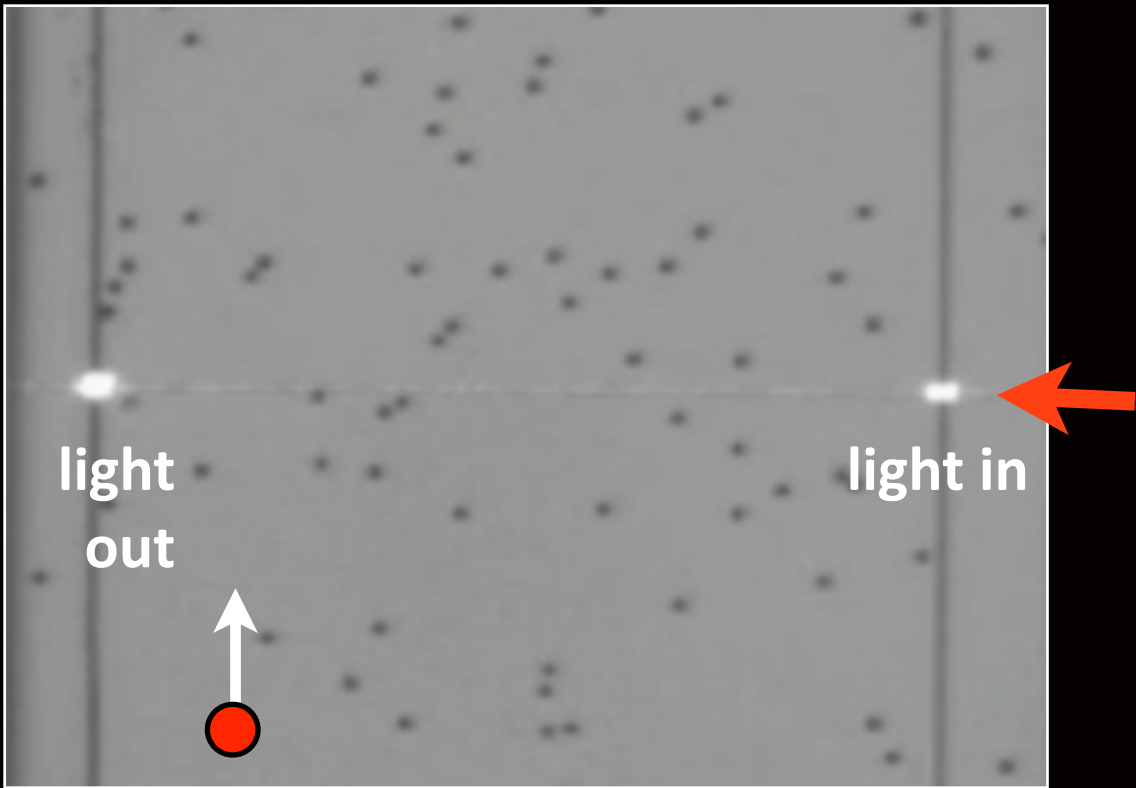




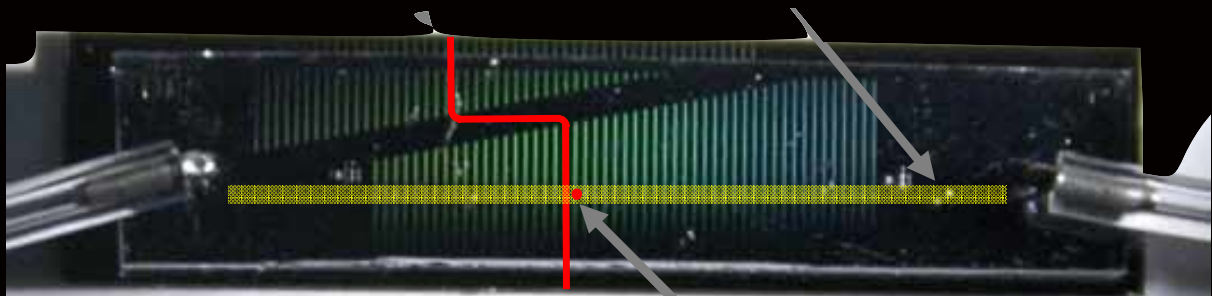
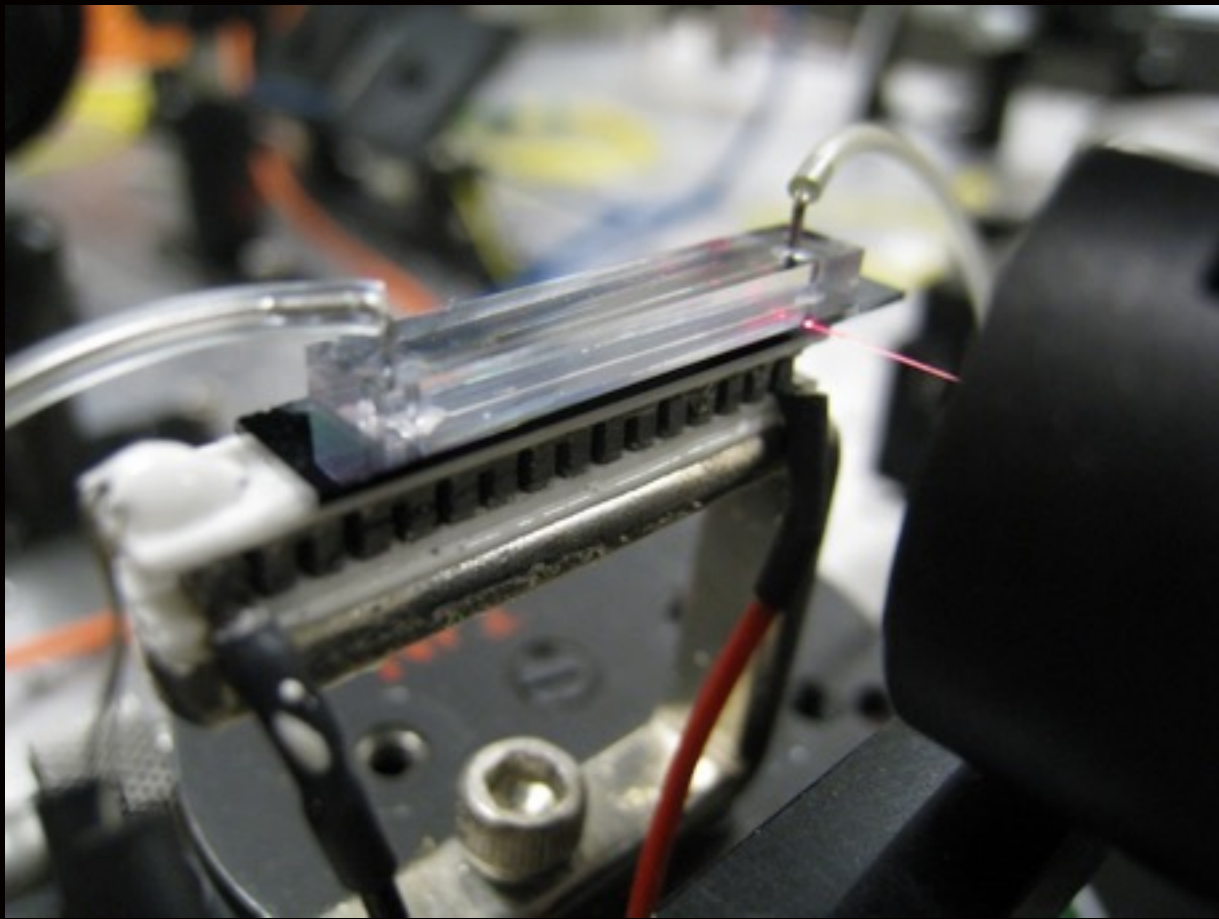
Optical Forces in the Microscopic World



waveguide cross-section



A. Nitkowski et al, Optics Letters 2009

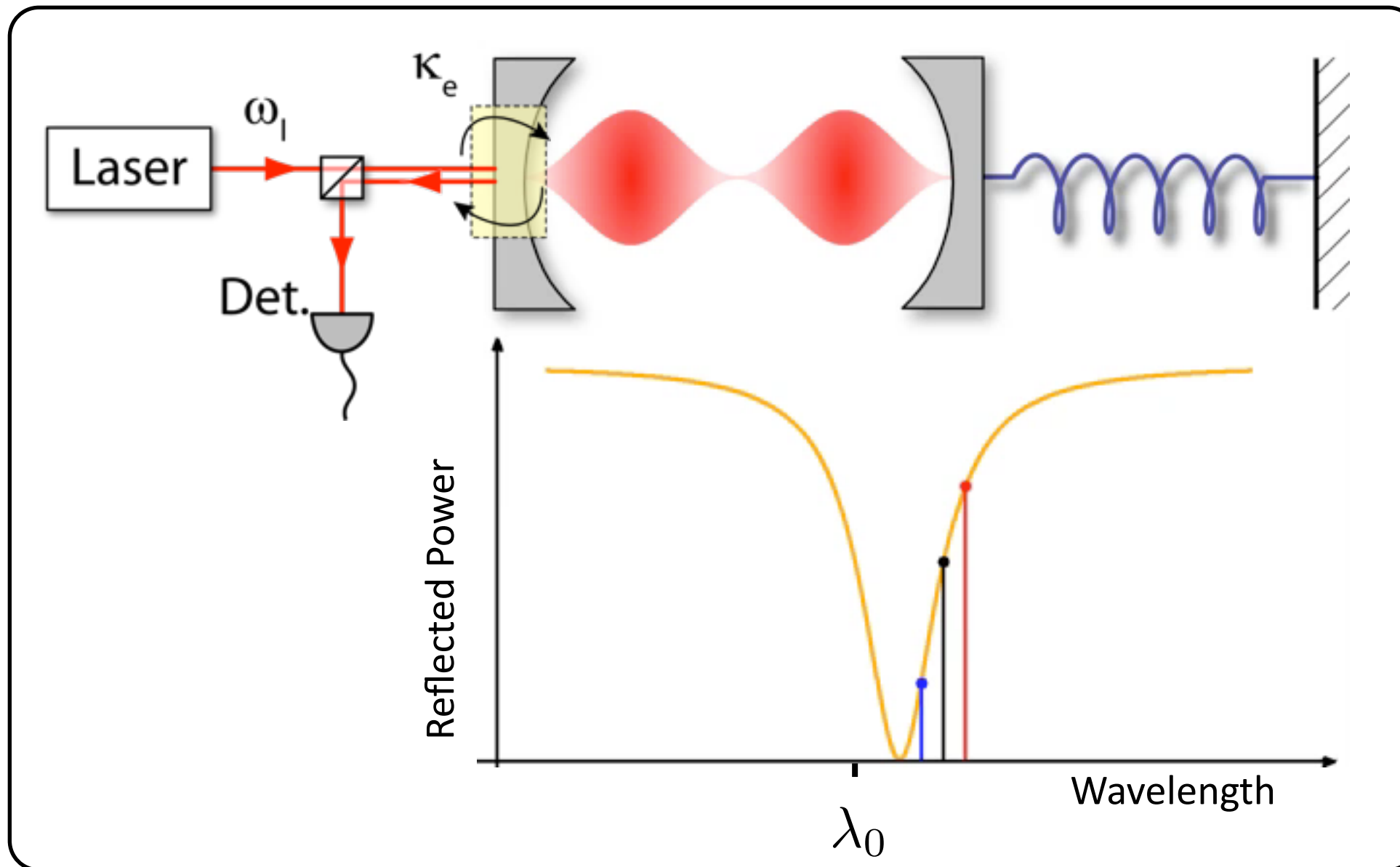


3 cm

Integrated particle trapping setup at Cornell
(Lipson group)



Mechanically Susceptible Cavity



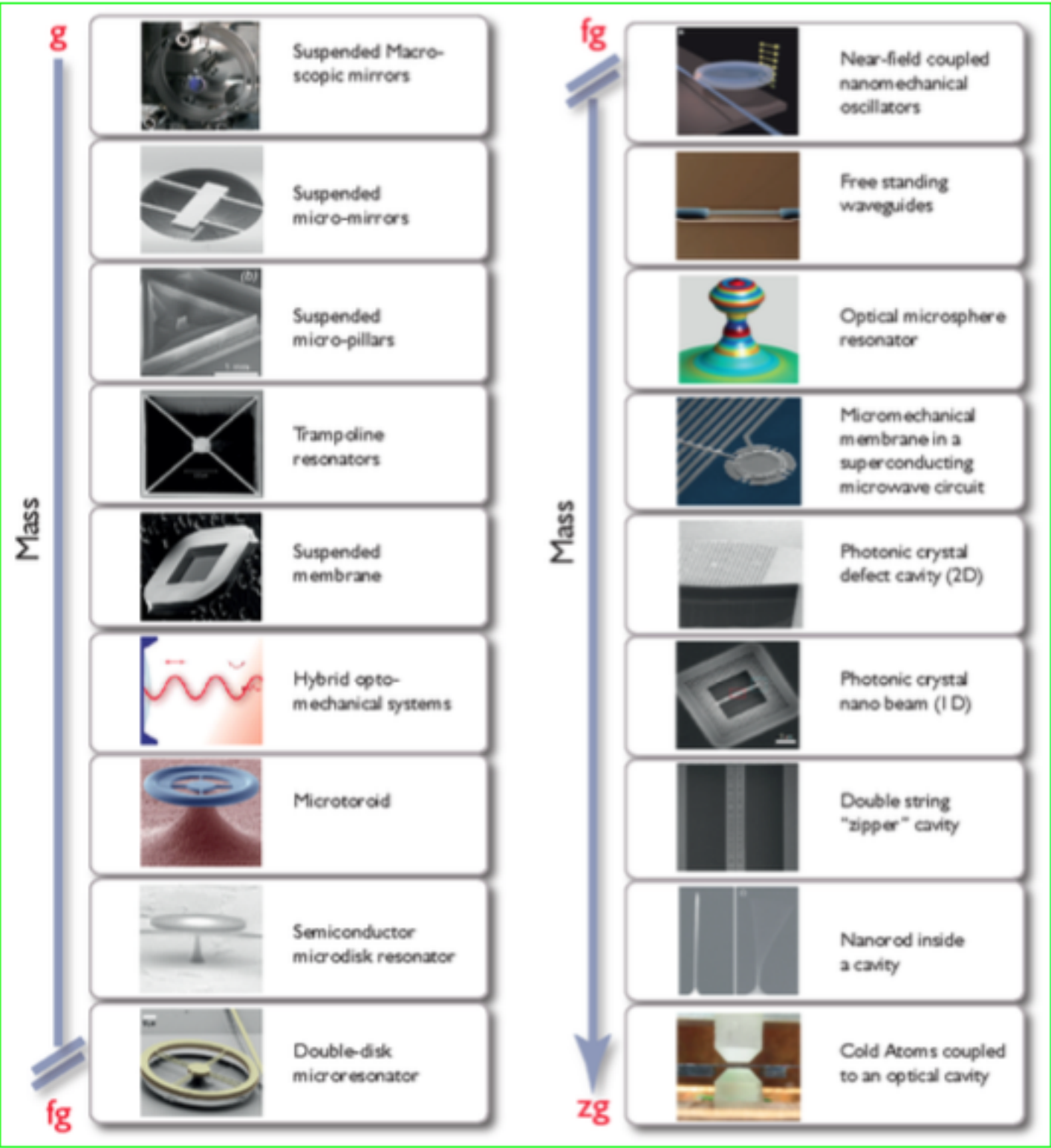
$$\begin{aligned}\phi_{\text{RT}} &= m(2\pi) \\ &= \frac{\omega_l}{c}(2L)\end{aligned}$$

$$\Rightarrow \omega_m = m \frac{\pi c}{L}$$

$$L = L_0 + x(t) \Rightarrow \omega_m(t) = m \frac{\pi c}{L_0 + x(t)} \approx \omega_m - \underbrace{\left(\frac{\omega_m}{L} \right)}_{g_{\text{OM}}} x(t)$$



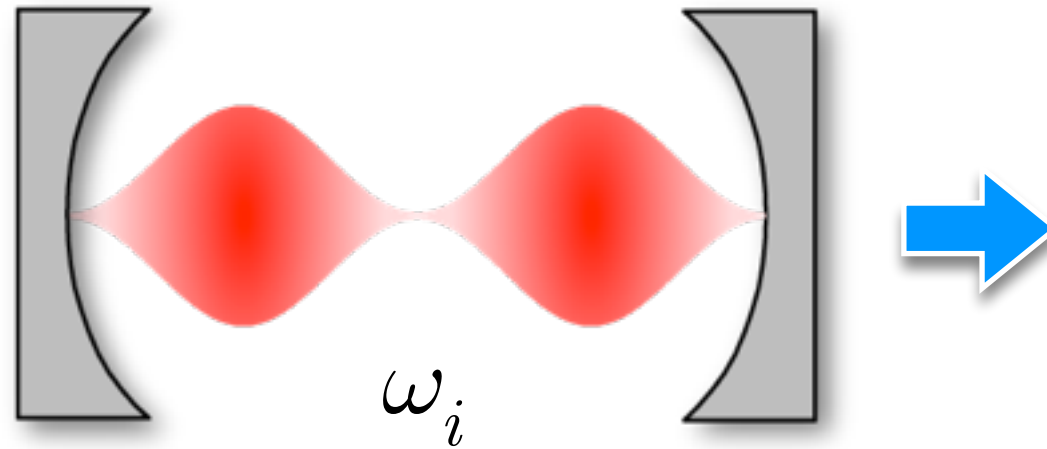
Mass-spring Fabry-Perot systems?





Modal field expansion

Maxwell's equations
+
boundary conditions



Optical Modes

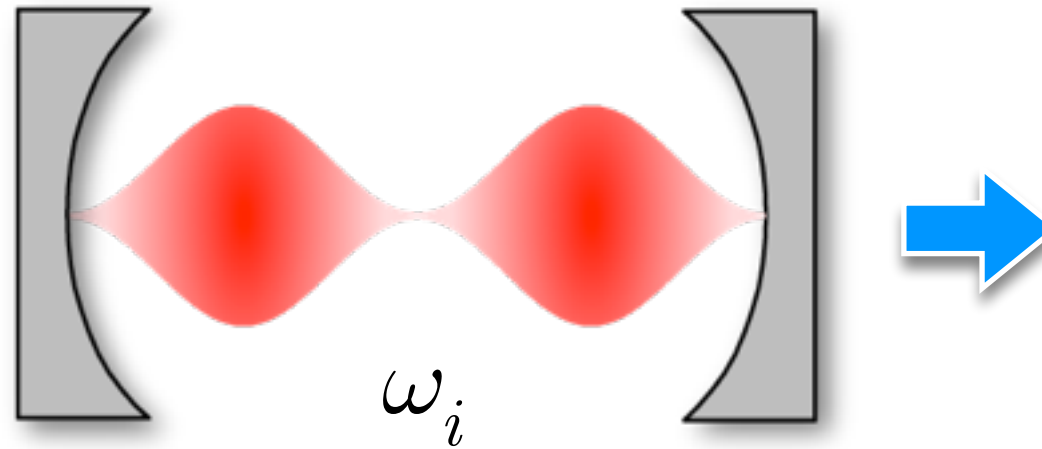
$$\mathbf{E}_j(\mathbf{r}, t) = \mathcal{E}_j(\mathbf{r})e^{-i\omega_j t}$$

$$\mathbf{B}_j(\mathbf{r}, t) = \mathcal{H}_j(\mathbf{r})e^{-i\omega_j t}$$



Modal field expansion

Maxwell's equations
+
boundary conditions



Optical Modes

$$\mathbf{E}_j(\mathbf{r}, t) = \mathcal{E}_j(\mathbf{r})e^{-i\omega_j t}$$

$$\mathbf{B}_j(\mathbf{r}, t) = \mathcal{H}_j(\mathbf{r})e^{-i\omega_j t}$$

Orthogonality relation

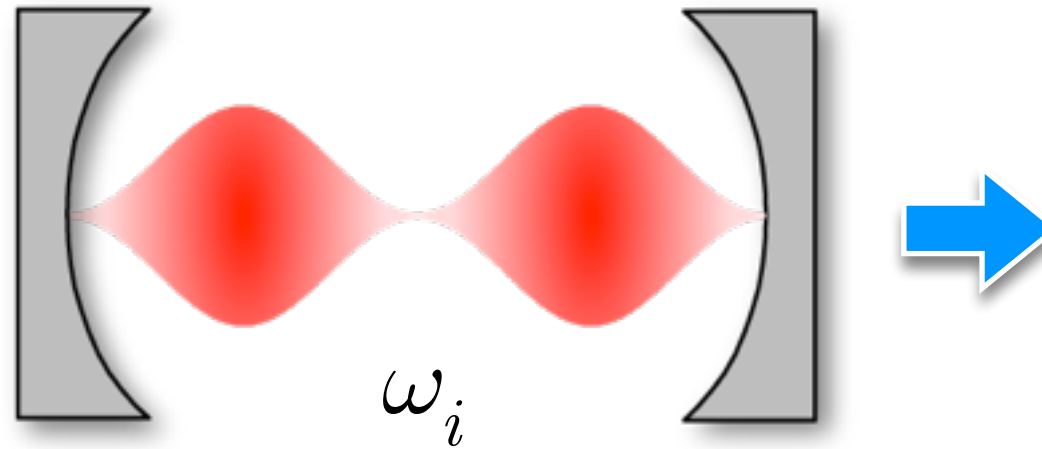
$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

$$\int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV = \delta_{ij}$$



Modal field expansion

Maxwell's equations
+
boundary conditions



Optical Modes

$$\mathbf{E}_j(\mathbf{r}, t) = \mathcal{E}_j(\mathbf{r})e^{-i\omega_j t}$$

$$\mathbf{B}_j(\mathbf{r}, t) = \mathcal{H}_j(\mathbf{r})e^{-i\omega_j t}$$

Orthogonality relation

$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

$$\int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV = \delta_{ij}$$

Spatial solution

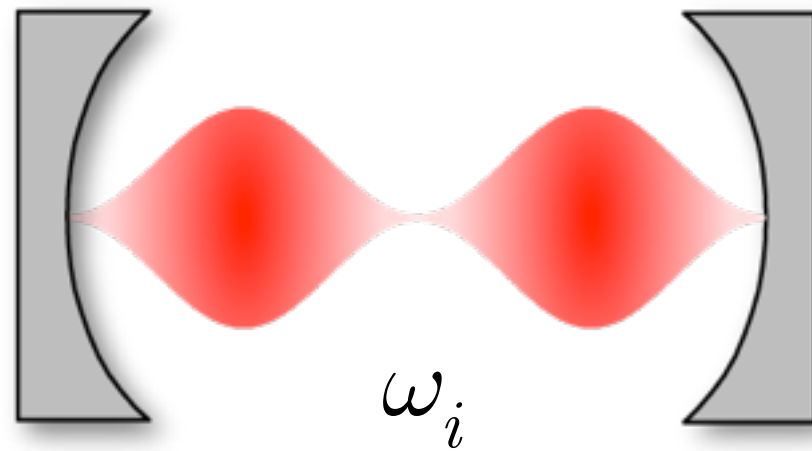
$$\nabla \times \mathcal{H}_j = -i\omega_j \epsilon_0 \mathcal{E}_j$$

$$\nabla \times \mathcal{E}_j = i\mu_0 \omega_j \mathcal{H}_j$$



Modal field expansion

Maxwell's equations
+
boundary conditions



Optical Modes

$$\mathbf{E}_j(\mathbf{r}, t) = \mathcal{E}_j(\mathbf{r})e^{-i\omega_j t}$$

$$\mathbf{B}_j(\mathbf{r}, t) = \mathcal{H}_j(\mathbf{r})e^{-i\omega_j t}$$

Orthogonality relation

$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

$$\int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV = \delta_{ij}$$

Expand total optical field as
linear superposition of the
optical modes

Spatial solution

$$\nabla \times \mathcal{H}_j = -i\omega_j \epsilon_0 \mathcal{E}_j$$

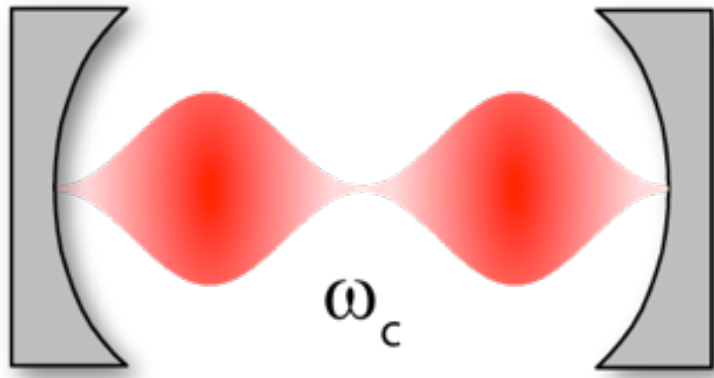
$$\nabla \times \mathcal{E}_j = i\mu_0 \omega_j \mathcal{H}_j$$

$$\mathbf{E}(\mathbf{r}, t) = \sum a_j(t) \mathcal{E}_j(\mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \sum a_j(t) \mathcal{H}_j(\mathbf{r})$$



Derivation of field amplitude lumped model



$$\mathbf{E}(\mathbf{r}, t) = \sum a_j(t) \mathcal{E}_j(\mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \sum a_j(t) \mathcal{H}_j(\mathbf{r})$$

Maxwell's equations

$$\nabla \times \mathbf{E}(\mathbf{r}, t) = -\mu_0 \partial_t \mathbf{H}(\mathbf{r}, t)$$

$$\nabla \times \mathbf{H}(\mathbf{r}, t) = \epsilon \partial_t \mathbf{E}(\mathbf{r}, t)$$

Orthogonality relation

$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

$$\int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV = \delta_{ij}$$

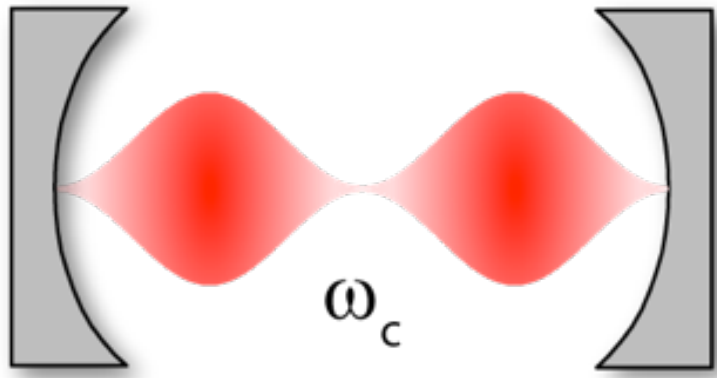
Spatial solution

$$\nabla \times \mathcal{H}_j = -i\omega_j \epsilon_0 \mathcal{E}_j$$

$$\nabla \times \mathcal{E}_j = i\mu_0 \omega_j \mathcal{H}_j$$



Derivation of field amplitude lumped model



$$\mathbf{E}(\mathbf{r}, t) = \sum a_j(t) \mathcal{E}_j(\mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \sum a_j(t) \mathcal{H}_j(\mathbf{r})$$

Orthogonality relation

$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

$$\int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV = \delta_{ij}$$

Spatial solution

$$\nabla \times \mathcal{H}_j = -i\omega_j \epsilon_0 \mathcal{E}_j$$

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Maxwell's equations

$$\nabla \times \mathbf{E}(\mathbf{r}, t) = -\mu_0 \partial_t \mathbf{H}(\mathbf{r}, t)$$

$$\nabla \times \mathbf{H}(\mathbf{r}, t) = \epsilon \partial_t \mathbf{E}(\mathbf{r}, t)$$



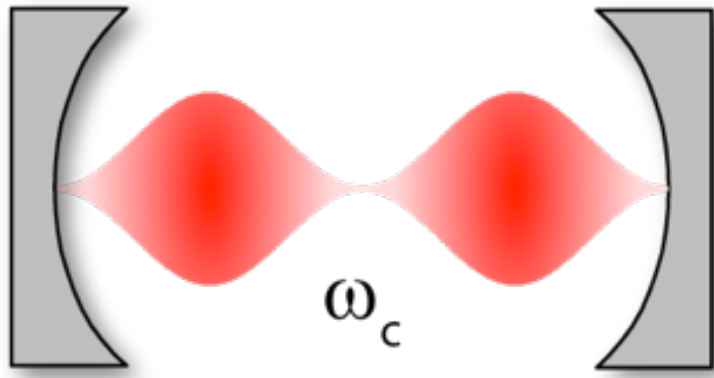
To account for cavity loss we assume that:

$$\epsilon = \epsilon_r + i\epsilon_i$$

$$\epsilon_r \gg \epsilon_i \rightarrow \epsilon \approx \epsilon_r$$



Derivation of field amplitude lumped model



$$\mathbf{E}(\mathbf{r}, t) = \sum a_j(t) \mathcal{E}_j(\mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \sum a_j(t) \mathcal{H}_j(\mathbf{r})$$

Orthogonality relation

$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

$$\int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV = \delta_{ij}$$

Spatial solution

$$\nabla \times \mathcal{H}_j = -i\omega_j \epsilon_0 \mathcal{E}_j$$

$$\nabla \times \mathcal{E}_j = i\mu_0 \omega_j \mathcal{H}_j$$

Maxwell's equations

$$\nabla \times \mathbf{E}(\mathbf{r}, t) = -\mu_0 \partial_t \mathbf{H}(\mathbf{r}, t)$$

$$\nabla \times \mathbf{H}(\mathbf{r}, t) = \epsilon \partial_t \mathbf{E}(\mathbf{r}, t)$$



To account for cavity
loss we assume that:

$$\epsilon = \epsilon_r + i\epsilon_i$$

$$\epsilon_r \gg \epsilon_i \rightarrow \epsilon \approx \epsilon_r$$

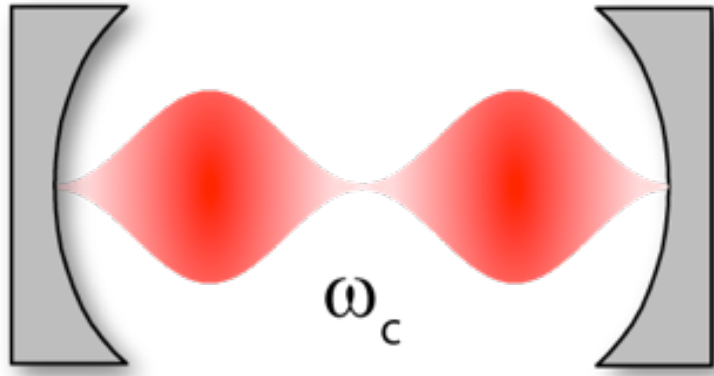
Therefore:

$$\nabla \times \left(\sum a_j(t) \mathcal{H}_j(\mathbf{r}) \right) = (\epsilon_r + i\epsilon_i) \partial_t \left(\sum a_j(t) \mathcal{E}_j(\mathbf{r}) \right)$$

$$\sum a_j(t) (\nabla \times \mathcal{H}_j(\mathbf{r})) = \sum (\epsilon_r + i\epsilon_i) \mathcal{E}_j(\mathbf{r}) \dot{a}_j(t)$$



Derivation of field amplitude lumped model



$$\mathbf{E}(\mathbf{r}, t) = \sum a_j(t) \mathcal{E}_j(\mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \sum a_j(t) \mathcal{H}_j(\mathbf{r})$$

Therefore:

$$\sum a_j(t) (\nabla \times \mathcal{H}_j(\mathbf{r})) = \sum (\epsilon_r + i\epsilon_i) \mathcal{E}_j(\mathbf{r}) \dot{a}_j(t)$$

Orthogonality relation

$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

$$\int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV = \delta_{ij}$$

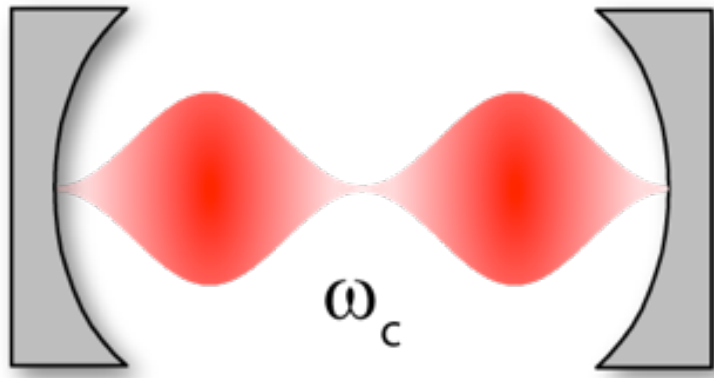
Spatial solution

$$\nabla \times \mathcal{H}_j = -i\omega_j \epsilon_0 \mathcal{E}_j$$

$$\nabla \times \mathcal{E}_j = i\mu_0 \omega_j \mathcal{H}_j$$



Derivation of field amplitude lumped model



$$\begin{aligned}\mathbf{E}(\mathbf{r}, t) &= \sum a_j(t) \mathcal{E}_j(\mathbf{r}) \\ \mathbf{H}(\mathbf{r}, t) &= \sum a_j(t) \mathcal{H}_j(\mathbf{r})\end{aligned}$$

Therefore:

$$\sum a_j(t) (\nabla \times \mathcal{H}_j(\mathbf{r})) = \sum (\epsilon_r + i\epsilon_i) \mathcal{E}_j(\mathbf{r}) \dot{a}_j(t)$$

$$\int \left[\sum \epsilon \mathcal{E}_j(\mathbf{r}) \left[\left(1 + i \frac{\epsilon_i}{\epsilon} \right) \dot{a}_j(t) + i\omega_j a_j(t) \right] = 0 \right] \mathcal{E}_n^*(\mathbf{r}) dV$$

Orthogonality relation

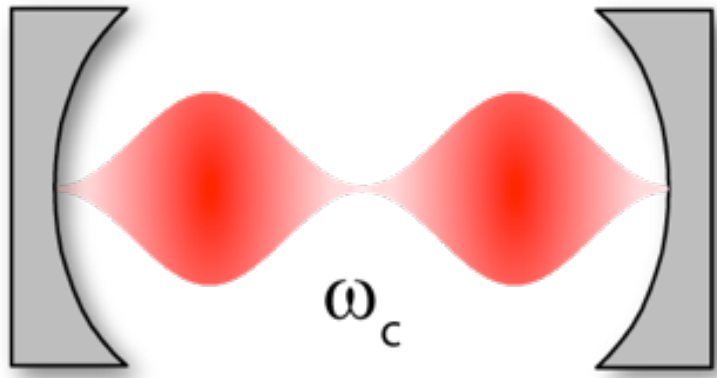
$$\begin{aligned}\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV &= \delta_{ij} \\ \int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV &= \delta_{ij}\end{aligned}$$

Spatial solution

$$\begin{aligned}\nabla \times \mathcal{H}_j &= -i\omega_j \epsilon_0 \mathcal{E}_j \\ \nabla \times \mathcal{E}_j &= i\mu_0 \omega_j \mathcal{H}_j\end{aligned}$$



Derivation of field amplitude lumped model



$$\begin{aligned}\mathbf{E}(\mathbf{r}, t) &= \sum a_j(t) \mathcal{E}_j(\mathbf{r}) \\ \mathbf{H}(\mathbf{r}, t) &= \sum a_j(t) \mathcal{H}_j(\mathbf{r})\end{aligned}$$

Therefore:

$$\sum a_j(t) (\nabla \times \mathcal{H}_j(\mathbf{r})) = \sum (\epsilon_r + i\epsilon_i) \mathcal{E}_j(\mathbf{r}) \dot{a}_j(t)$$

$$\int \left[\sum \epsilon \mathcal{E}_j(\mathbf{r}) \left[\left(1 + i \frac{\epsilon_i}{\epsilon} \right) \dot{a}_j(t) + i\omega_j a_j(t) \right] = 0 \right] \mathcal{E}_n^*(\mathbf{r}) dV$$

Orthogonality relation

$$\begin{aligned}\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV &= \delta_{ij} \\ \int \epsilon (\mathcal{E}_i^* \cdot \mathcal{E}_j) dV &= \delta_{ij}\end{aligned}$$

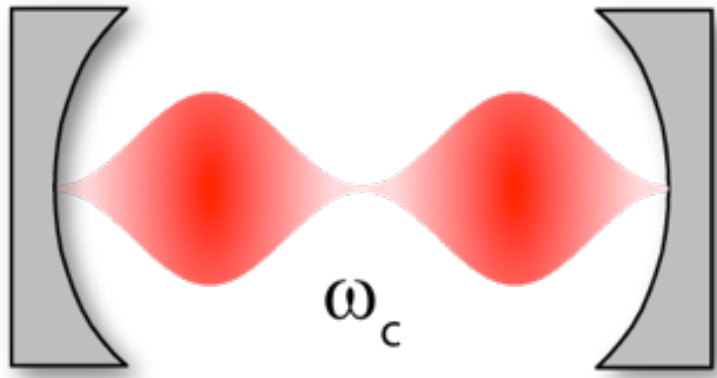
$$\dot{a}_n(t) = -i\omega_n a_n(t) - \left(\omega_n \frac{\epsilon_i}{\epsilon} \right) a_n(t)$$

Spatial solution

$$\begin{aligned}\nabla \times \mathcal{H}_j &= -i\omega_j \epsilon_0 \mathcal{E}_j \\ \nabla \times \mathcal{E}_j &= i\mu_0 \omega_j \mathcal{H}_j\end{aligned}$$



Derivation of field amplitude lumped model



$$\mathbf{E}(\mathbf{r}, t) = \sum a_j(t) \mathcal{E}_j(\mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \sum a_j(t) \mathcal{H}_j(\mathbf{r})$$

Therefore:

$$\mathbf{E}(\mathbf{r}, t) = \sum a_j(t) \mathcal{E}_j(\mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \sum a_j(t) \mathcal{H}_j(\mathbf{r})$$

Orthogonality relation

$$\int \frac{1}{\mu} (\mathcal{H}_i^* \cdot \mathcal{H}_j) dV = \delta_{ij}$$

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$$\dot{a}_n(t) = -i\omega_n a_n(t) - \left(\omega_n \frac{\epsilon_i}{\epsilon} \right) a_n(t)$$

Spatial solution

$$\nabla \times \mathcal{H}_j = -i\omega_j \epsilon_0 \mathcal{E}_j$$

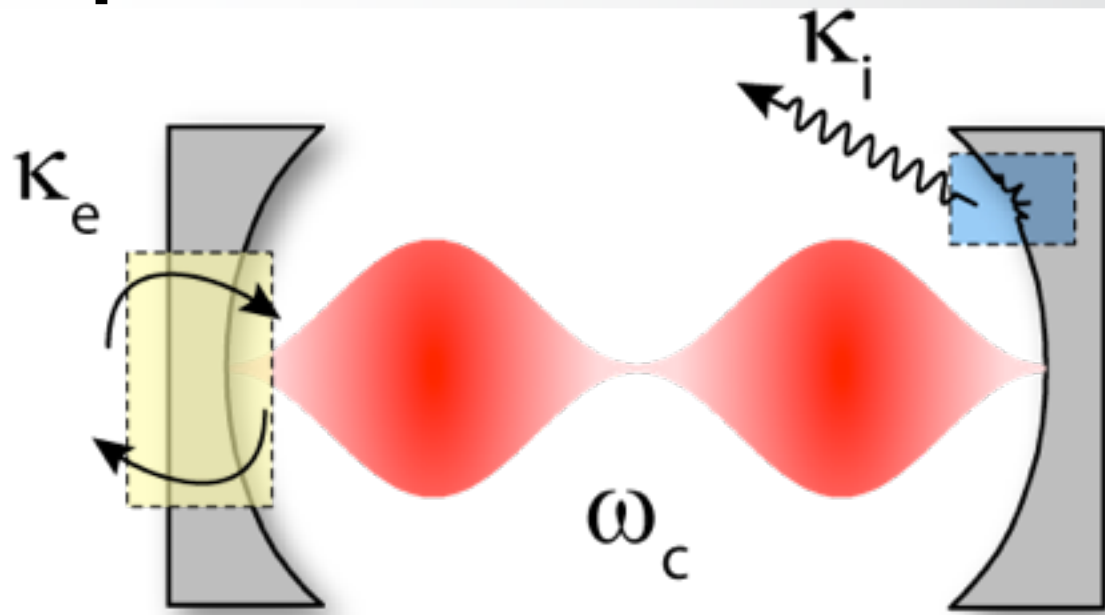
$$\nabla \times \mathcal{E}_j = i\mu_0 \omega_j \mathcal{H}_j$$

Lumped Model

$$\dot{a}(t) = -i\omega_c a(t) - \frac{\kappa}{2} a(t)$$



Equation for the field amplitude

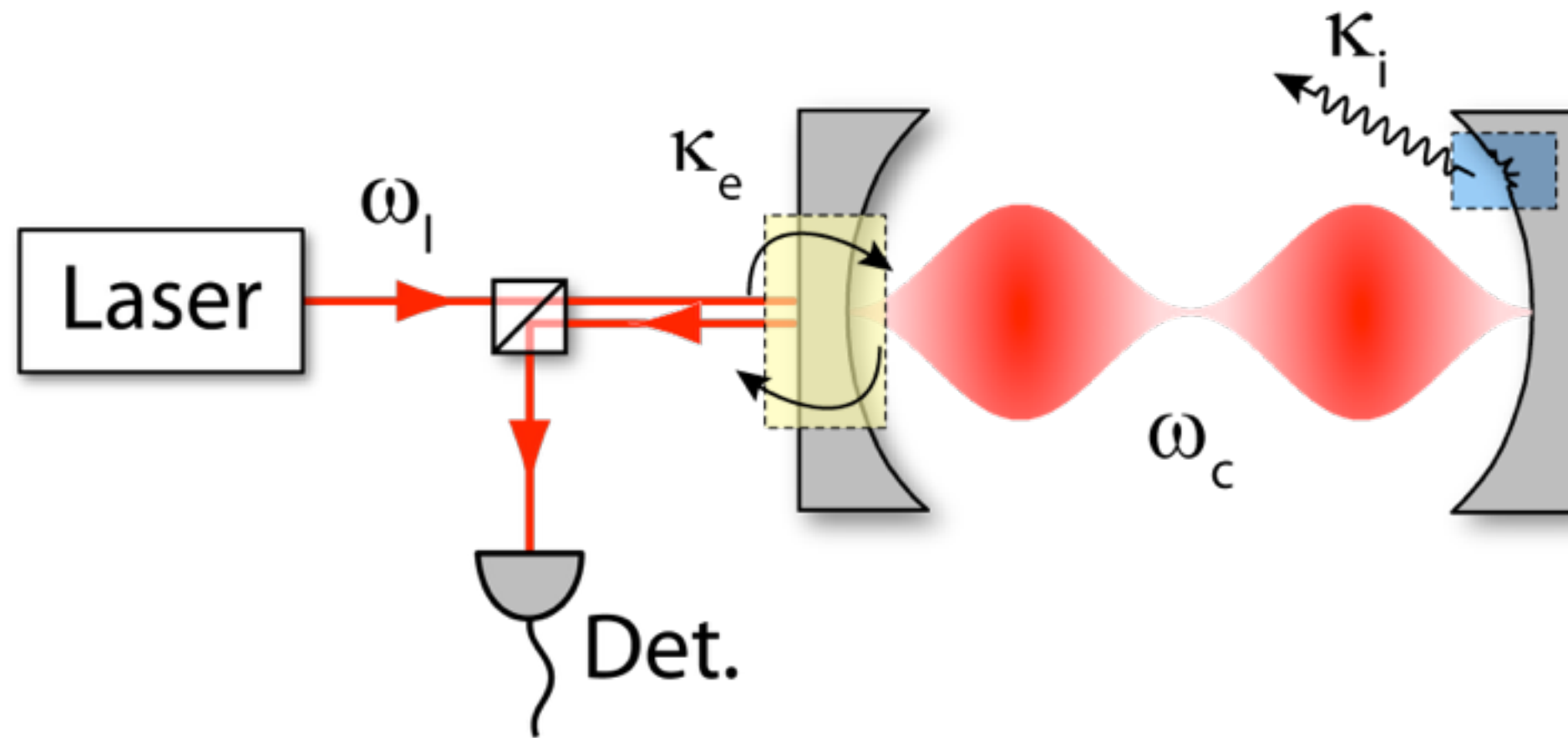


$$\kappa = \kappa_e + \kappa_i$$

$$\dot{a}(t) = -i\omega_c a(t) - \frac{\kappa}{2} a(t) + \sqrt{\kappa_e} \alpha_{\text{in}} e^{-i\omega_l t}$$



Optical Time Domain Equation

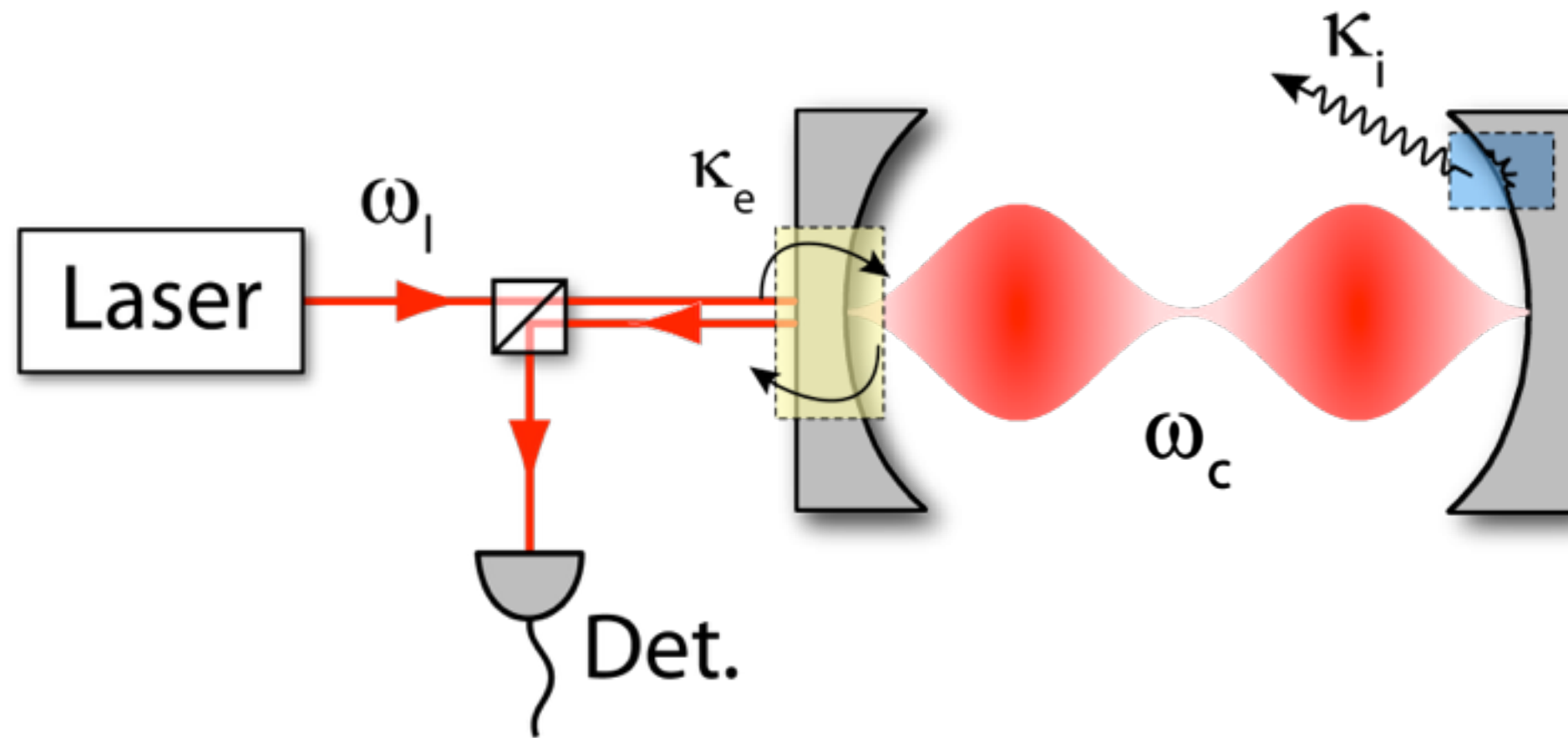


$$\dot{a}(t) = -i\omega_c a(t) - \frac{\kappa}{2} a(t) + \sqrt{\kappa_e} \alpha_{\text{in}} e^{-i\omega_l t}$$

$\alpha_{\text{in}} = \text{Optical Pump Field Rate}$



Optical Time Domain Equation



$$\dot{a}(t) = -i\omega_c a(t) - \frac{\kappa}{2} a(t) + \sqrt{\kappa_e} \alpha_{\text{in}} e^{-i\omega_l t}$$

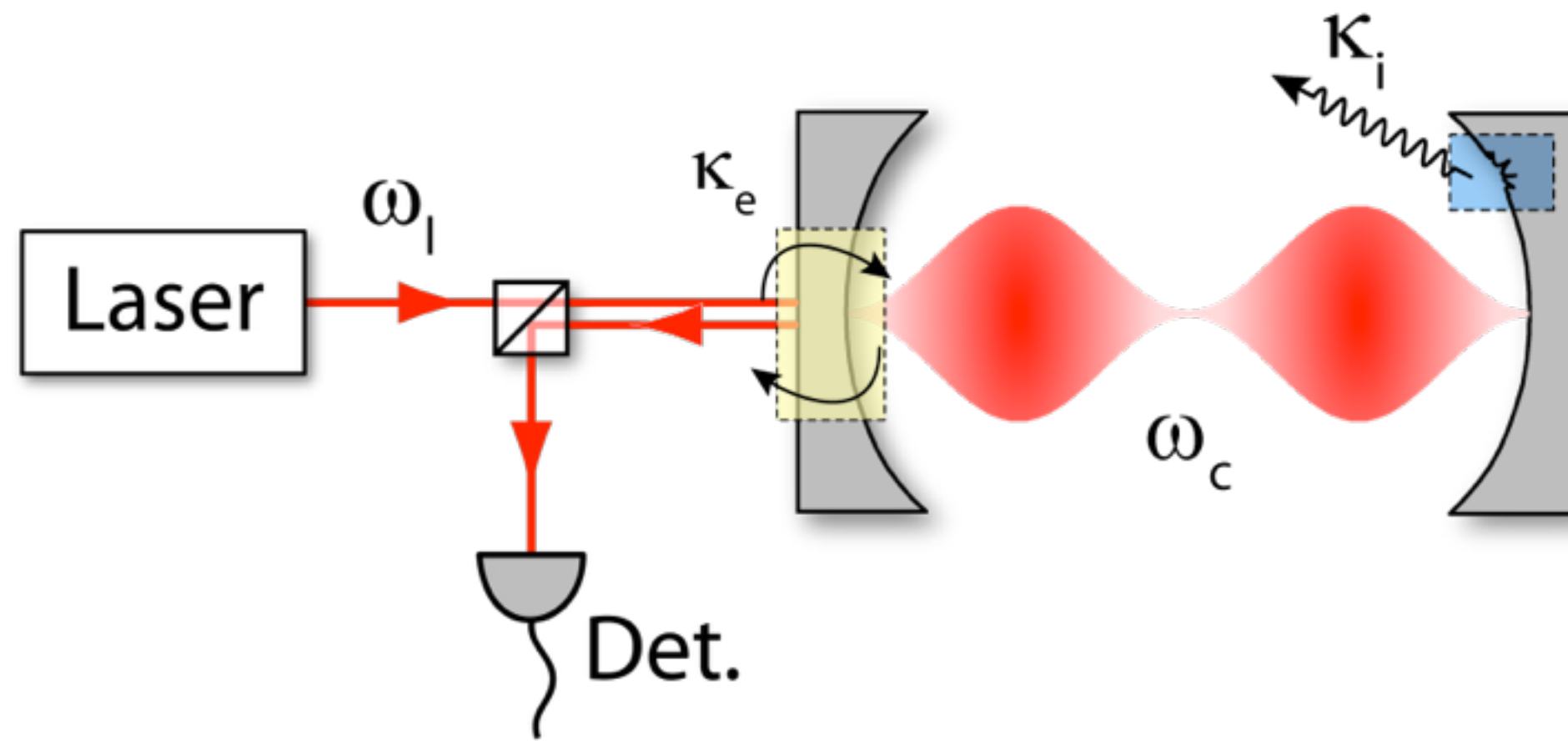
Normalization

α_{in} = Optical Pump Field Rate

$$\hbar\omega_l |\alpha_{\text{in}}|^2 = P_{\text{in}}$$



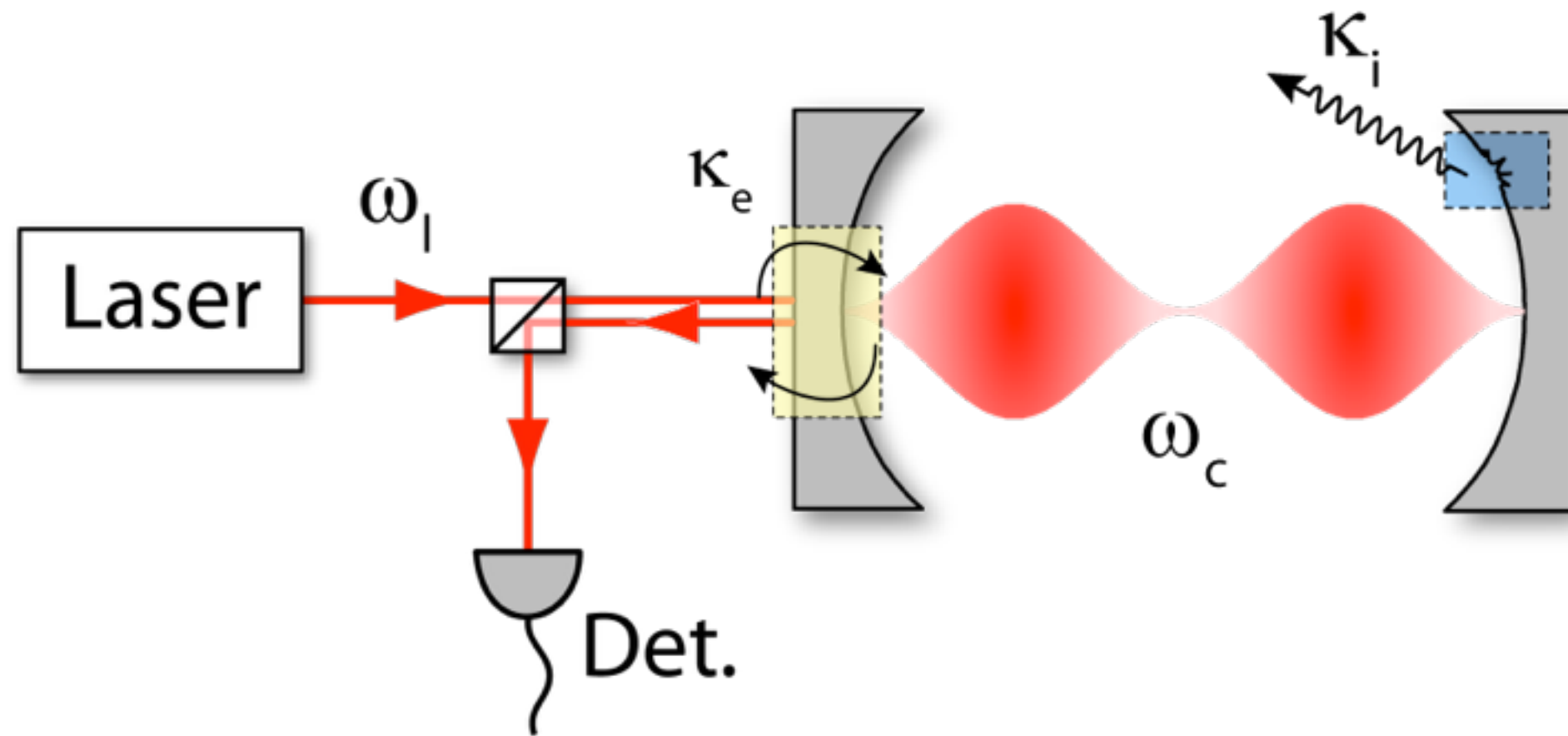
Rotating wave approximation



$$\dot{a}(t) = -i\omega_c a(t) - \frac{\kappa}{2} a(t) + \sqrt{\kappa_e} \alpha_{\text{in}} e^{-i\omega_l t}$$



Rotating wave approximation

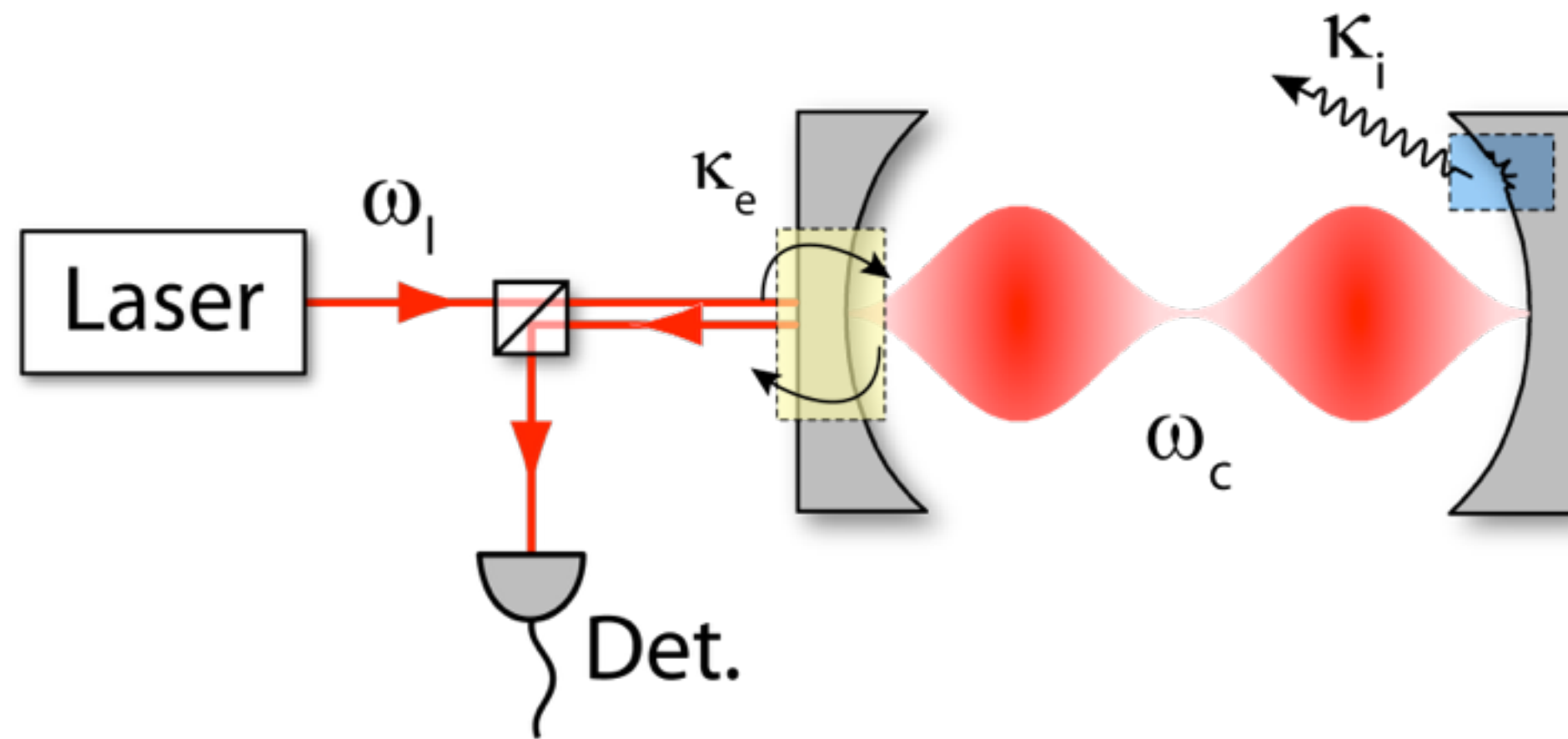


$$\dot{a}(t) = -i\omega_c a(t) - \frac{\kappa}{2} a(t) + \sqrt{\kappa_e} \alpha_{\text{in}} e^{-i\omega_l t}$$

$$a_{\text{old}}(t) = a(t) e^{-i\omega_l t}$$



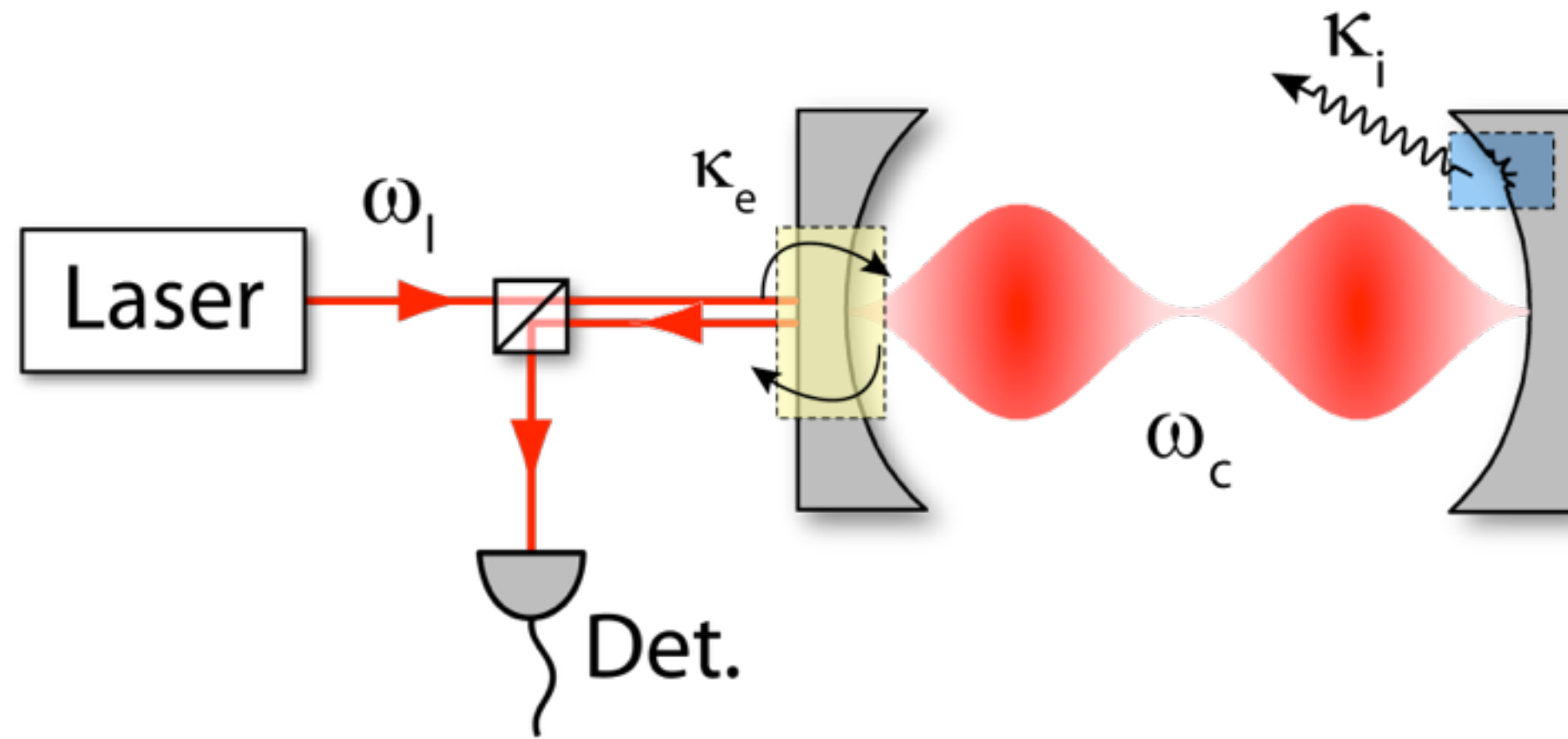
Rotating wave approximation



$$\dot{a}(t) = \boxed{-i\omega_c} a(t) - \frac{\kappa}{2} a(t) + \sqrt{\kappa_e} \alpha_{\text{in}} \boxed{e^{-i\omega_l t}}$$



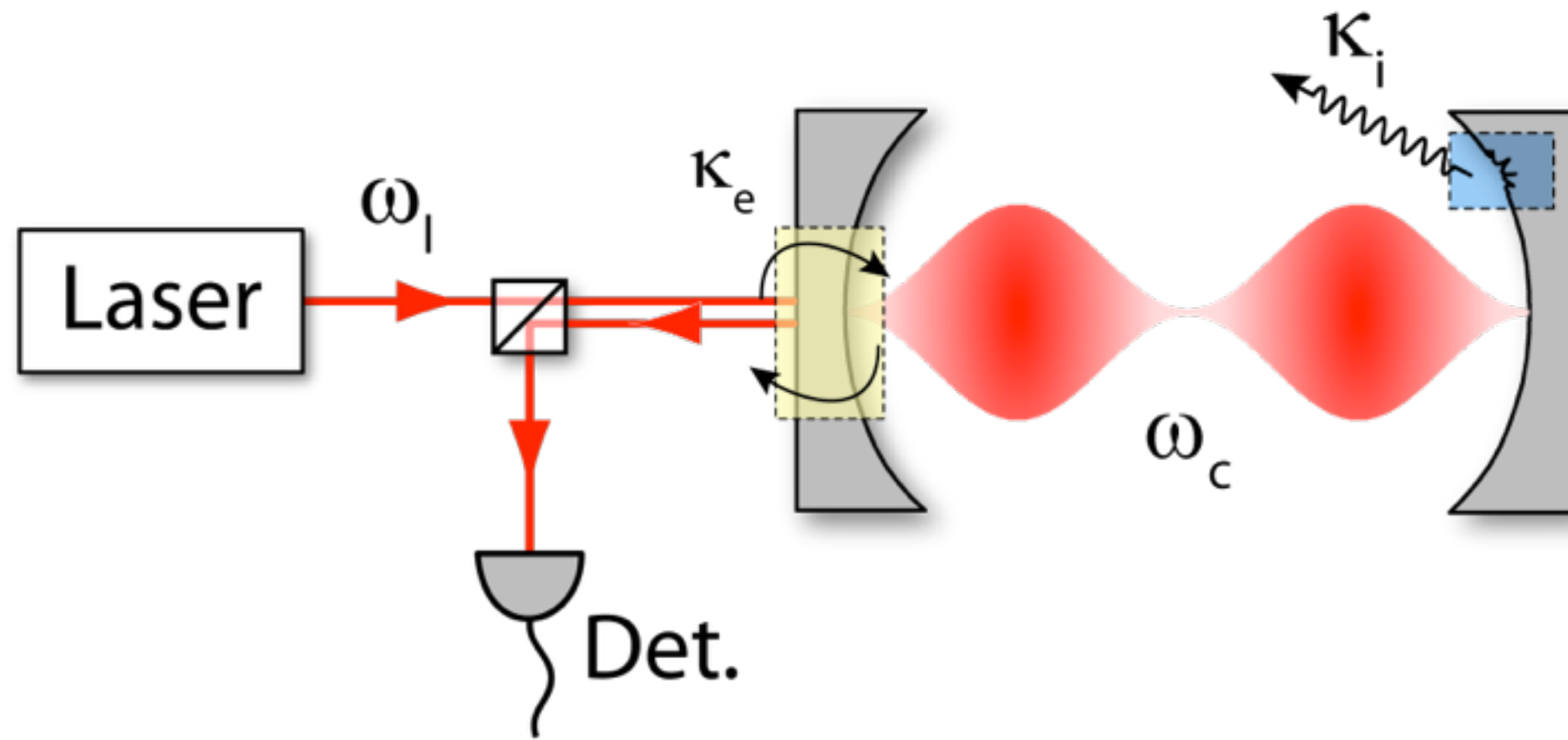
Optical Time Domain Equation



$$\dot{a}(t) = i(\omega_l - \omega_c)a - \frac{\kappa}{2}a + \sqrt{\kappa_e}a_{\text{in}}$$



Optical Time Domain Equation

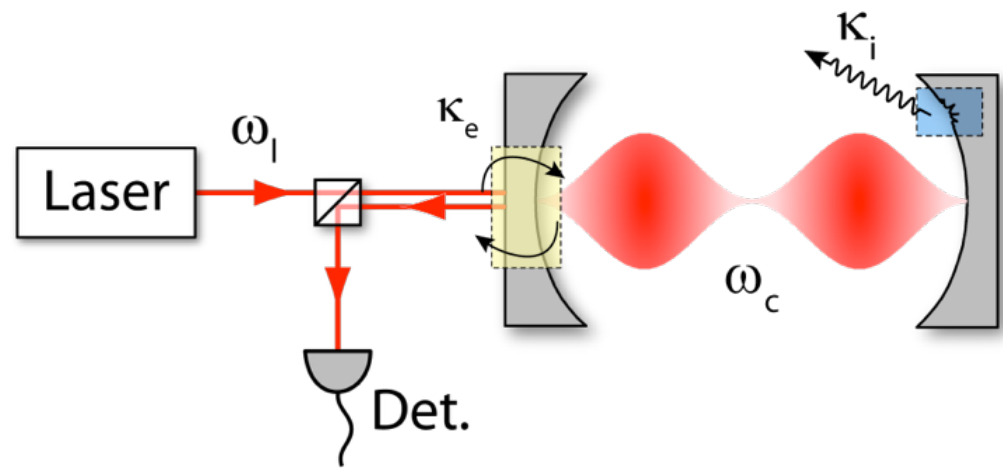


$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2}a + \sqrt{\kappa_e}a_{\text{in}}$$

Optical Amplitude Equation



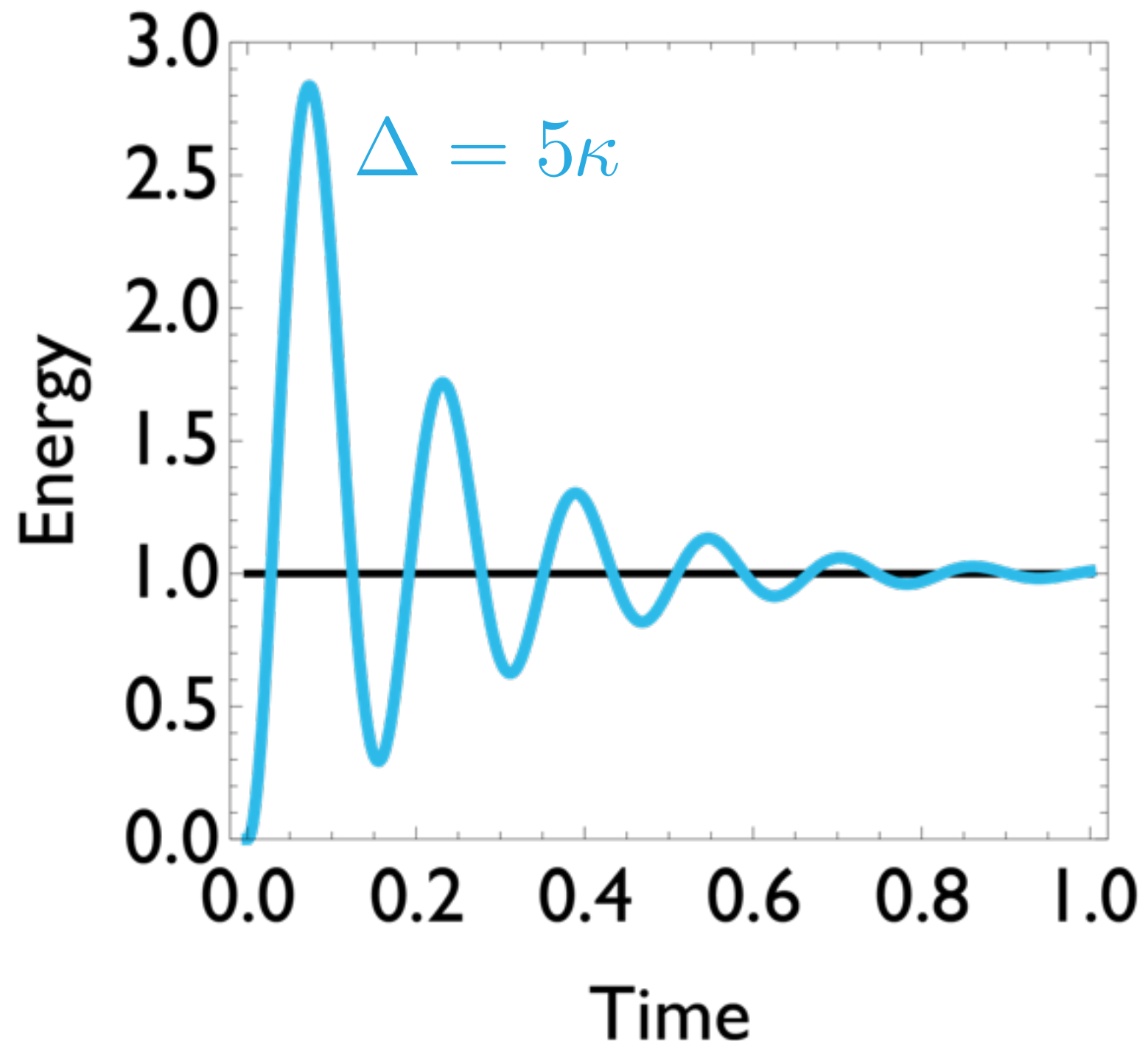
Optical Time Domain Equation



Solution

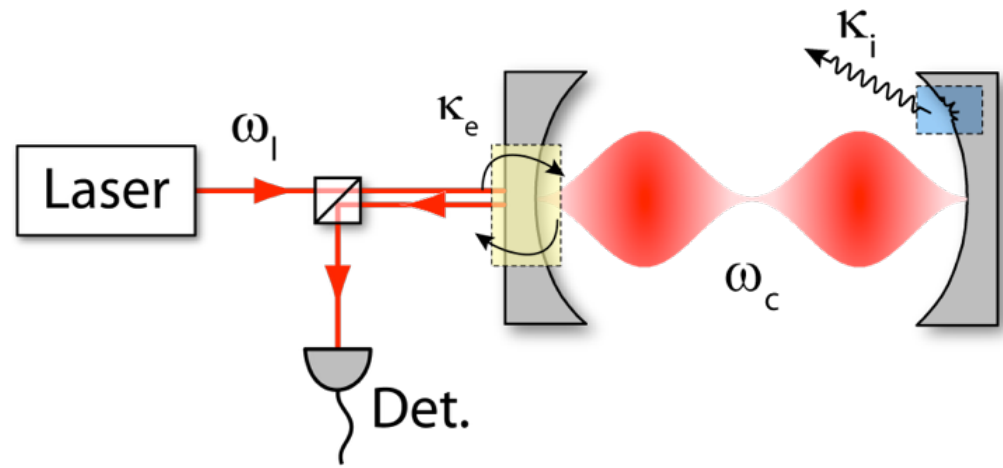
$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2}a + \sqrt{\kappa_e}\alpha_{\text{in}}$$

$$|a(t)|^2 \propto \text{Energy}$$





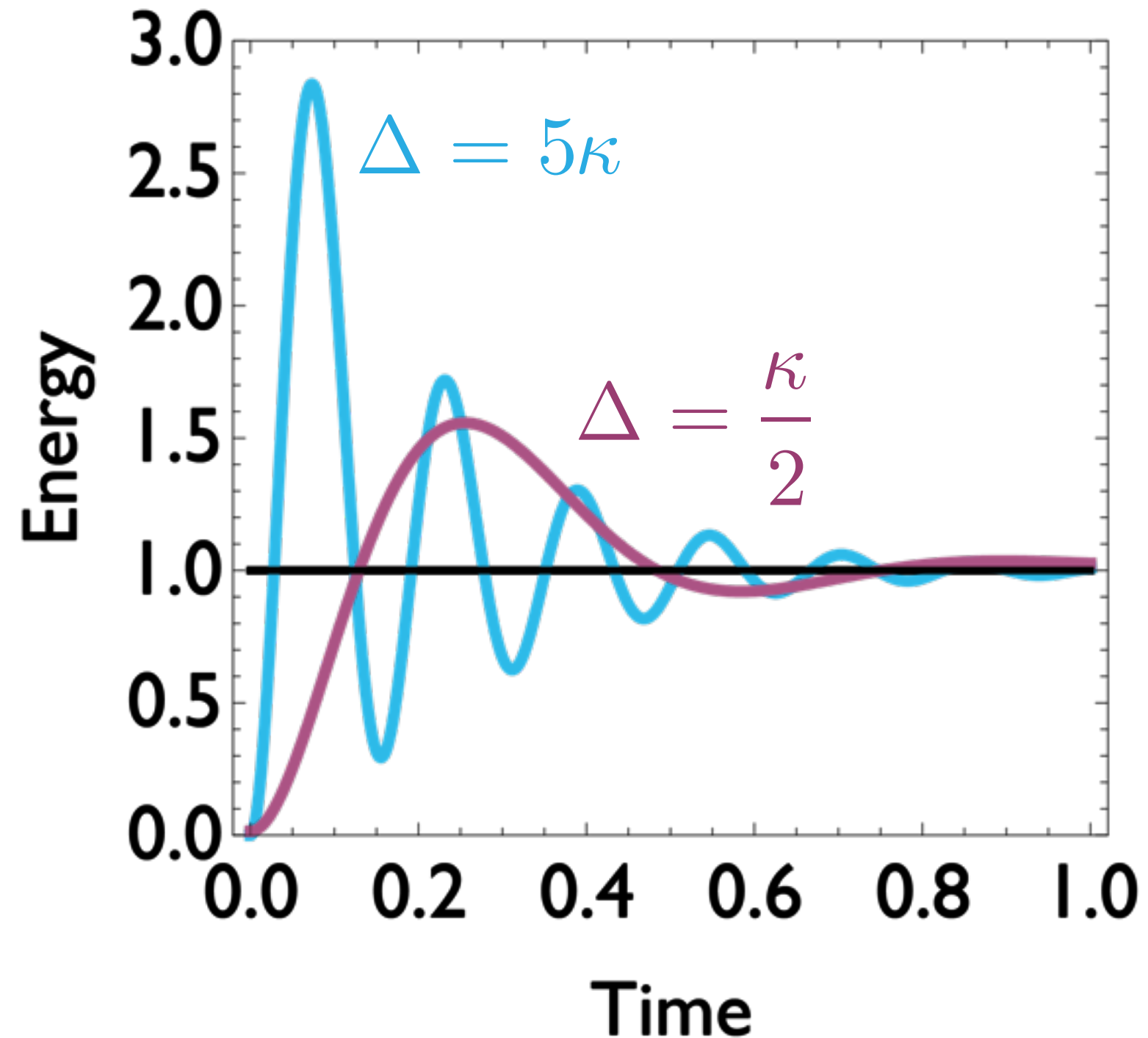
Optical Time Domain Equation



Solution

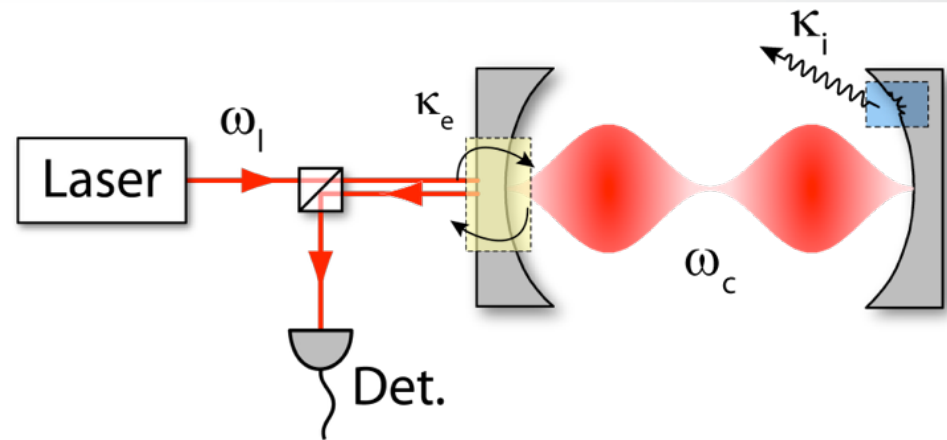
$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2}a + \sqrt{\kappa_e}\alpha_{\text{in}}$$

$$|a(t)|^2 \propto \text{Energy}$$





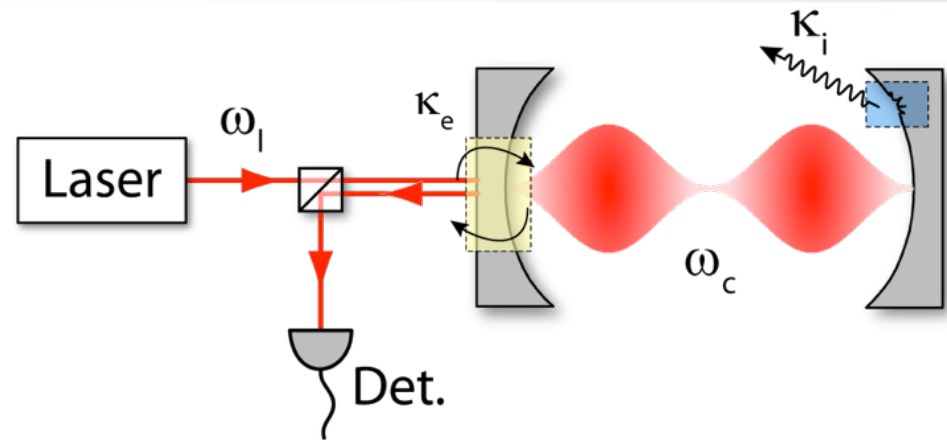
Optical - Frequency Domain



$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2}a + \sqrt{\kappa_e}\alpha_{\text{in}}$$



Optical - Frequency Domain



$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2}a + \sqrt{\kappa_e}\alpha_{\text{in}}$$

Steady state

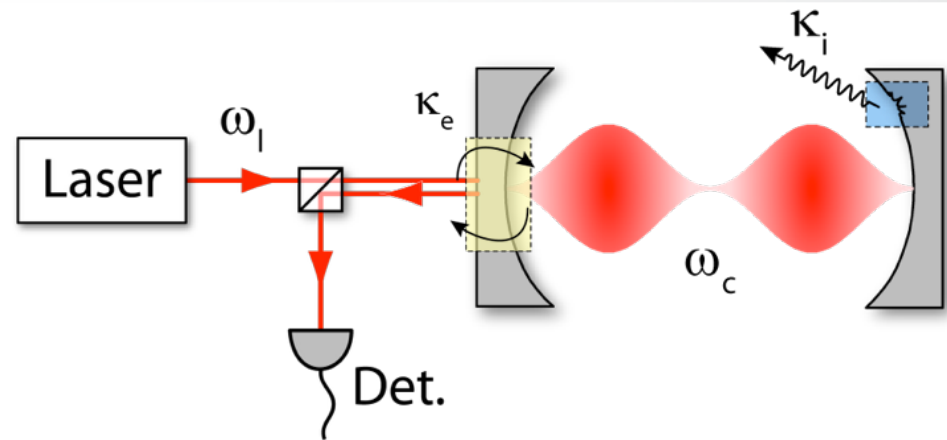
$$\dot{a}(t) = 0$$



$$\langle a \rangle = \frac{\sqrt{\kappa_e}\alpha_{\text{in}}}{\frac{\kappa}{2} - i\Delta}$$



Optical - Frequency Domain



Input-output relation

$$\alpha_{\text{out}} = \alpha_{\text{in}} - \sqrt{\kappa_e} a$$

$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2} a + \sqrt{\kappa_e} \alpha_{\text{in}}$$

Steady state

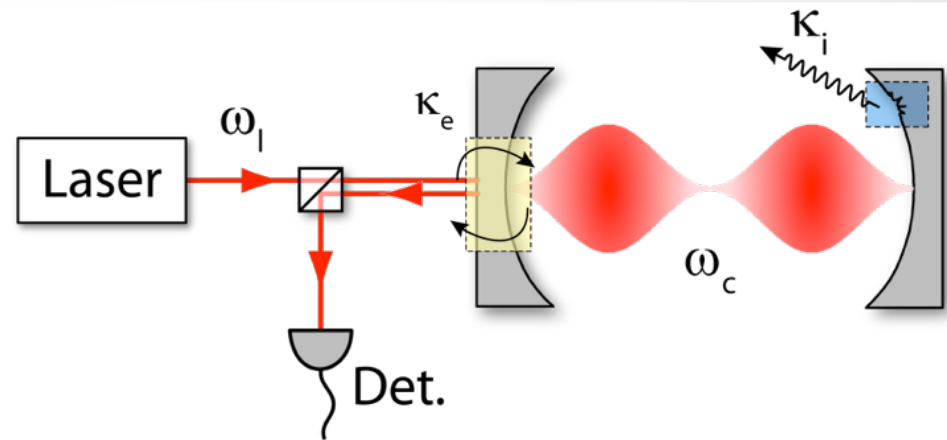
$$\dot{a}(t) = 0$$



$$\langle a \rangle = \frac{\sqrt{\kappa_e} \alpha_{\text{in}}}{\frac{\kappa}{2} - i\Delta}$$



Optical - Frequency Domain



Input-output relation

$$\alpha_{\text{out}} = \alpha_{\text{in}} - \sqrt{\kappa_e} a$$

$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2} a + \sqrt{\kappa_e} \alpha_{\text{in}}$$

Steady state

$$\dot{a}(t) = 0$$

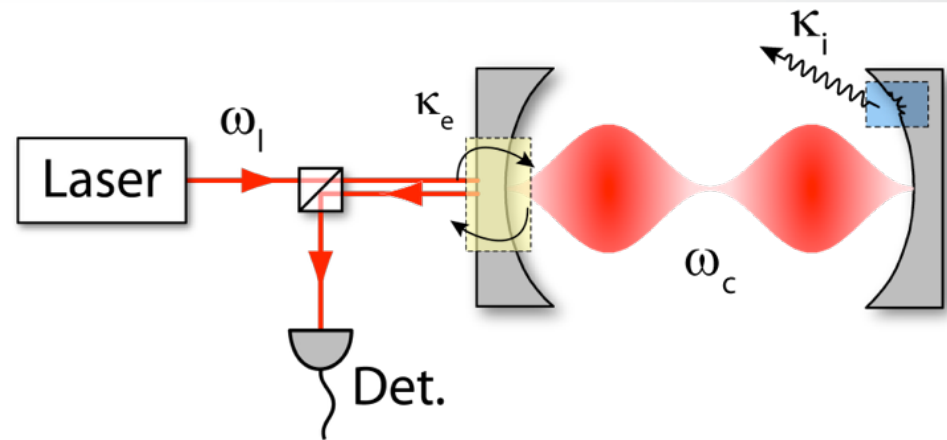


$$\langle a \rangle = \frac{\sqrt{\kappa_e} \alpha_{\text{in}}}{\frac{\kappa}{2} - i\Delta}$$

$$R = \frac{\alpha_{\text{out}}}{\alpha_{\text{in}}} = \frac{(\kappa_i - \kappa_e) / 2 - i\Delta}{(\kappa_i + \kappa_e) / 2 - i\Delta}$$



Optical - Frequency Domain



Input-output relation

$$\alpha_{\text{out}} = \alpha_{\text{in}} - \sqrt{\kappa_e} a$$

$$\dot{a}(t) = i\Delta a - \frac{\kappa}{2} a + \sqrt{\kappa_e} \alpha_{\text{in}}$$

Steady state

$$\dot{a}(t) = 0$$



$$\langle a \rangle = \frac{\sqrt{\kappa_e} \alpha_{\text{in}}}{\frac{\kappa}{2} - i\Delta}$$

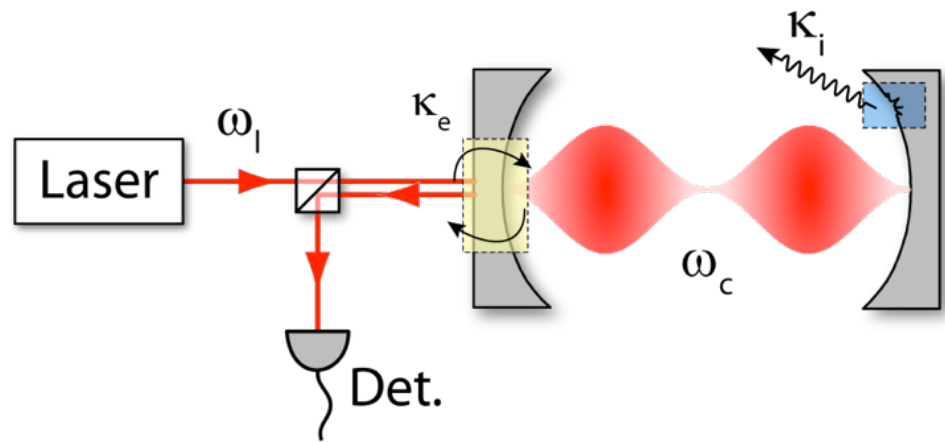
$$R = \frac{\alpha_{\text{out}}}{\alpha_{\text{in}}} = \frac{(\kappa_i - \kappa_e) / 2 - i\Delta}{(\kappa_i + \kappa_e) / 2 - i\Delta}$$

Optical Spectra
(Probability)

$$|R(\omega)|^2 = \frac{(\kappa_i - \kappa_e)^2 + 4\Delta^2}{(\kappa_i + \kappa_e)^2 + 4\Delta^2}$$

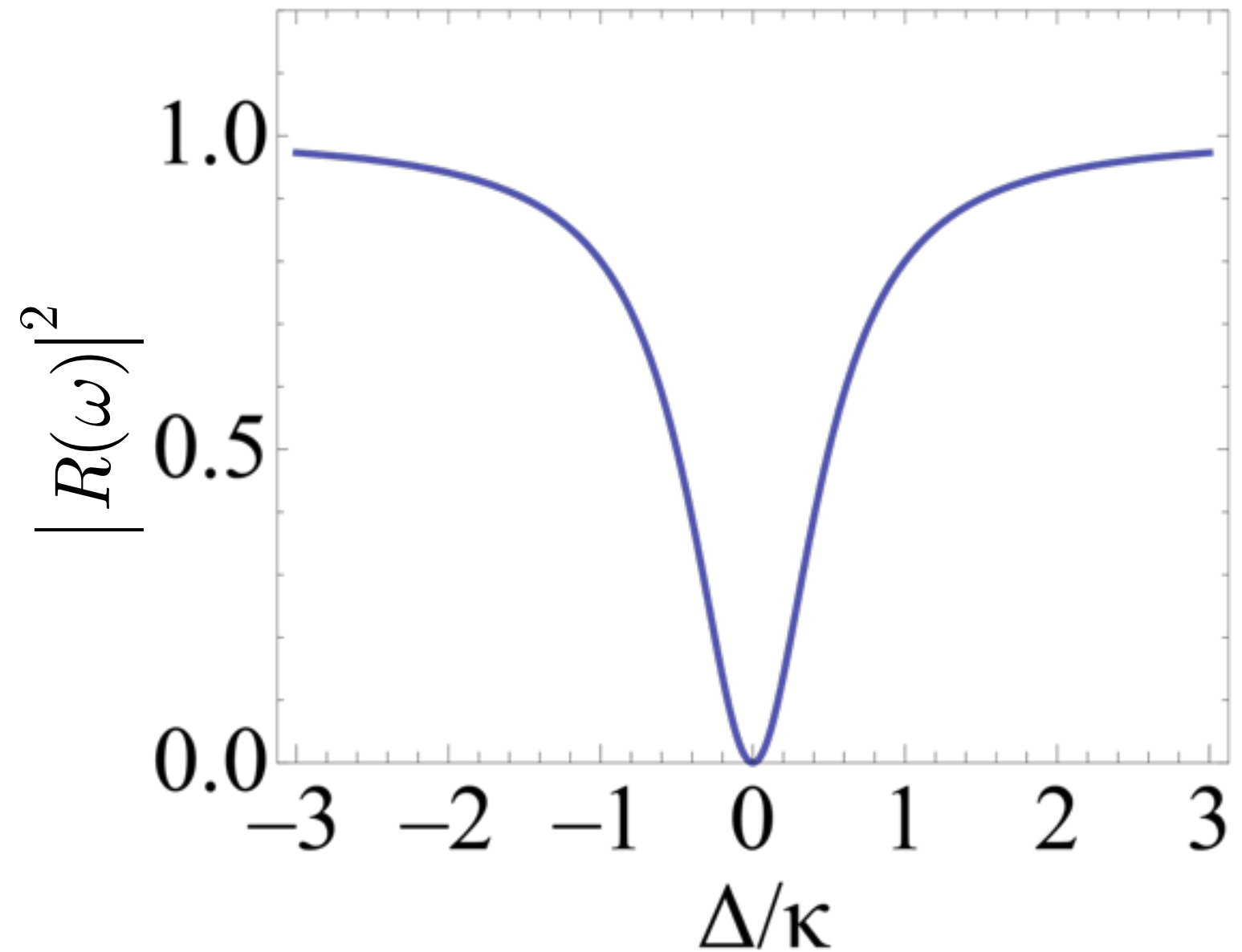


Optical - Frequency Domain



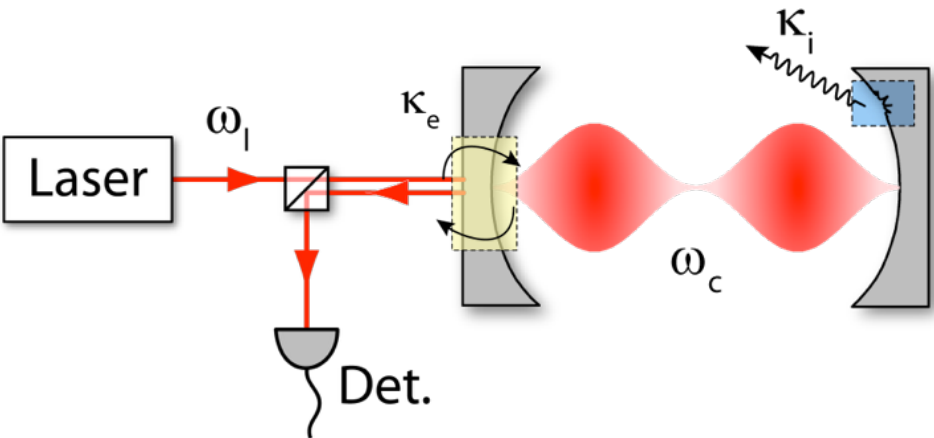
Steady state

$$|R(\omega)|^2 = \frac{(\kappa_i - \kappa_e)^2 + 4\Delta^2}{(\kappa_i + \kappa_e)^2 + 4\Delta^2}$$



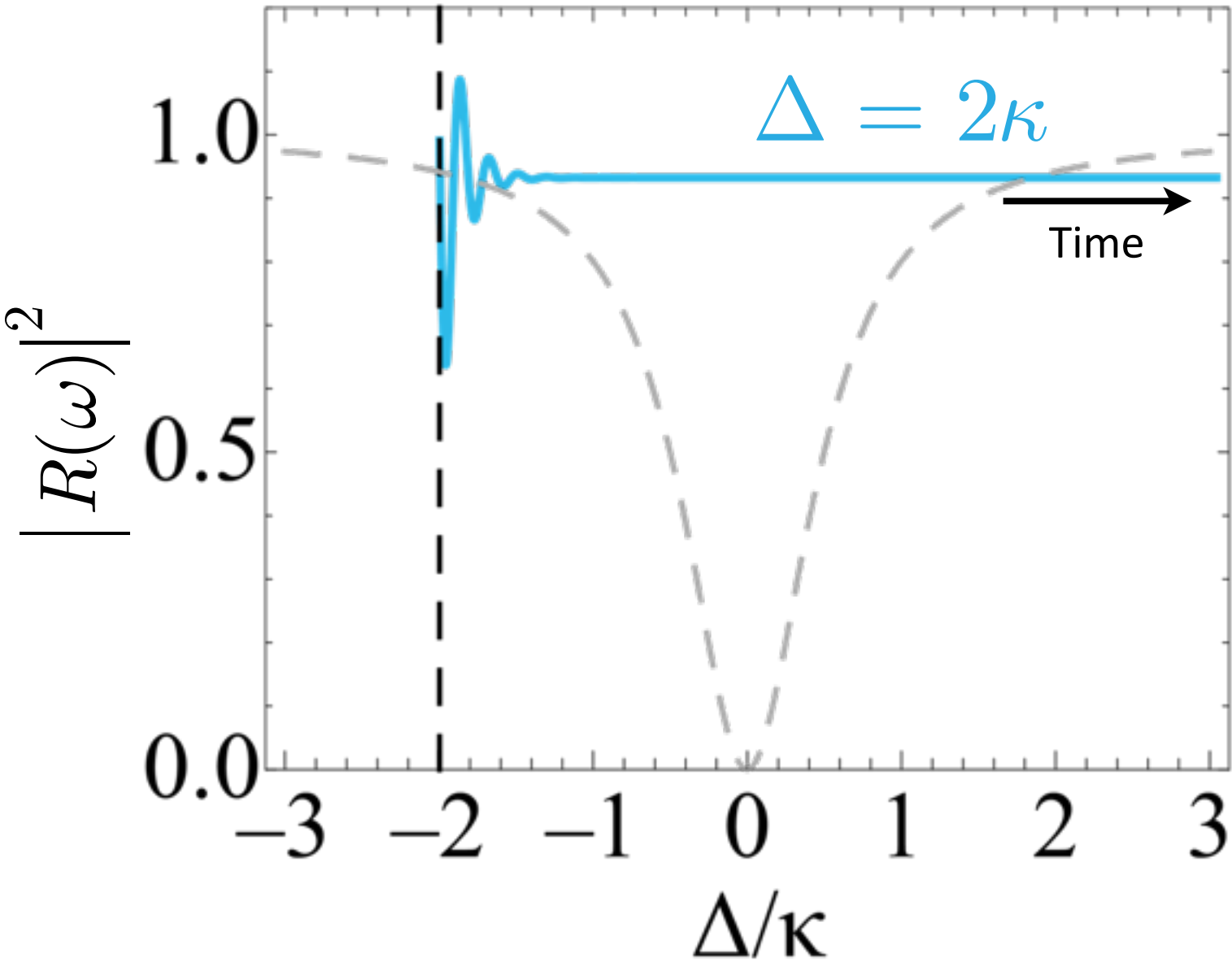


Optical - Frequency Domain



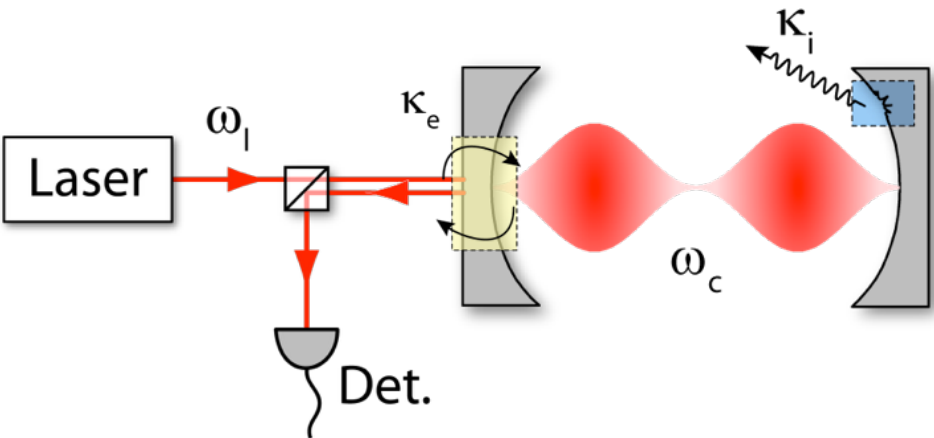
Steady state

$$|R(\omega)|^2 = \frac{(\kappa_i - \kappa_e)^2 + 4\Delta^2}{(\kappa_i + \kappa_e)^2 + 4\Delta^2}$$



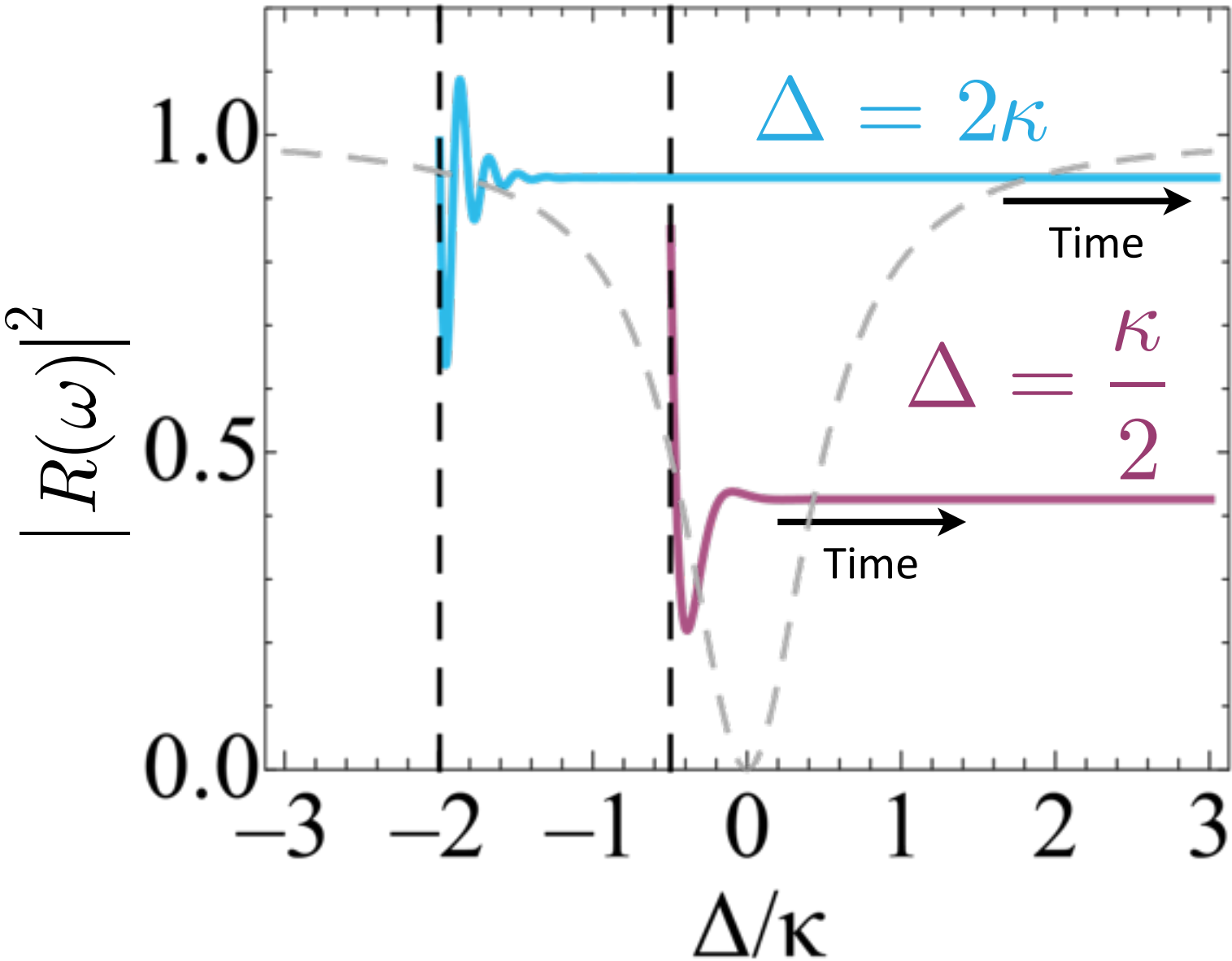


Optical - Frequency Domain



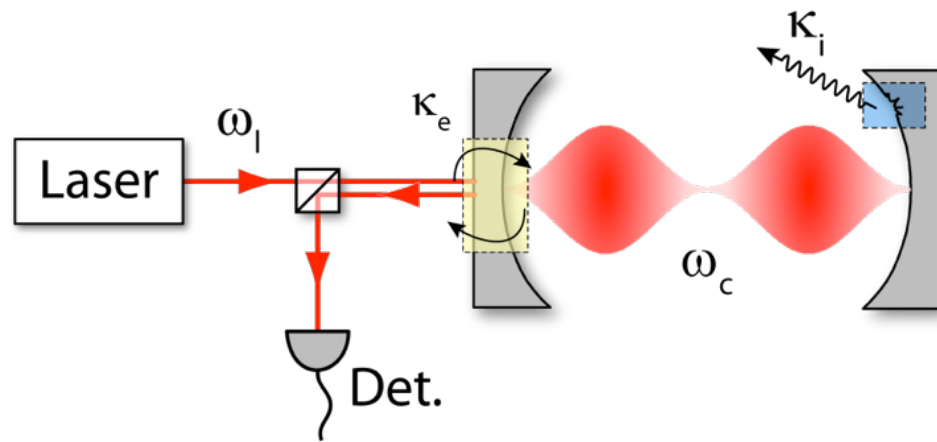
Steady state

$$|R(\omega)|^2 = \frac{(\kappa_i - \kappa_e)^2 + 4\Delta^2}{(\kappa_i + \kappa_e)^2 + 4\Delta^2}$$



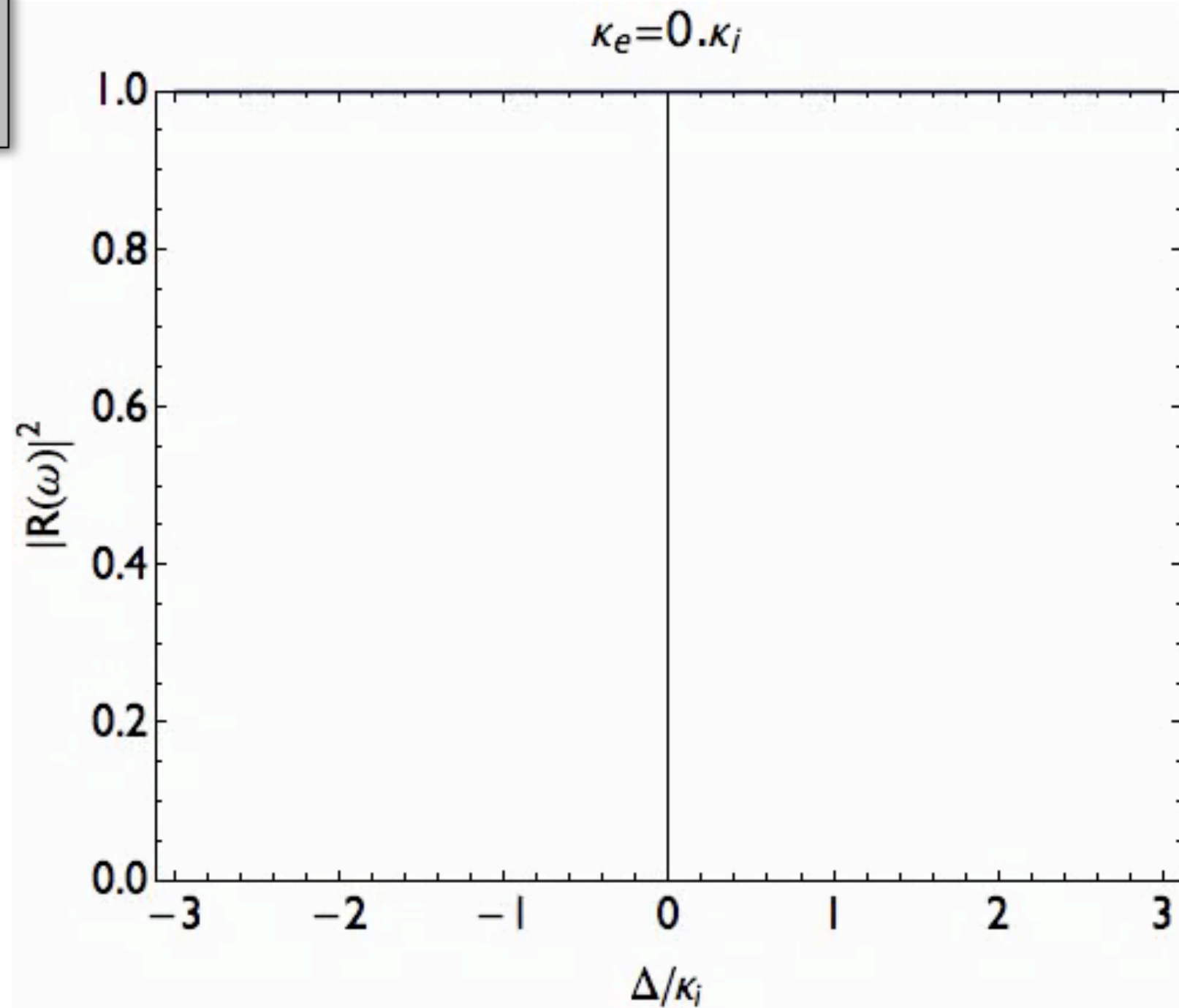


Optical - Frequency Domain



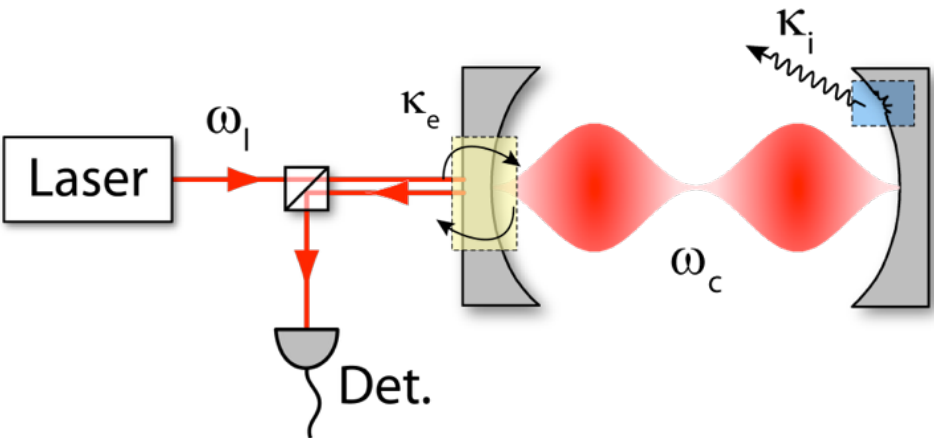
Steady state

$$|R(\omega)|^2 = \frac{(\kappa_i - \kappa_e)^2 + 4\Delta^2}{(\kappa_i + \kappa_e)^2 + 4\Delta^2}$$



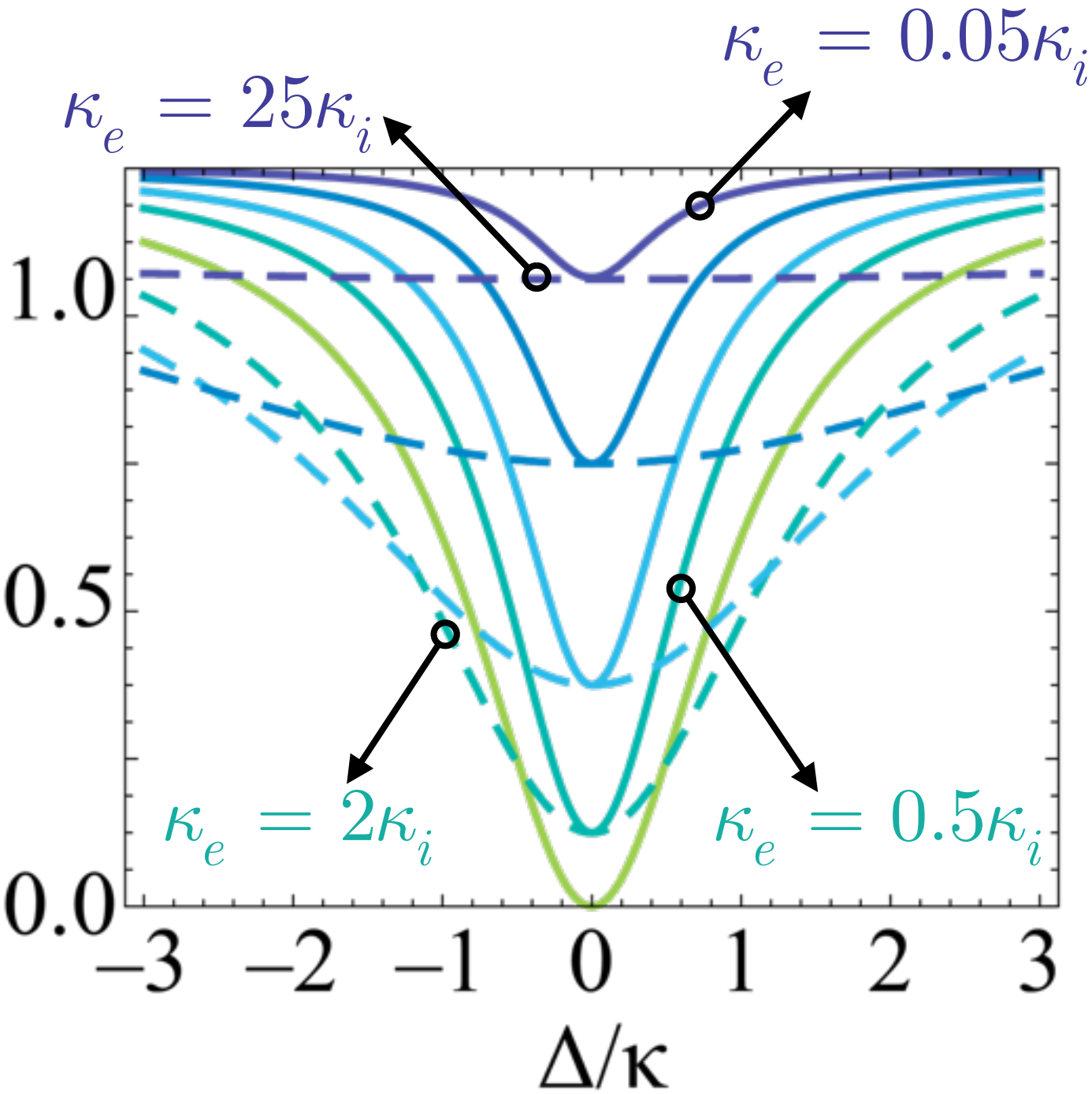


Optical - Frequency Domain



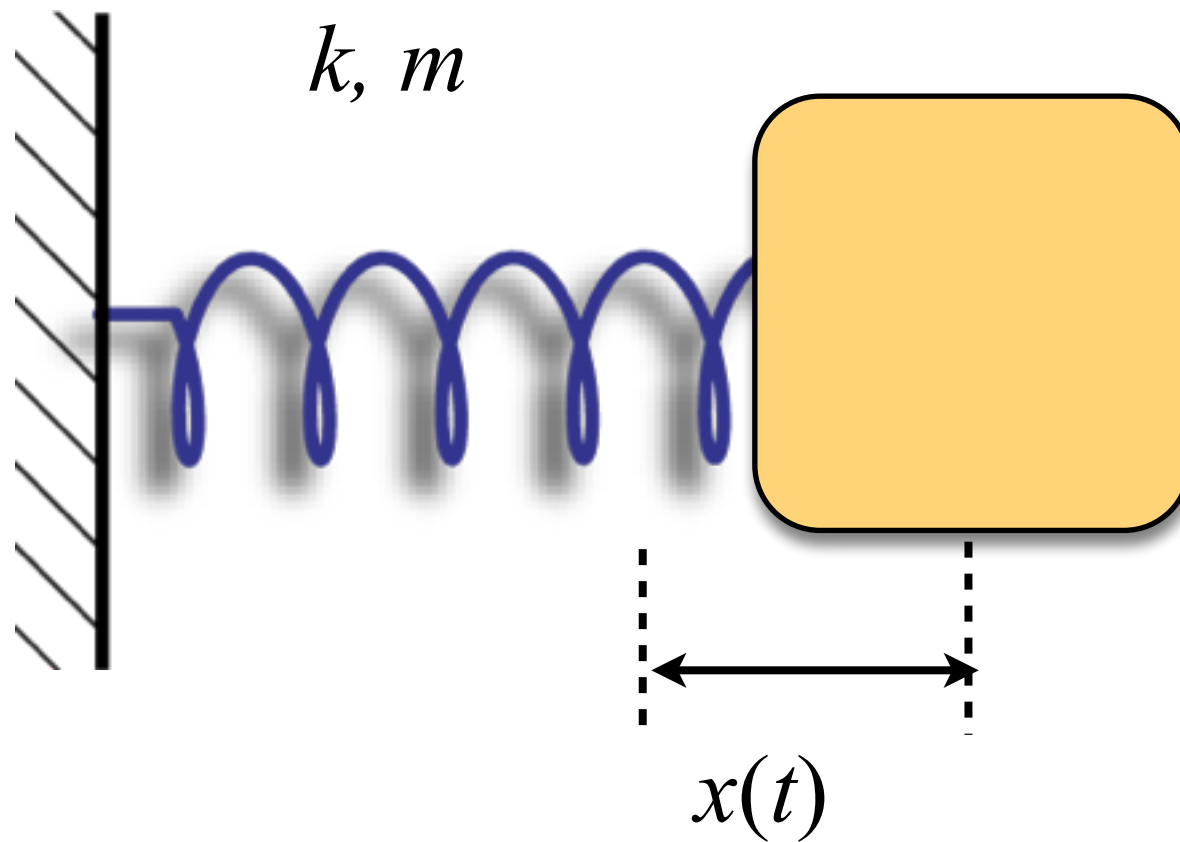
Steady state

$$|R(\omega)|^2 = \frac{(\kappa_i - \kappa_e)^2 + 4\Delta^2}{(\kappa_i + \kappa_e)^2 + 4\Delta^2}$$





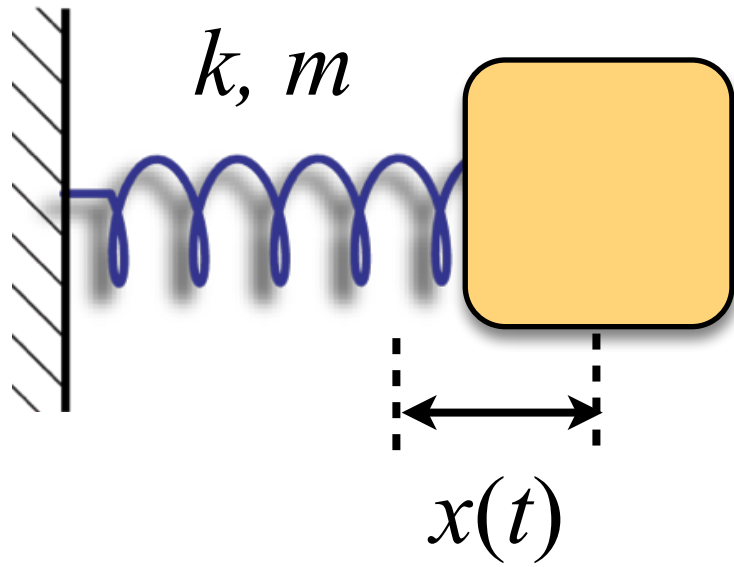
Mechanical modes





Mechanics of continuous

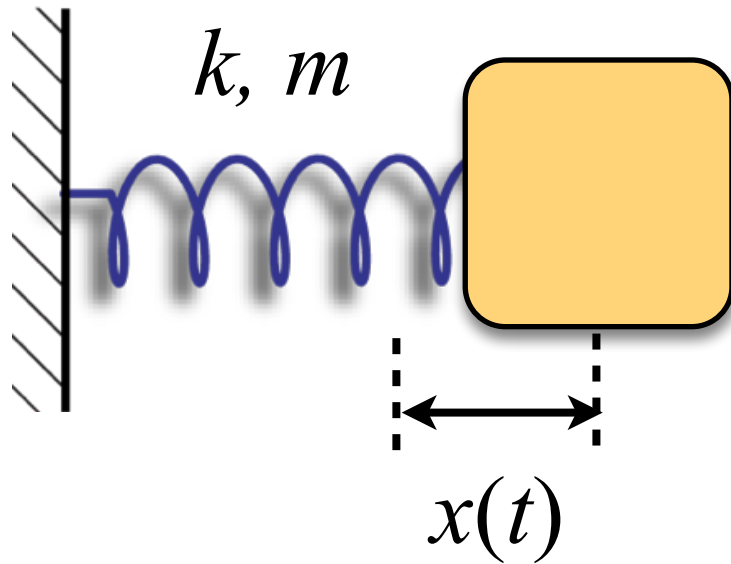
Single particle





Mechanics of continuous

Single particle



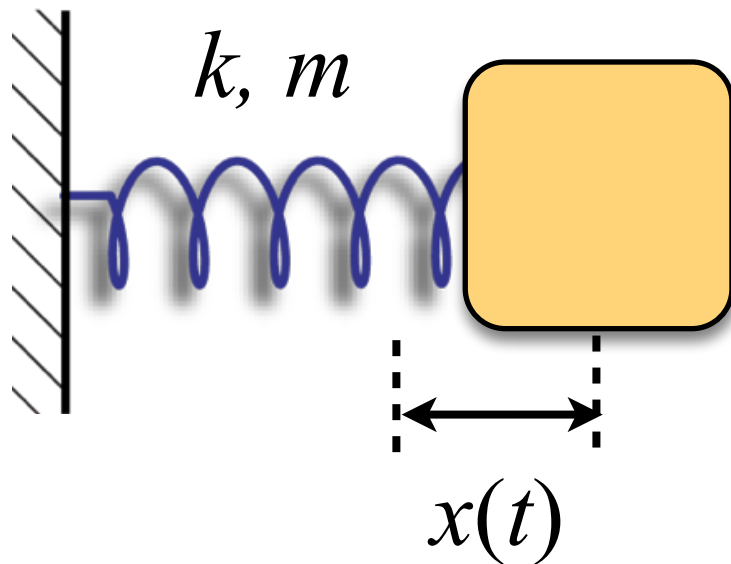
Newton's Law

$$F_x = m \frac{d^2 x}{dt^2}$$



Mechanics of continuous

Single particle



Newton's Law

$$F_x = m \frac{d^2 x}{dt^2}$$

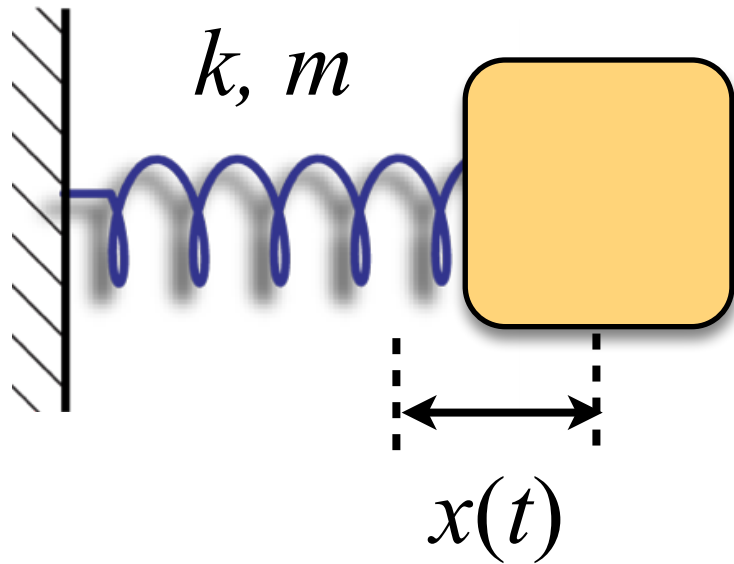
Hooke's Law

$$F_x = -kx$$



Mechanics of continuous

Single particle



Newton's Law

$$F_x = m \frac{d^2 x}{dt^2}$$

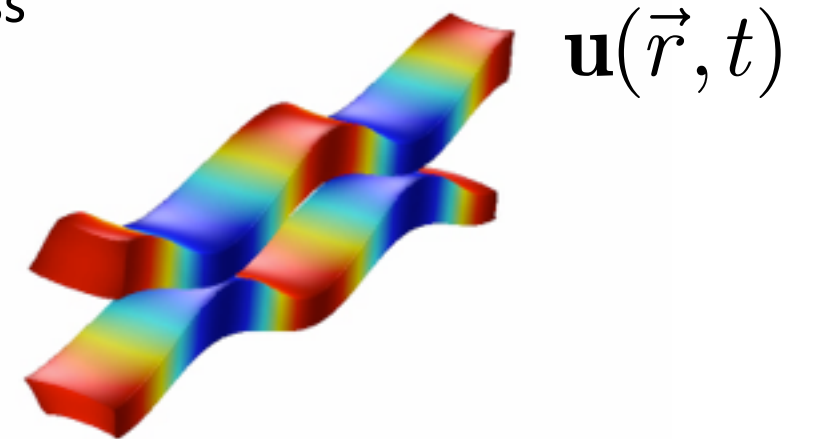
Hooke's Law

$$F_x = -kx$$

Solid mechanics

\mathbf{c} : elastic stiffness

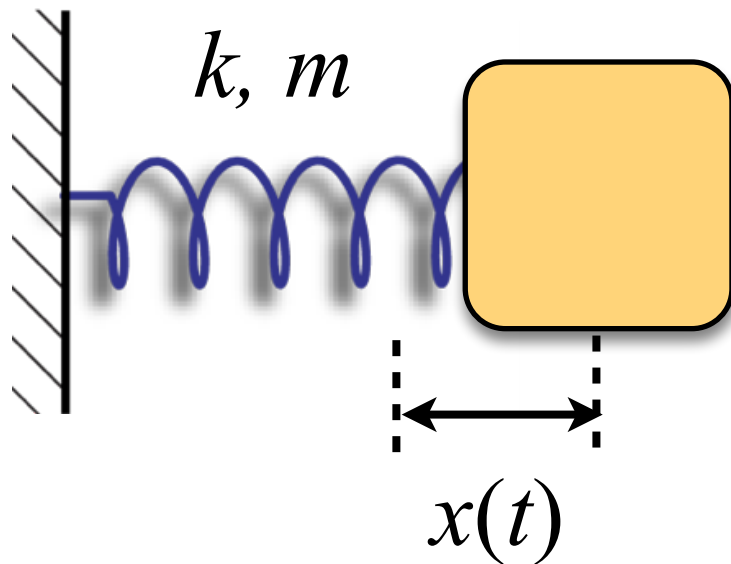
ρ : density





Mechanics of continuous

Single particle



Newton's Law

$$F_x = m \frac{d^2 x}{dt^2}$$

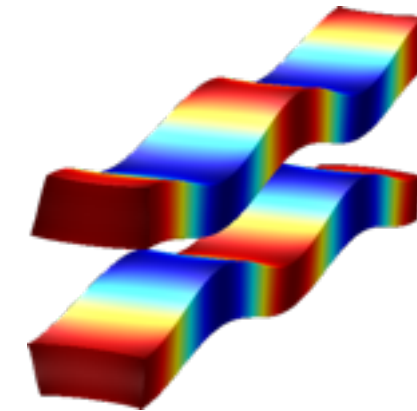
Hooke's Law

$$F_x = -kx$$

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

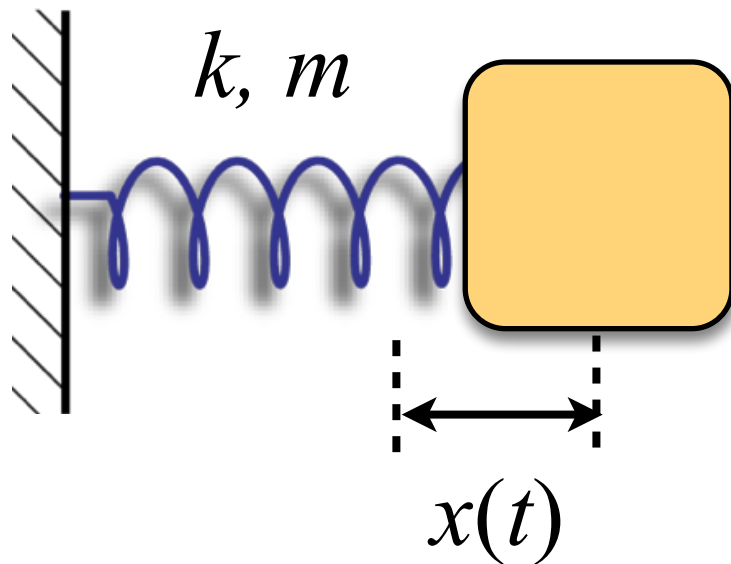
Newton's Law (solid)

$$\nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$



Mechanics of continuous

Single particle



Newton's Law

$$F_x = m \frac{d^2 x}{dt^2}$$

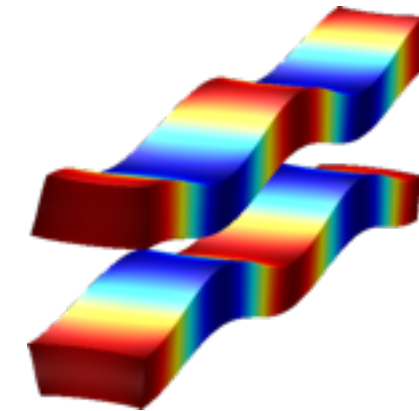
Hooke's Law

$$F_x = -kx$$

Solid mechanics

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ρ : density



$\mathbf{u}(\vec{r}, t)$

Newton's Law (solid)

$$\nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law

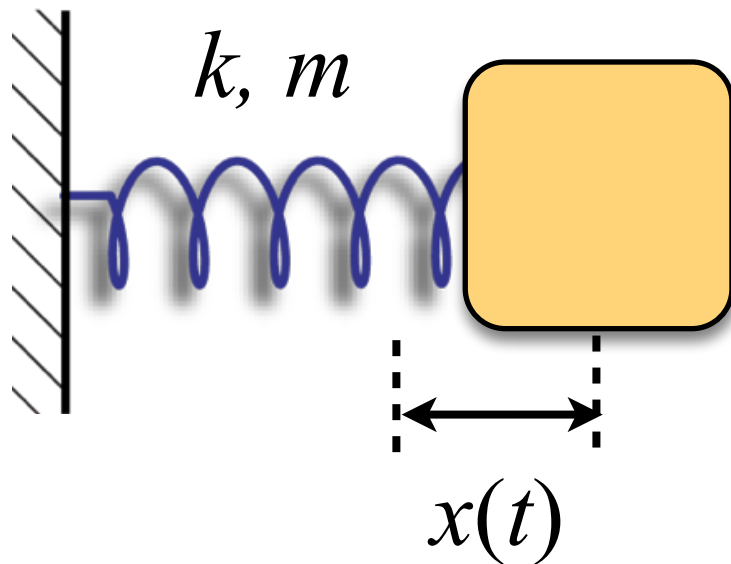
$$\mathbf{T} = \mathbf{c} : \mathbf{S}$$

$$\left(\mathbf{S}_{\text{strain}} = \nabla_s \mathbf{u} \right)$$



Mechanics of continuous

Single particle



Newton's Law

$$F_{\text{ext}} + F_x = m \frac{d^2 x}{dt^2}$$

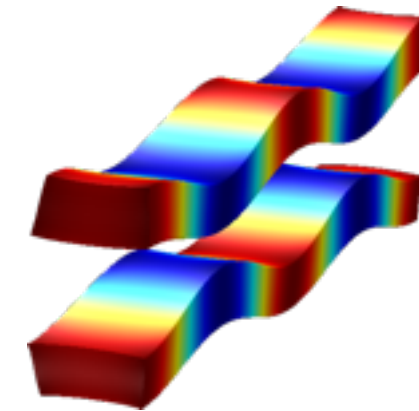
Hooke's Law

$$F_x = -kx$$

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Newton's Law (solid)

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law

$$\nabla \cdot \mathbf{T} = \nabla \cdot (\mathbf{c} : \mathbf{S})$$

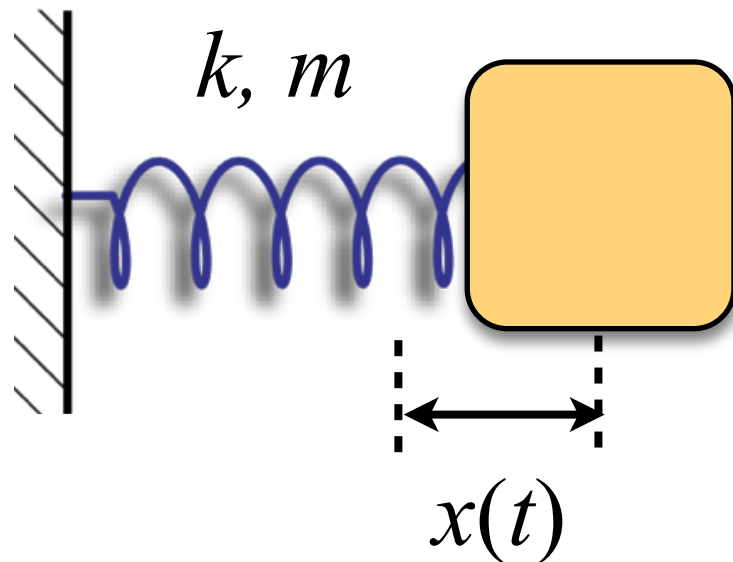
$$\left(\mathbf{S} = \nabla_s \mathbf{u} \right)$$

strain



Mechanics of continuous

Single particle



Newton's Law

$$F_{\text{ext}} + F_x = m \frac{d^2 x}{dt^2}$$

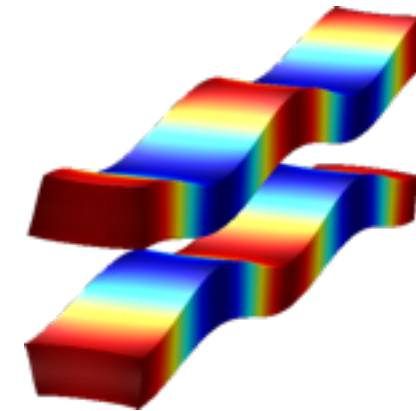
Hooke's Law

$$F_x = -kx$$

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Newton's Law

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot (\mathbf{c} : \mathbf{S}) = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law

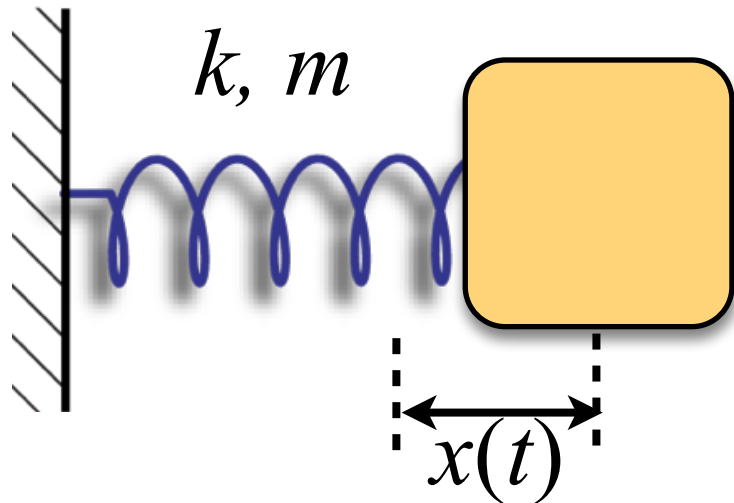
$$\nabla \cdot \mathbf{T} = \nabla \cdot (\mathbf{c} : \mathbf{S})$$

$$\left(\mathbf{S} = \underset{\text{strain}}{\nabla_s \mathbf{u}} \right)$$



Mechanics of continuous

Single particle



Newton's Law

$$F_{\text{ext}} + F_x = m \frac{d^2 x}{dt^2}$$

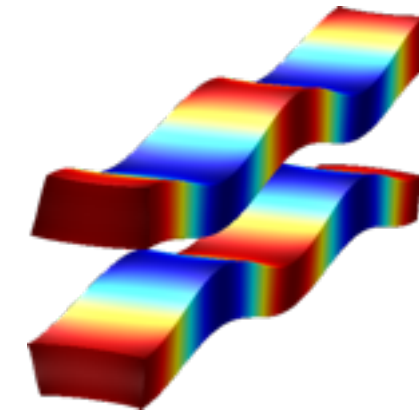
Hooke's Law (single mode)

$$F_x = -m \left(\frac{k}{m} \right) x = -m \Omega_j^2 x$$

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Newton's Law

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot (\mathbf{c} : \mathbf{S}) = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law (single mode)

$$S_j(\vec{r}) = \nabla_s U_j(\vec{r})$$

$$\nabla \cdot (\mathbf{c} : S_j) = -\rho \Omega_j^2 U_j$$



Mechanics of continuous

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Newton's Law

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law
(single mode)

$$S_j(\vec{r}) = \nabla_s U_j(\vec{r})$$

$$\nabla \cdot (\mathbf{c} : S_j) = -\rho \Omega_j^2 U_j$$

Orthogonality relation

$$\int \rho (U_i^* U_j) dV = \delta_{ij} m_{\text{eff}}$$



Mechanics of continuous

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Mechanical Field

$$\mathbf{u}(\vec{r}, t) = \sum b_j(t) U_j(\vec{r})$$

$$\mathbf{S}(\vec{r}, t) = \sum b_j(t) S_j(\vec{r})$$

Newton's Law

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law
(single mode)

$$S_j(\vec{r}) = \nabla_s U_j(\vec{r})$$

$$\nabla \cdot (\mathbf{c} : S_j) = -\rho \Omega_j^2 U_j$$

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Mechanics of continuous

Solid mechanics

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$\mathbf{u}(\vec{r}, t)$

Mechanical Field

$$\mathbf{u}(\vec{r}, t) = \sum b_j(t) U_j(\vec{r})$$

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Newton's Law

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law
(single mode)

$$S_j(\vec{r}) = \nabla_s U_j(\vec{r})$$

$$\nabla \cdot (\mathbf{c} : S_j) = -\rho \Omega_j^2 U_j$$

$$\mathbf{f}_{\text{ext}} + \nabla \cdot (\mathbf{c} : \sum b_j(t) S_j) = \rho \sum \frac{d^2}{dt^2} (b_j(t) U_j)$$

Orthogonality relation

$$\int \rho (U_i^* U_j) dV = \delta_{ij} m_{\text{eff}}$$



Mechanics of continuous

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Mechanical Field

$$\mathbf{u}(\vec{r}, t) = \sum b_j(t) U_j(\vec{r})$$

$$\mathbf{S}(\vec{r}, t) = \sum b_j(t) S_j(\vec{r})$$

Newton's Law

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law
(single mode)

$$S_j(\vec{r}) = \nabla_s U_j(\vec{r})$$

$$\nabla \cdot (\mathbf{c} : S_j) = -\rho \Omega_j^2 U_j$$

$$\mathbf{f}_{\text{ext}} + \nabla \cdot (\mathbf{c} : \sum b_j(t) S_j) = \rho \sum \frac{d^2}{dt^2} (b_j(t) U_j)$$

$$\int \left[\mathbf{f}_{\text{ext}} - \sum \Omega_j^2 b_j(t) \rho U_j = \sum \ddot{b}_j(t) \rho U_j \right] U_i^* dV$$

Orthogonality relation

$$\int \rho (U_i^* U_j) dV = \delta_{ij} m_{\text{eff}}$$



Mechanics of continuous

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Mechanical Field

$$\mathbf{u}(\vec{r}, t) = \sum b_j(t) U_j(\vec{r})$$

$$\mathbf{S}(\vec{r}, t) = \sum b_j(t) S_j(\vec{r})$$

Newton's Law

$$\nabla \cdot \mathbf{T}_{\text{ext}} + \nabla \cdot \mathbf{T} = \rho \frac{d^2 \mathbf{u}(\vec{r}, t)}{dt^2}$$

Hooke's Law (single mode)

$$S_j(\vec{r}) = \nabla_s U_j(\vec{r})$$

$$\nabla \cdot (\mathbf{c} : S_j) = -\rho \Omega_j^2 U_j$$

Orthogonality relation

$$\int \rho (U_i^* U_j) dV = \delta_{ij} m_{\text{eff}}$$

$$\mathbf{f}_{\text{ext}} + \nabla \cdot (\mathbf{c} : \sum b_j(t) S_j) = \rho \sum \frac{d^2}{dt^2} (b_j(t) U_j)$$

$$\int \left[\mathbf{f}_{\text{ext}} - \sum \Omega_j^2 b_j(t) \rho U_j = \sum \ddot{b}_j(t) \rho U_j \right] U_i^* dV$$

Lumped Model

$$\ddot{b}_j(t) + \Omega_i^2 b_i(t) = \frac{\mathbf{F}_{i,\text{ext}}}{m_{\text{eff}}}$$



Lumped Model - Frequency Domain

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Mechanical Loss Channel

$$\ddot{b}(t) + \gamma_m \dot{b}(t) + \Omega_m^2 b(t) = \frac{\mathbf{F}_{\text{ext}}}{m_{\text{eff}}}$$



Lumped Model - Frequency Domain

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Mechanical Loss Channel

$$\ddot{b}(t) + \gamma_m \dot{b}(t) + \Omega_m^2 b(t) = \frac{\mathbf{F}_{\text{ext}}}{m_{\text{eff}}}$$

$$\tilde{b}(\Omega) = \int_{-\infty}^{+\infty} dt \, e^{i\Omega t} b(t)$$



Lumped Model - Frequency Domain

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Mechanical Loss Channel

$$\ddot{b}(t) + \gamma_m \dot{b}(t) + \Omega_m^2 b(t) = \frac{\mathbf{F}_{\text{ext}}}{m_{\text{eff}}}$$

$$\tilde{b}(\Omega) = \int_{-\infty}^{+\infty} dt \, e^{i\Omega t} b(t)$$

$$-\Omega^2 \tilde{b}(\Omega) - i\Omega \gamma_m \tilde{b}(\Omega) + \Omega_m^2 \tilde{b}(\Omega) = \frac{\tilde{\mathbf{F}}_{\text{ext}}(\Omega)}{m_{\text{eff}}}$$



Lumped Model - Frequency Domain

Solid mechanics

\mathbf{c} : elastic stiffness

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$\mathbf{u}(\vec{r}, t)$

Mechanical Loss Channel

$$\ddot{b}(t) + \gamma_m \dot{b}(t) + \Omega_m^2 b(t) = \frac{\mathbf{F}_{\text{ext}}}{m_{\text{eff}}}$$

$$\tilde{b}(\Omega) = \int_{-\infty}^{+\infty} dt e^{i\Omega t} b(t)$$

$$-\Omega^2 \tilde{b}(\Omega) - i\Omega \gamma_m \tilde{b}(\Omega) + \Omega_m^2 \tilde{b}(\Omega) = \frac{\tilde{\mathbf{F}}_{\text{ext}}(\Omega)}{m_{\text{eff}}}$$

Mechanical Frequency Response

$$\tilde{b}(\Omega) = \chi_{bb}(\Omega) \tilde{\mathbf{F}}_{\text{ext}}(\Omega)$$

Mechanical Susceptibility

$$\chi_{bb}(\Omega) = \left[m_{\text{eff}} \left(\Omega_m^2 - \Omega^2 - i\Omega \gamma_m \right) \right]^{-1}$$



Lumped Model - Frequency Domain

Solid mechanics

\mathbf{c} : elastic stiffness

ρ : density



$\mathbf{u}(\vec{r}, t)$

Mechanical Frequency Response

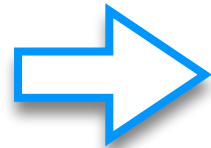
$$\tilde{\mathbf{b}}(\Omega) = \chi_{bb}(\Omega) \tilde{\mathbf{F}}_{\text{ext}}(\Omega)$$

Mechanical Susceptibility

$$\chi_{bb}(\Omega) = \left[m_{\text{eff}} \left(\Omega_m^2 - \Omega^2 - i\Omega\gamma_m \right) \right]^{-1}$$

DC response

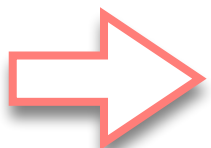
$$\Omega \approx 0$$



$$\chi_{bb}(0) = \left[m_{\text{eff}} \Omega_m^2 \right]^{-1} = 1 / k$$

Lorentzian shape

$$\Omega \approx \Omega_m$$

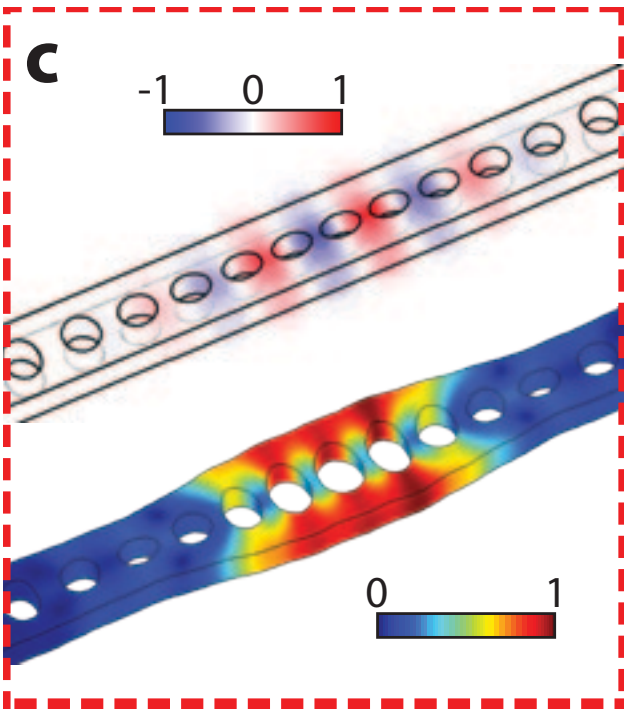
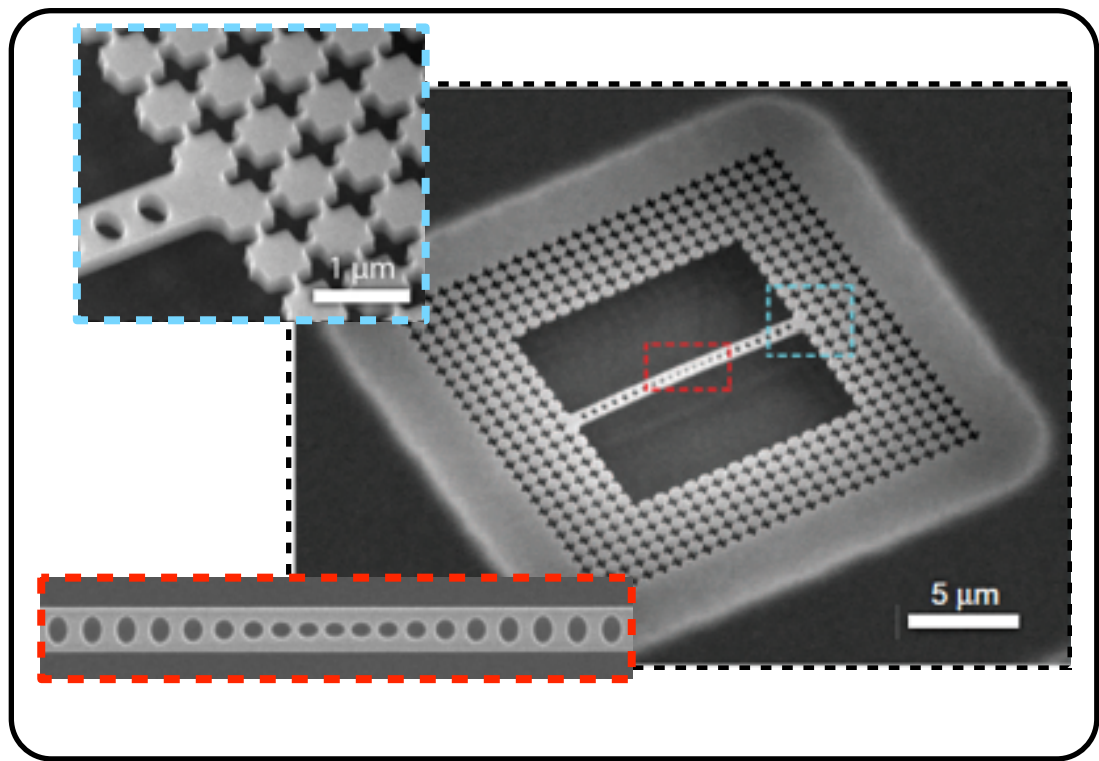


$$\chi_{bb}(\Omega) \approx \left[2m_{\text{eff}} \Omega_m \left((\Omega_m - \Omega) - i\gamma_m/2 \right) \right]^{-1}$$



Lumped Model - Frequency Domain

Optomechanical Crystal



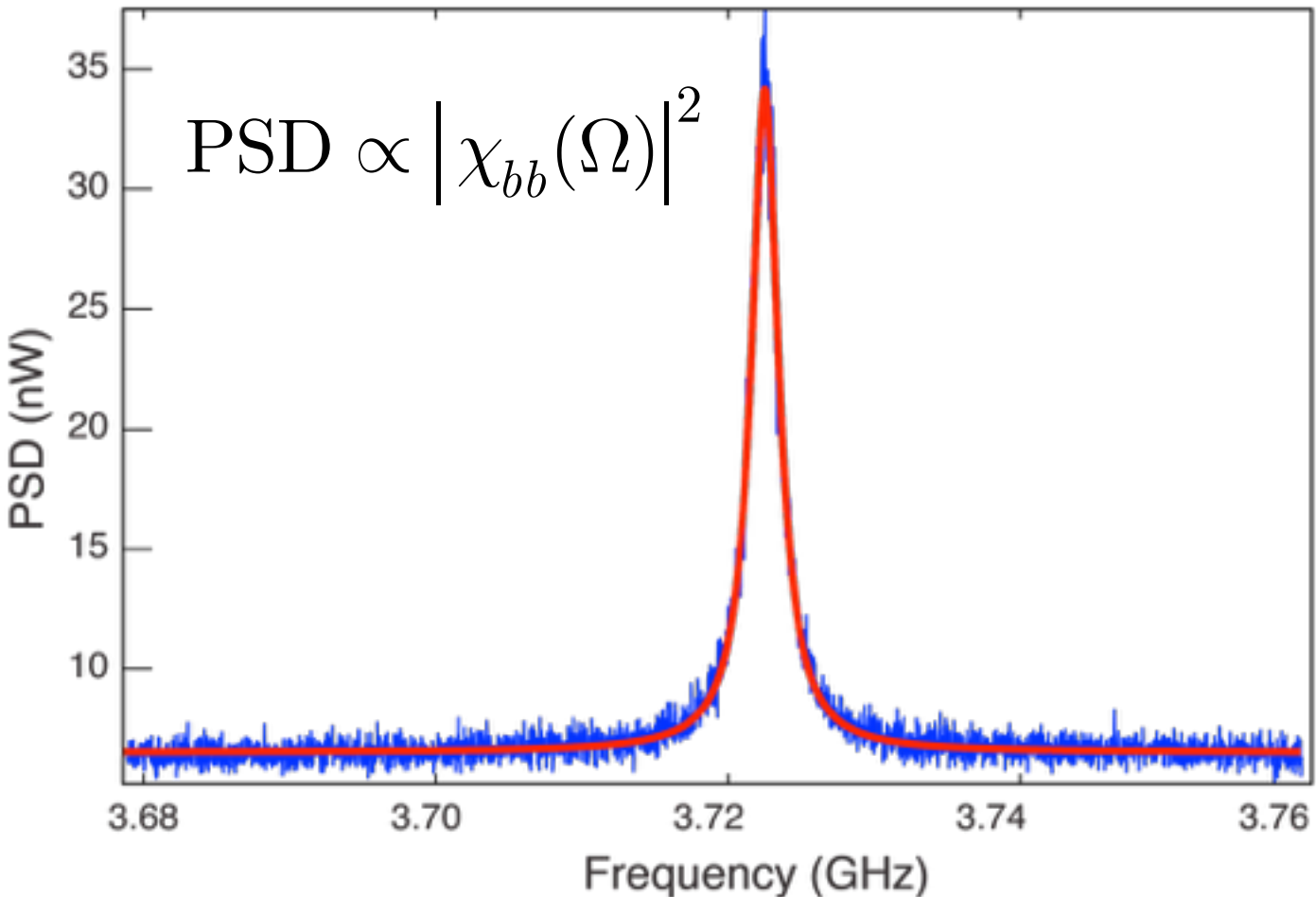
Mechanical Frequency Response

$$\tilde{b}(\Omega) = \chi_{bb}(\Omega) \tilde{\mathbf{F}}_{\text{ext}}(\Omega)$$

$$\chi_{bb}(\Omega) \approx \left[2m_{\text{eff}}\Omega_m \left((\Omega_m - \Omega) - i\gamma_m/2 \right) \right]^{-1}$$

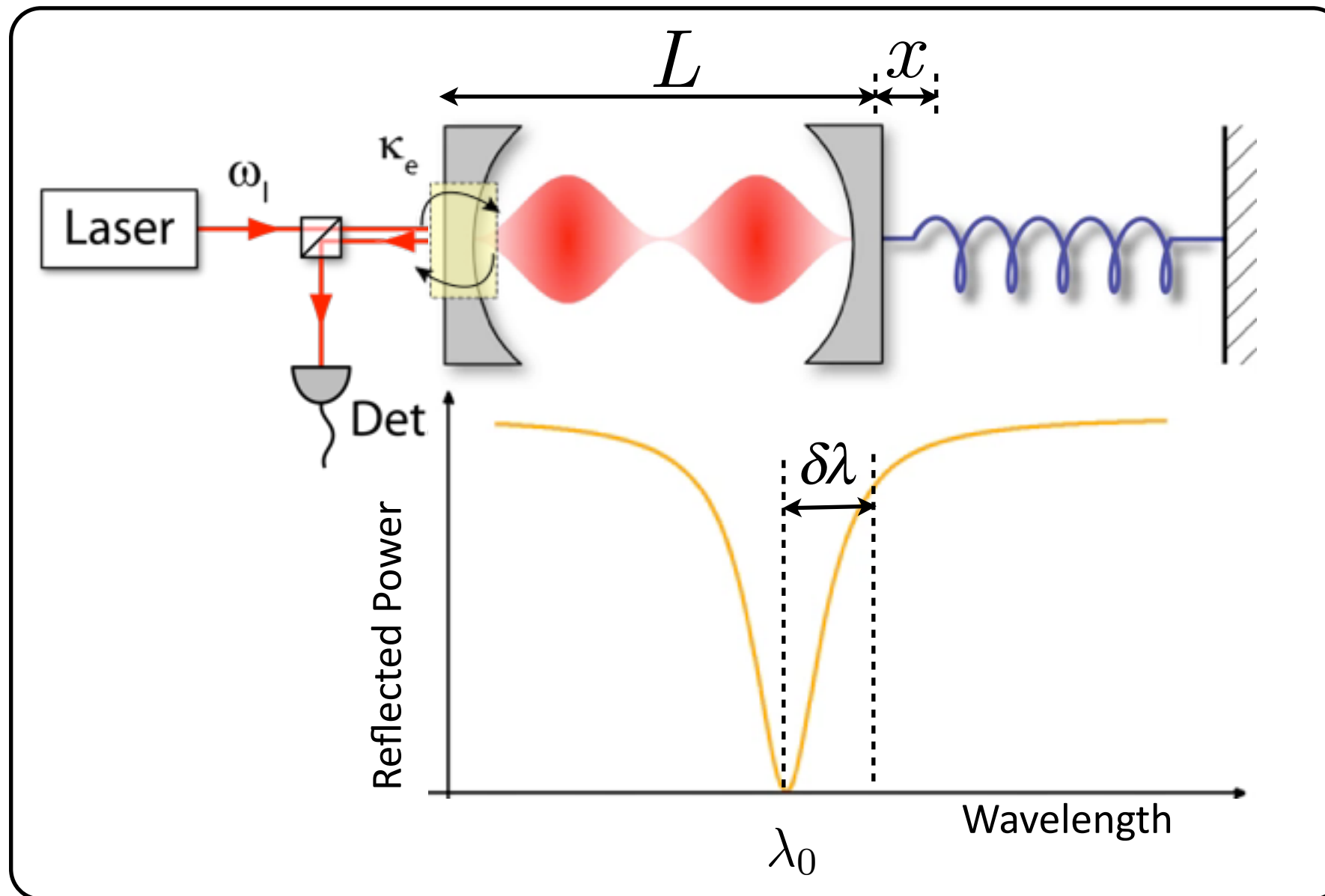
Lorentzian shape

$$|\chi_{bb}(\Omega)|^2 \approx \frac{1}{4m_{\text{eff}}^2\Omega_m^2 \left((\Omega_m - \Omega)^2 - (\gamma_m/2)^2 \right)}$$






Optomechanical cavity toy-model



$$\omega_c = n \frac{\pi c}{L}$$

$$\omega_c(x) \approx \omega_c + \frac{\partial \omega}{\partial x} x = \omega_c - \left(\frac{\omega}{L} \right) x$$

 g_{om}



Outline

- ★ Optical and acoustic mode interaction
- ★ **Optical force actuation**
- ★ Dynamical back-action
- ★ Optomechanical clocks
- ★ Bullseye - a case study
- ★ Outlook



Optical Cavities: Harnessing Light Force

$$\omega_c(x) \approx \omega_c - g_{\text{om}} x$$

$$g_{\text{om}} = \omega / L$$

Frequency pull
parameter

$$f = \frac{\Delta p}{\Delta t}$$

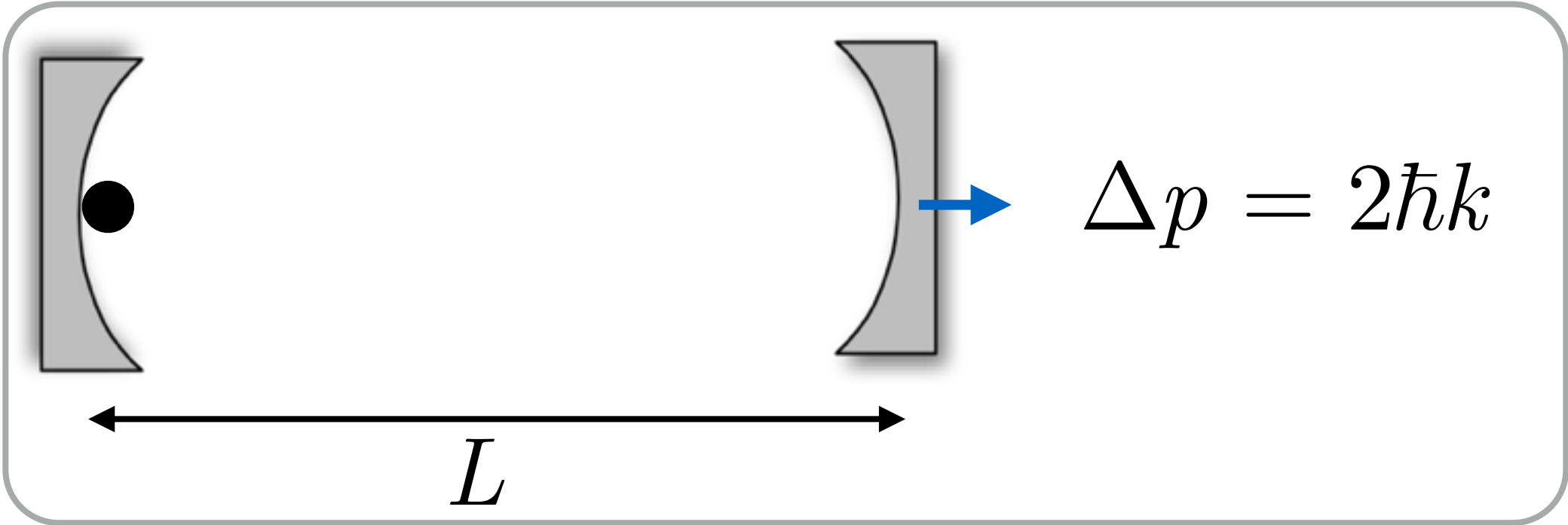
Single-photon
force

$$= \frac{2\hbar k}{\Delta t} = 2\hbar k \left(\frac{c}{2L} \right) = \hbar \frac{\omega_c}{L}$$

$$\hbar g_{\text{om}}$$

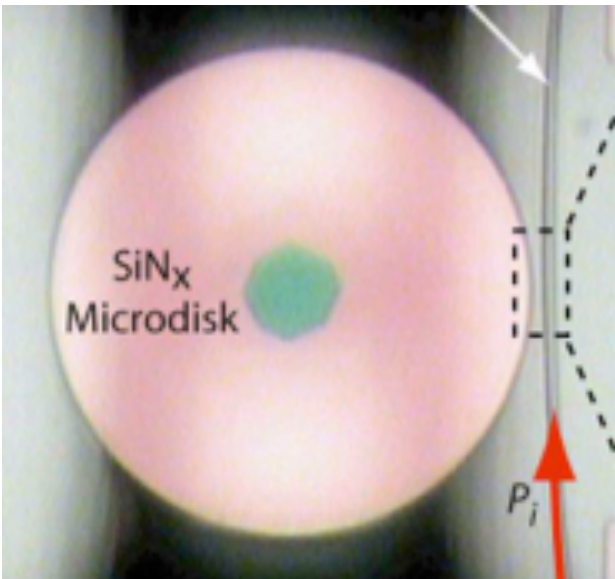
single-
photon
force

$$\approx 2 \text{ fN}$$

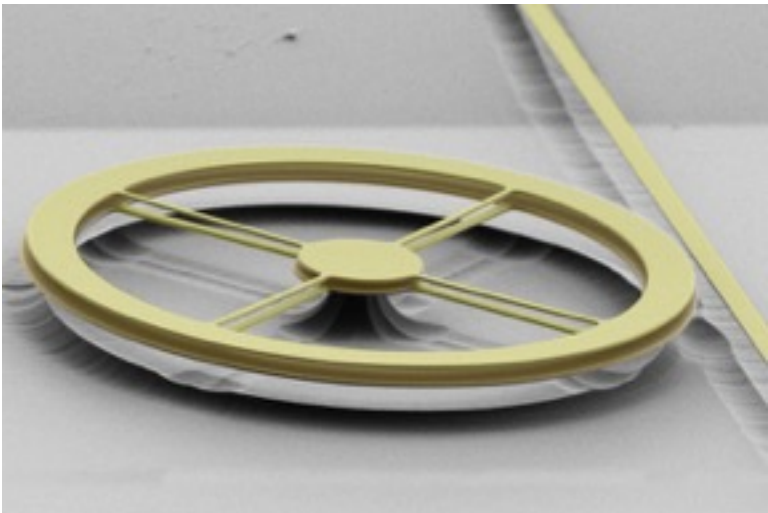




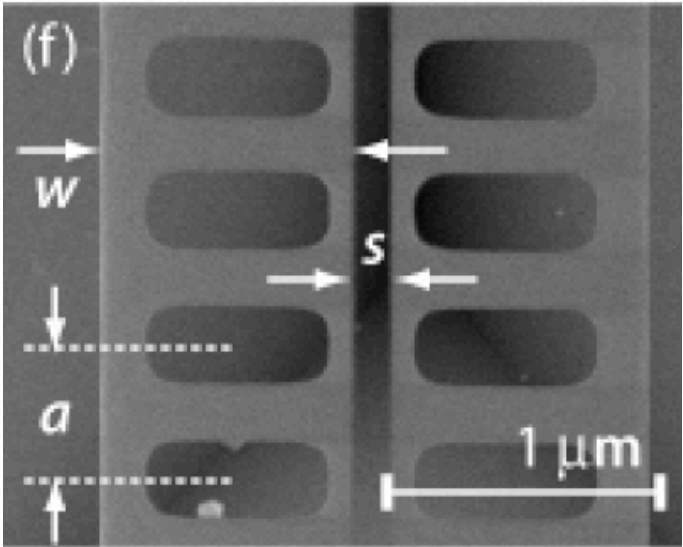
Optical Forces Among Guiding Structures



Eichenfield et al.
Nature Photonics (2007)



Wiederhecker et al.
Nature (2009)

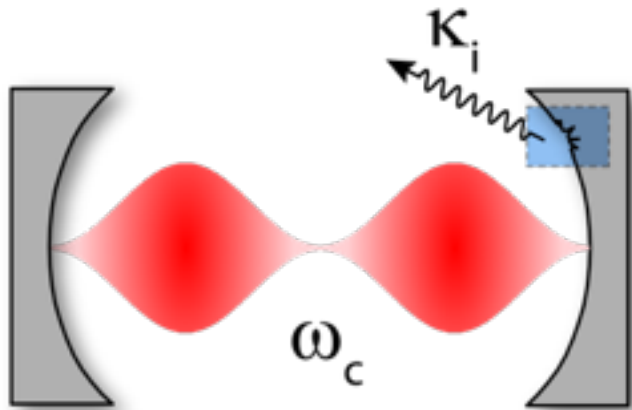


Eichenfield et al.
Nature (2009)



Optical Cavities: Harnessing Light Force

Fabry-Perot



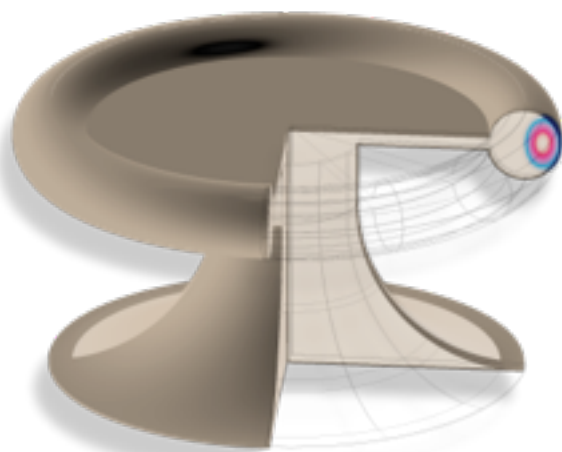
$$L = 2.5 \text{ cm}$$

$$Q = 10^6 \text{ (1 ns)}$$

$$\mathcal{F} \approx 10$$

$$10^{-5} \text{ pN}$$

Ring resonator



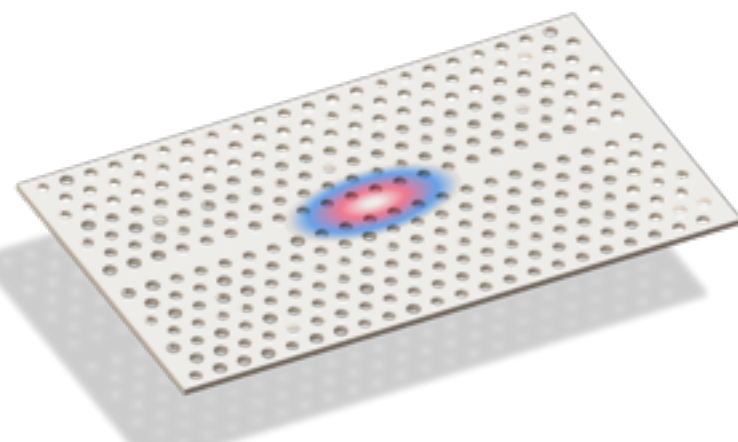
$$L = 2\pi R \approx 60 \text{ }\mu\text{m}$$

$$Q = 10^6 \text{ (1 ns)}$$

$$\mathcal{F} \approx 4 \times 10^3$$

$$10^{-3} \text{ pN}$$

PhC



$$L = 4 \text{ }\mu\text{m}$$

$$Q = 10^6 \text{ (1 ns)}$$

$$\mathcal{F} \approx 60 \times 10^3$$

$$10^{-2} \text{ pN}$$

$$f = \hbar g_{\text{om}}$$

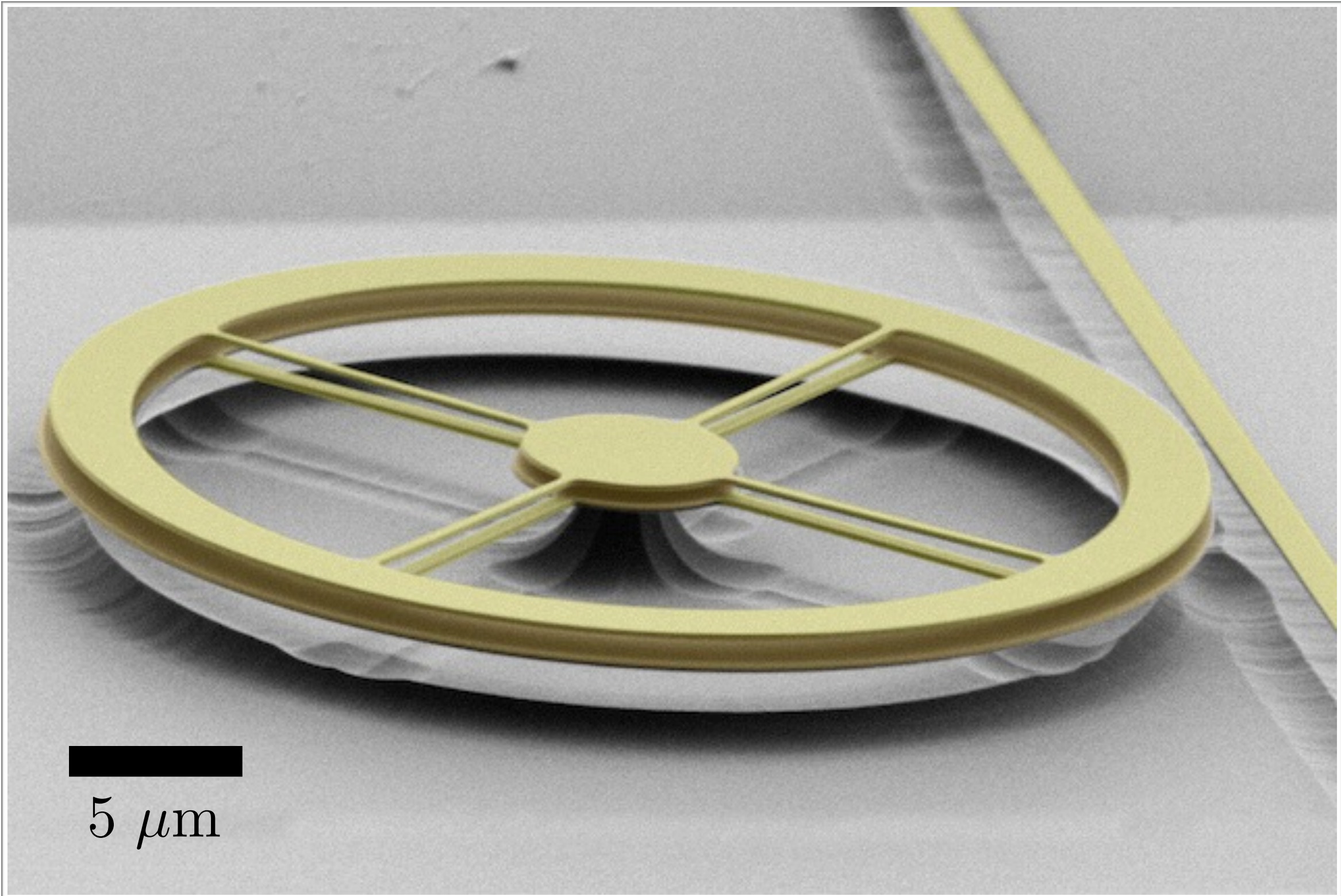
force due to single-photon

A standard laser pointer can load a cavity with as much as 10 million photons!

$$F = 10^6 \hbar g_{\text{om}} \approx \text{nN}$$

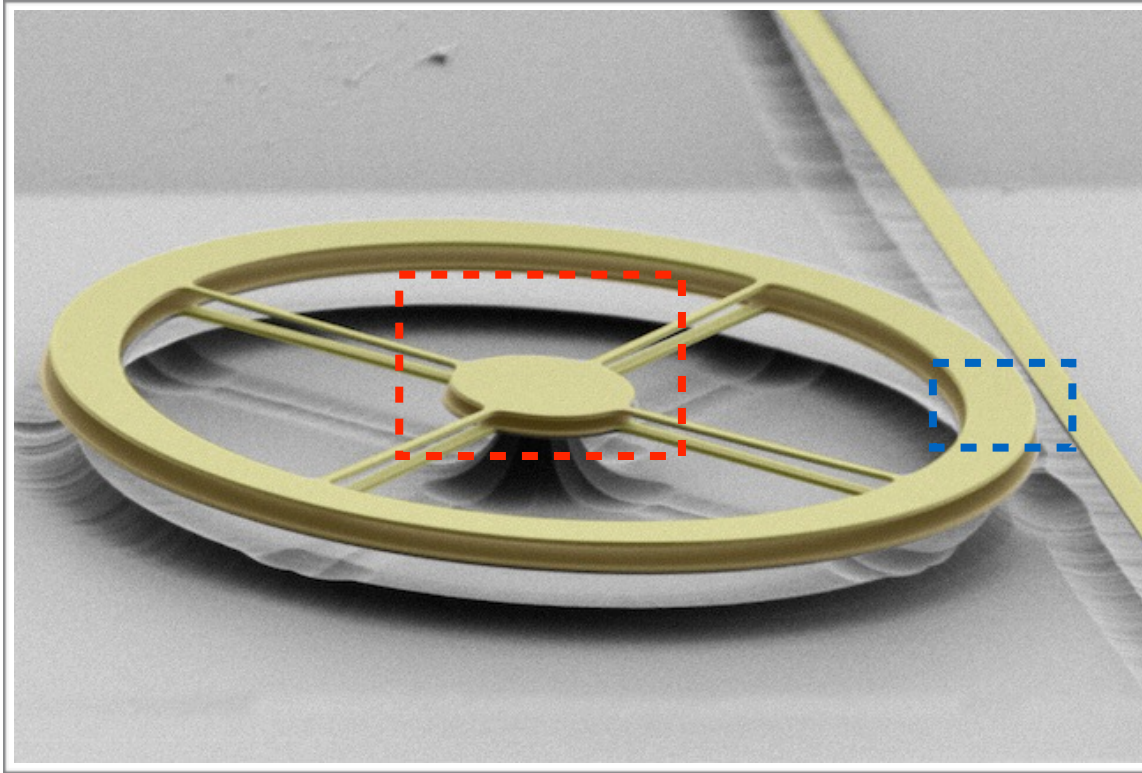


Controlling Cavities with Optical Forces



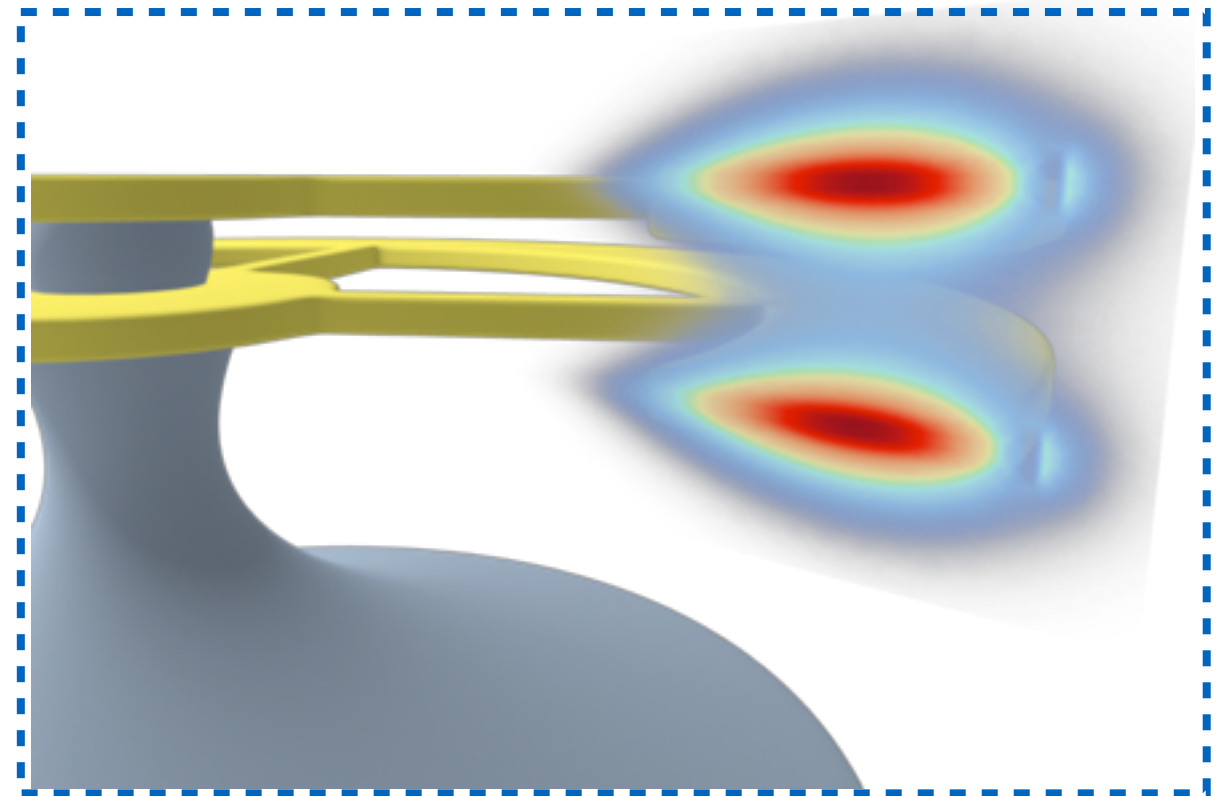
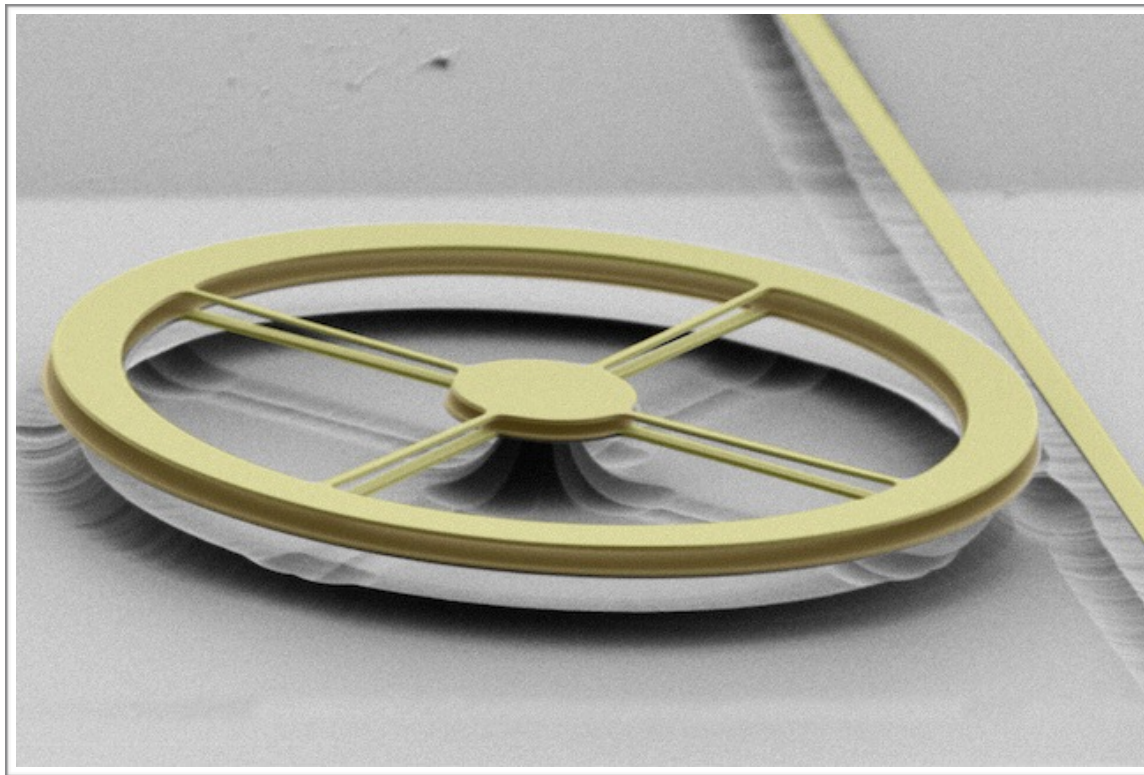


Controlling Cavities with Optical Forces

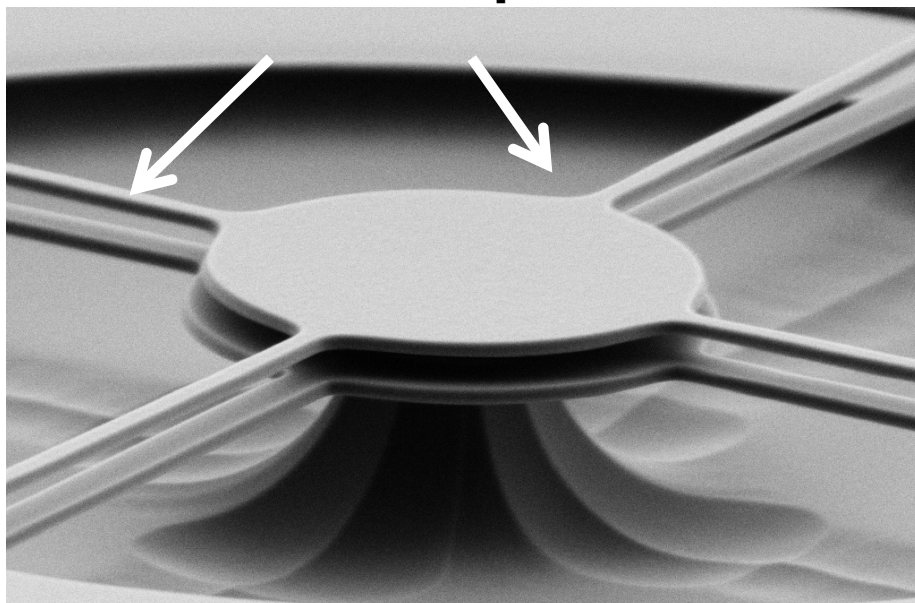




Controlling Cavities with Optical Forces

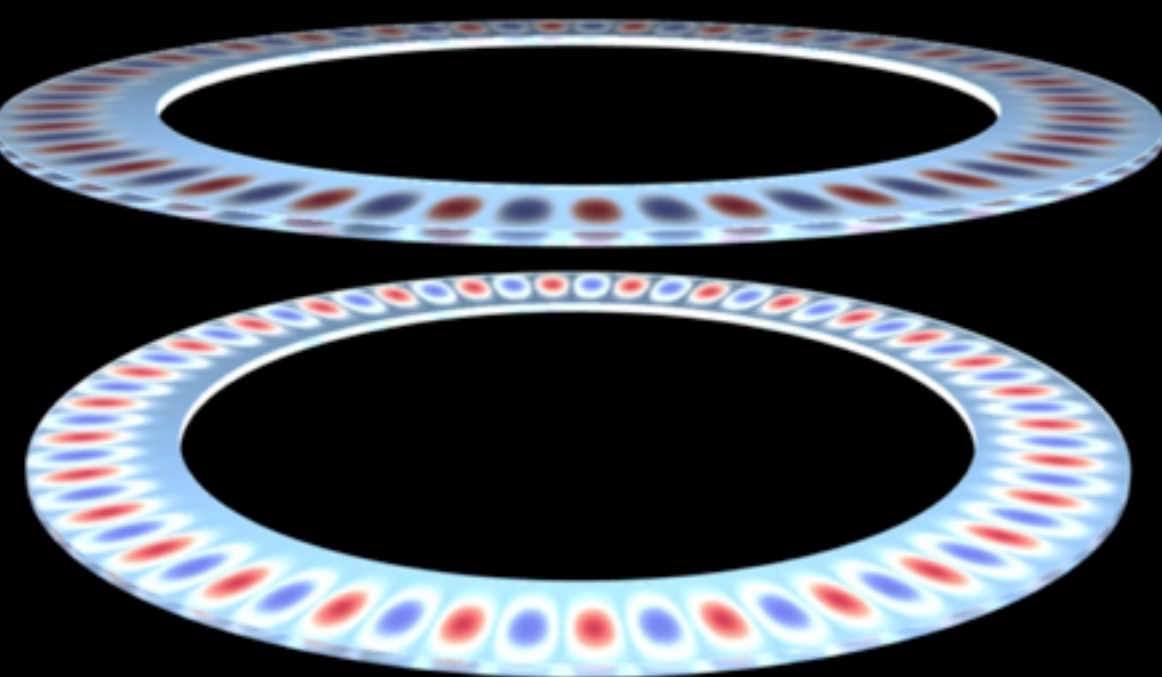


300 x 200 nm spokes



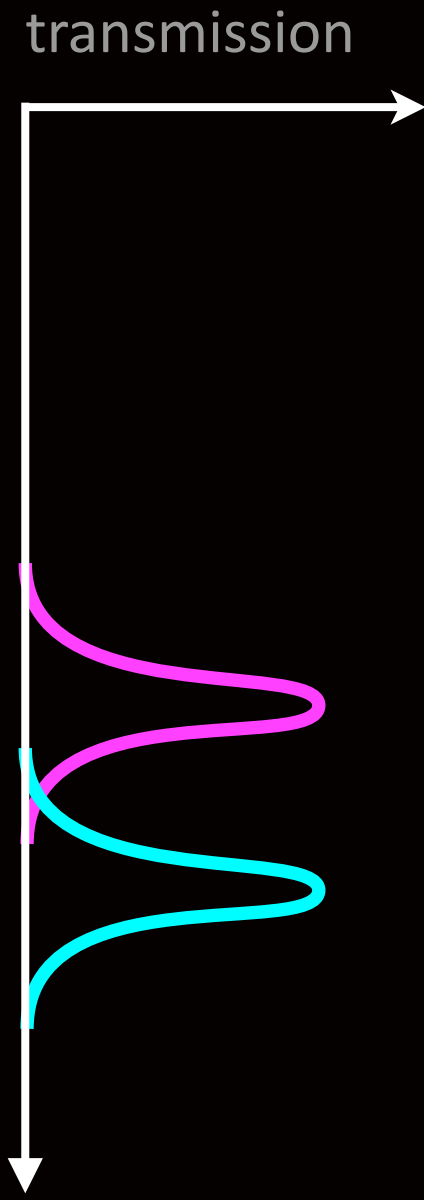
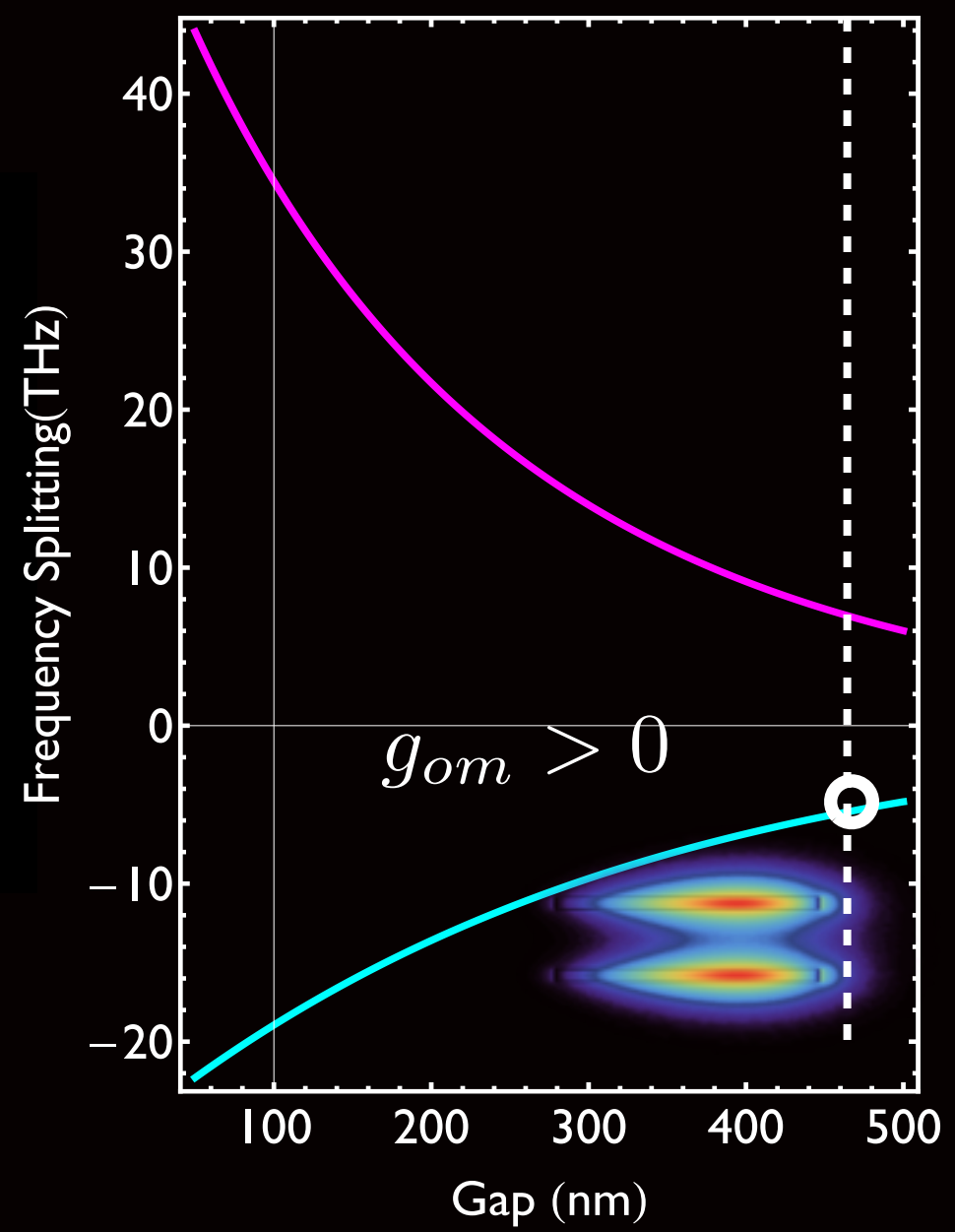


Attractive and Repulsive Forces



$$F \propto g_{om}$$

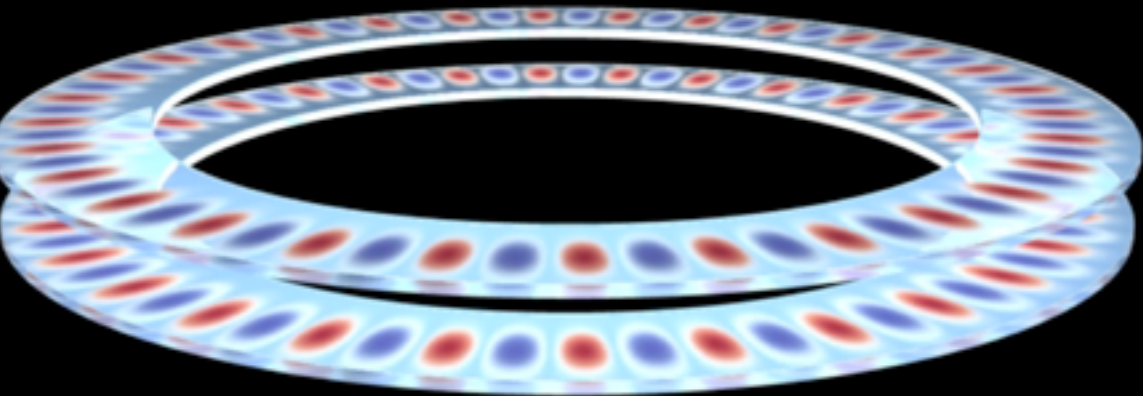
single-photon force



Wiederhecker et al, Nature 462 (2009)
Wiederhecker et al, OpEx 19, 2782 (2011)

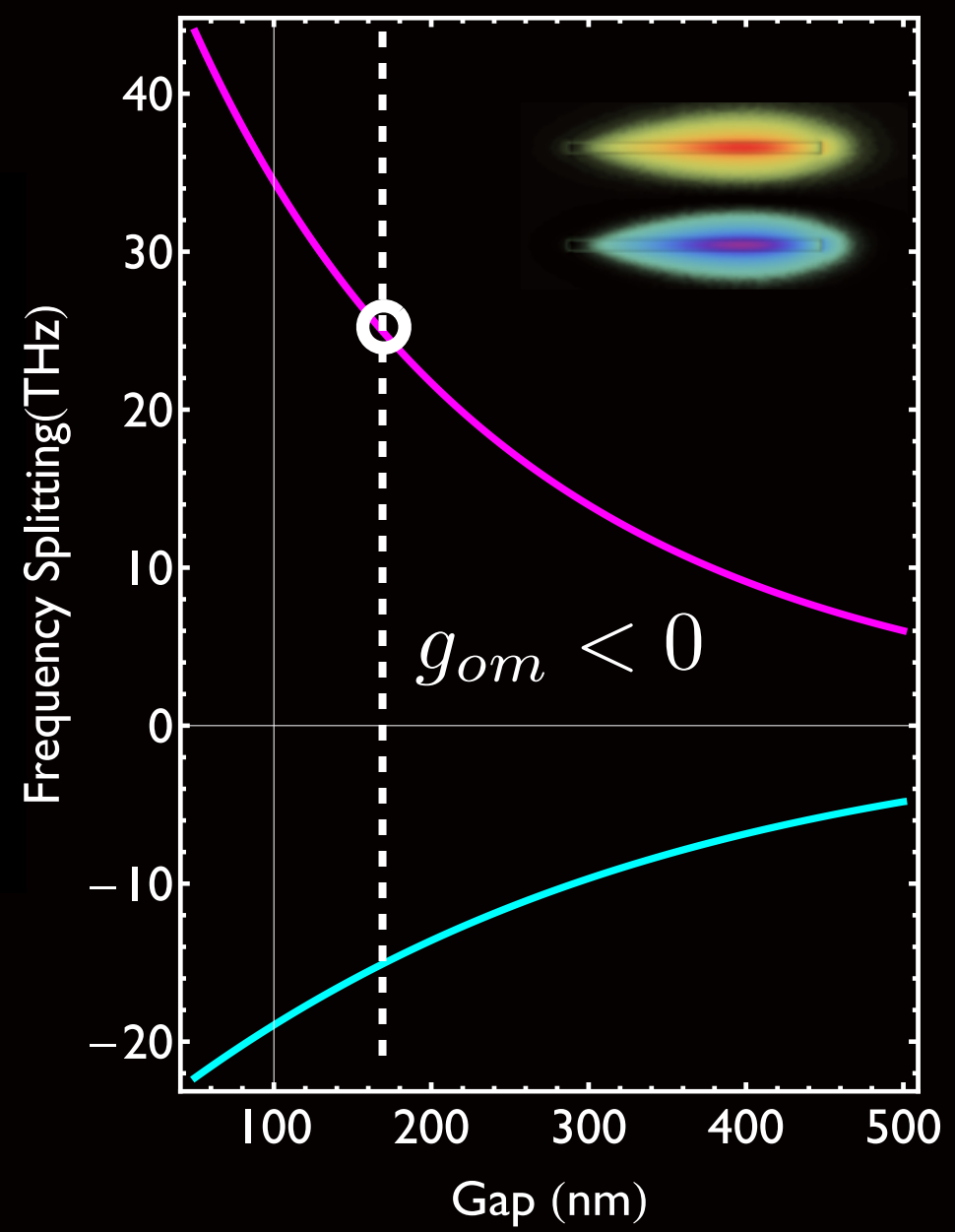


Attractive and Repulsive Forces

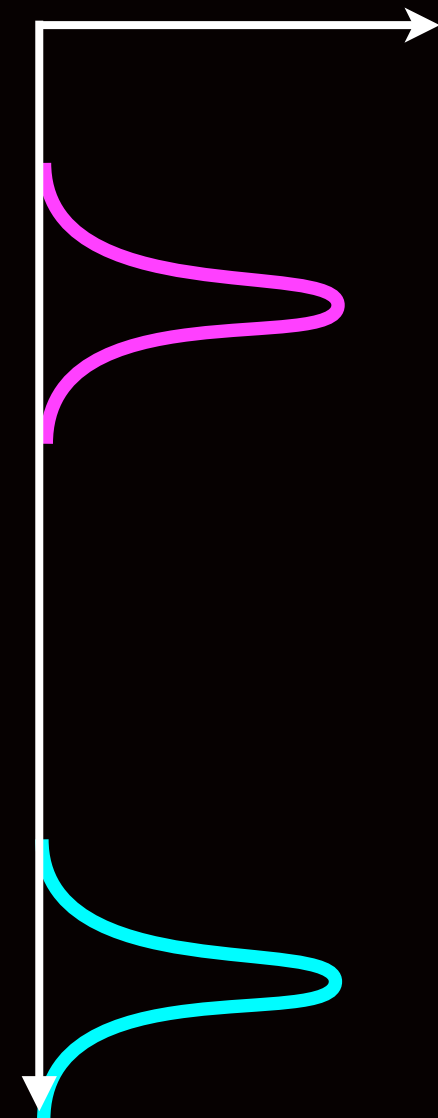


$$F \propto g_{\text{om}}$$

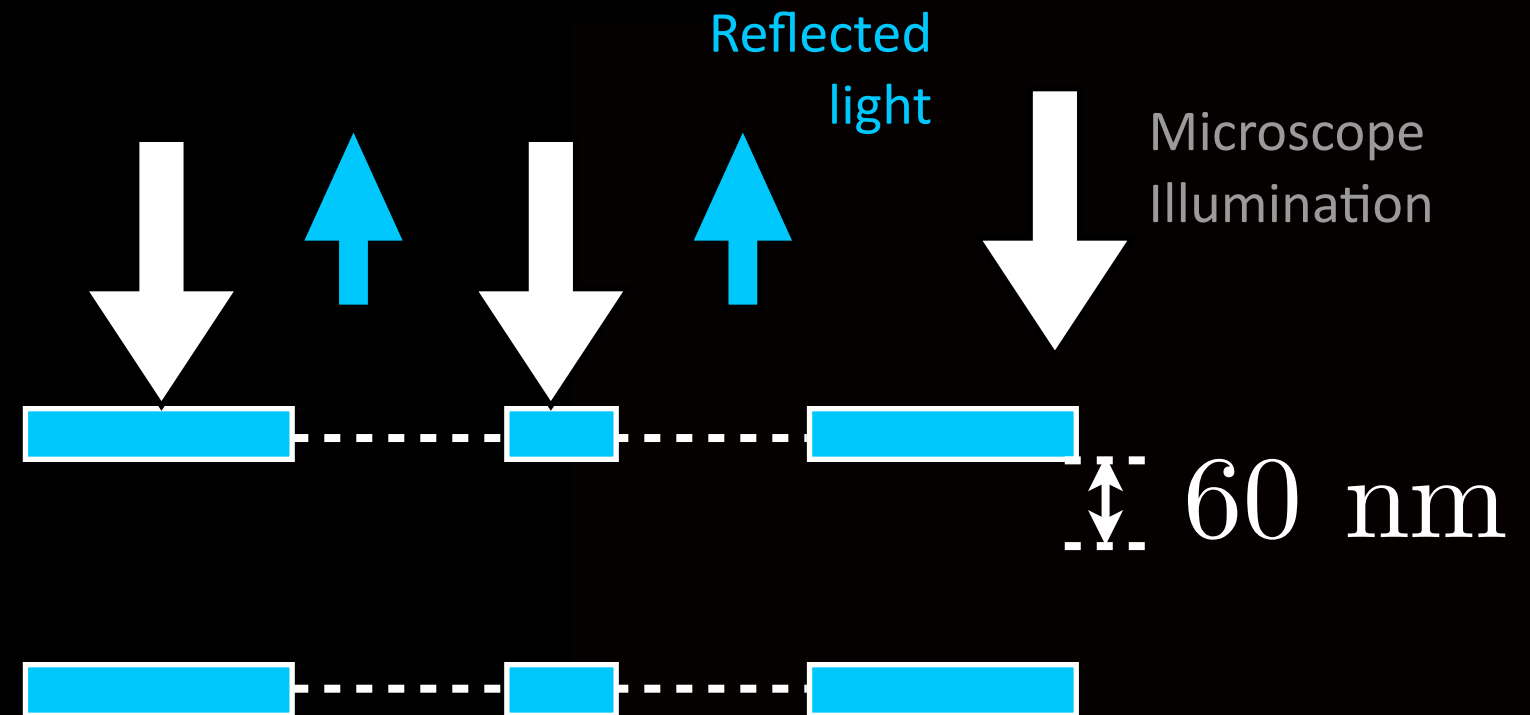
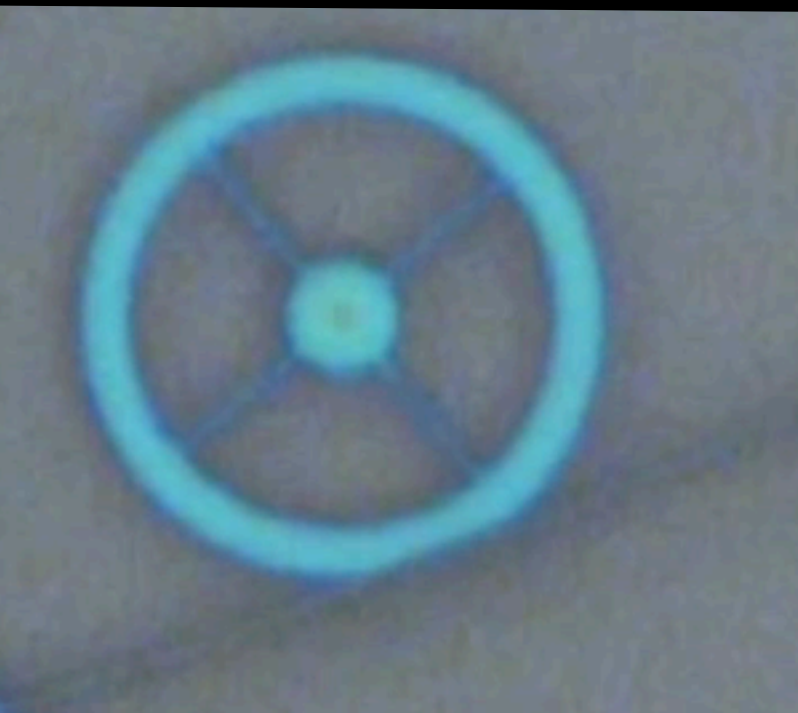
single-photon force



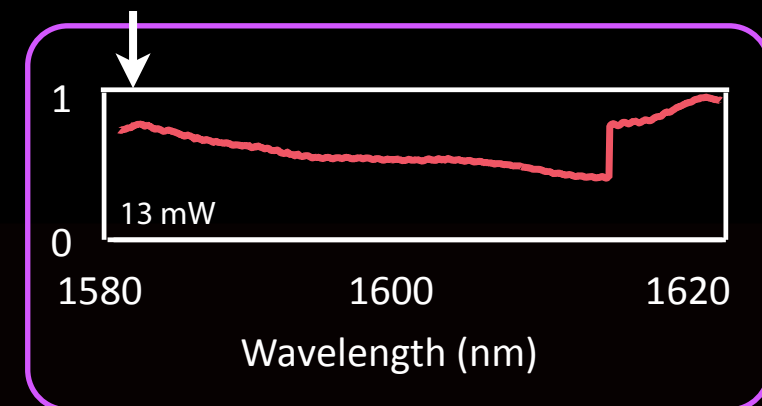
transmission



Wiederhecker et al, Nature 462 (2009)
Wiederhecker et al, OpEx 19, 2782 (2011)

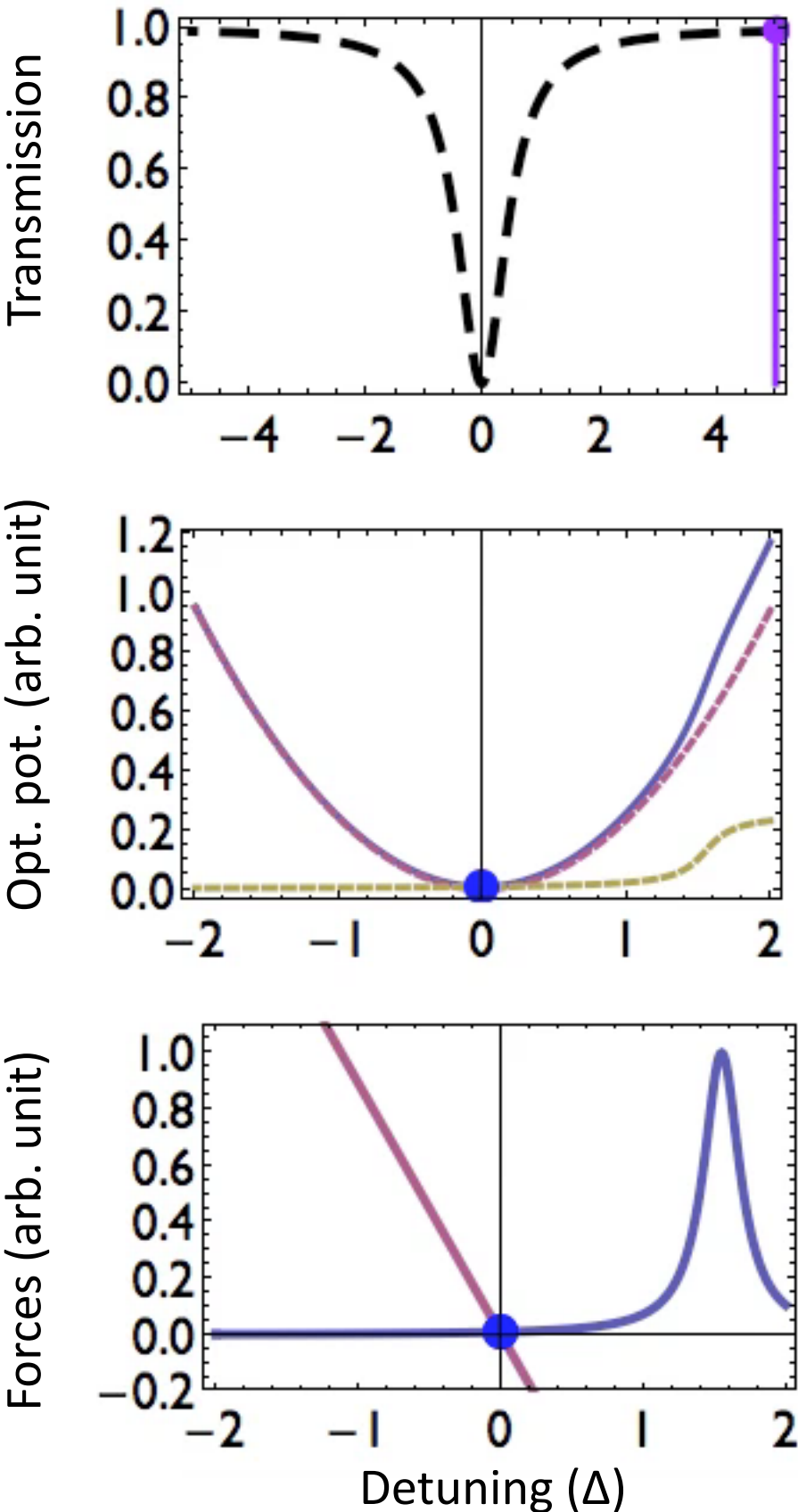


$$F_{\text{opt}} \approx 50 \text{ nN}$$





Static Bistability



Solve for displacement: cubic equation

$$\frac{m_{\text{eff}}\omega_0}{g_{\text{OM}}}x_0 = \frac{\alpha_{\text{in}}^2\kappa_e}{(\Delta - g_{\text{OM}}x_0)^2 + \kappa^2 / 4}$$

$$V \propto \int \frac{\alpha_{\text{in}}^2\kappa_e}{(\Delta - g_{\text{OM}}x)^2 + \kappa^2 / 4} dx$$

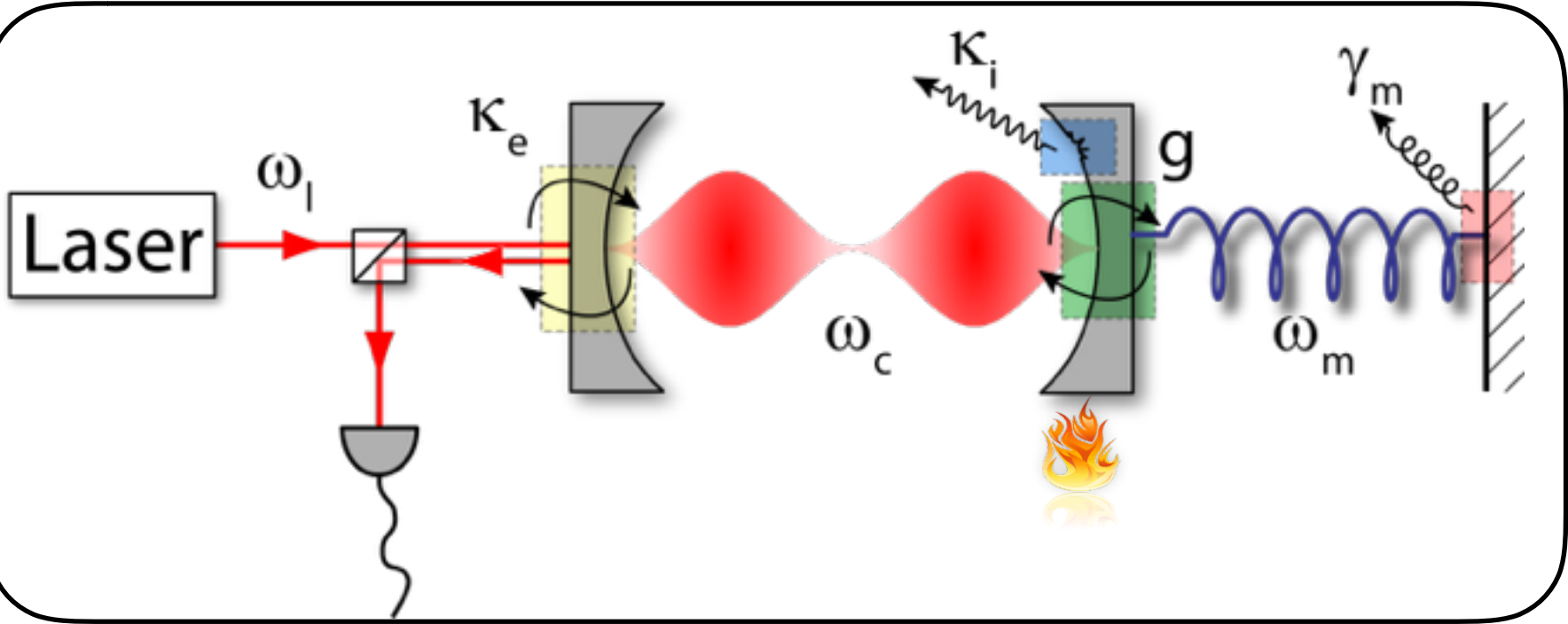


Outline

- ★ Optical and acoustic mode interaction
- ★ Optical force actuation
- ★ **Dynamical back-action**
- ★ Optomechanical clocks
- ★ Bullseye - a case study
- ★ Outlook



Dynamical Back-action



$m = 100 \text{ pg}$
 $\Omega / 2\pi = 1 \text{ MHz}$
 $\langle \delta x_T \rangle \approx 40 \text{ pm}$

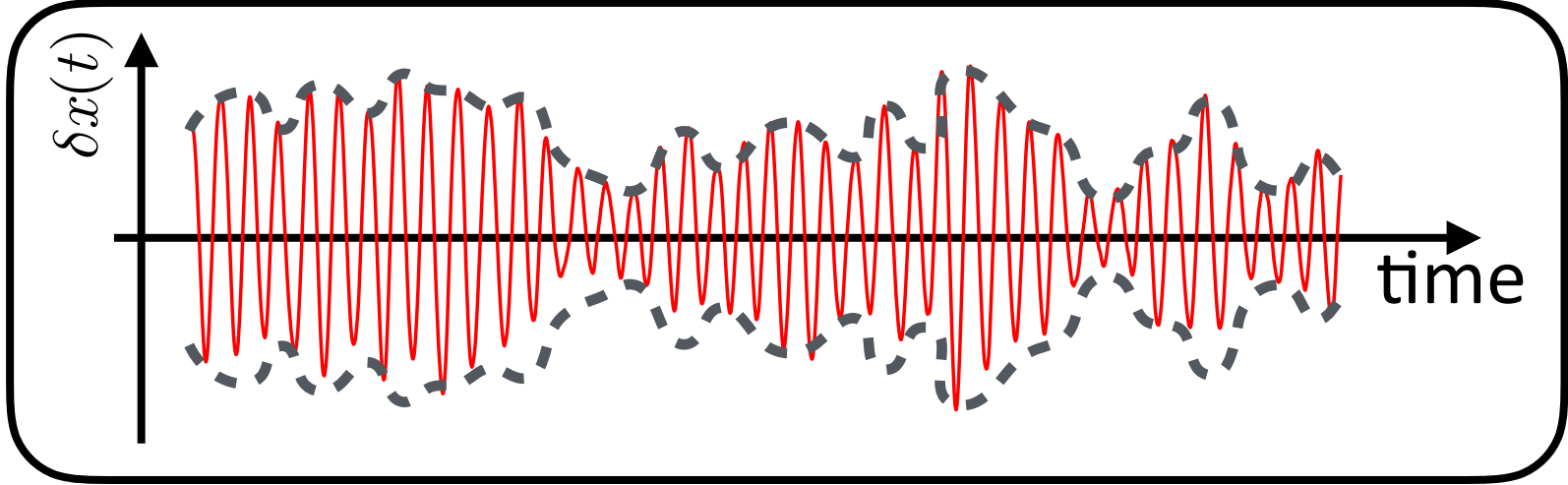
δF_{th}



Thermal noise

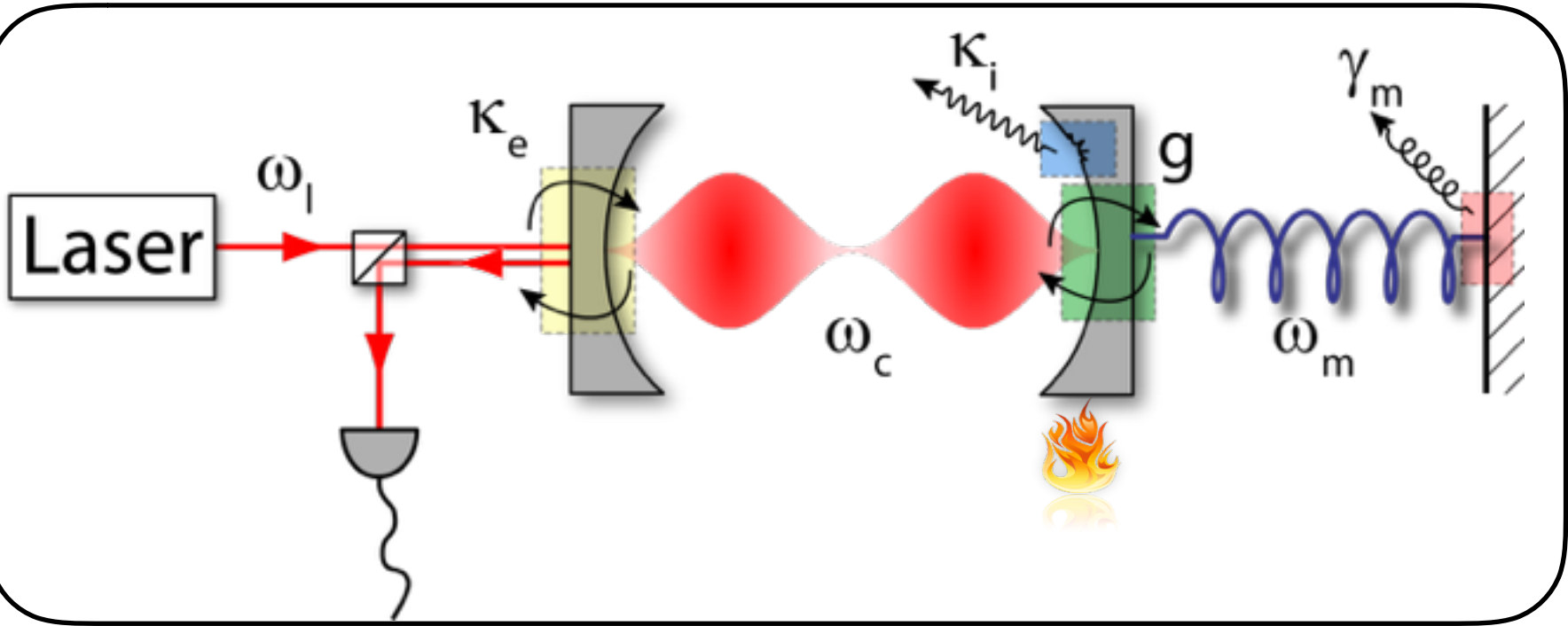
$\delta x(t)$

Displacement noise

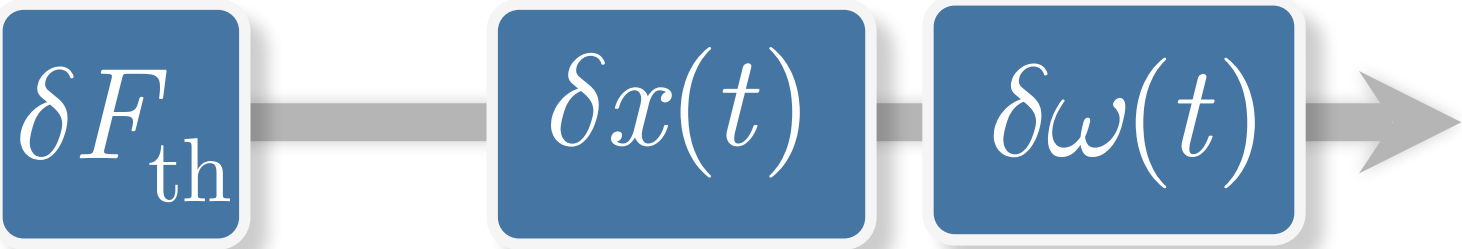




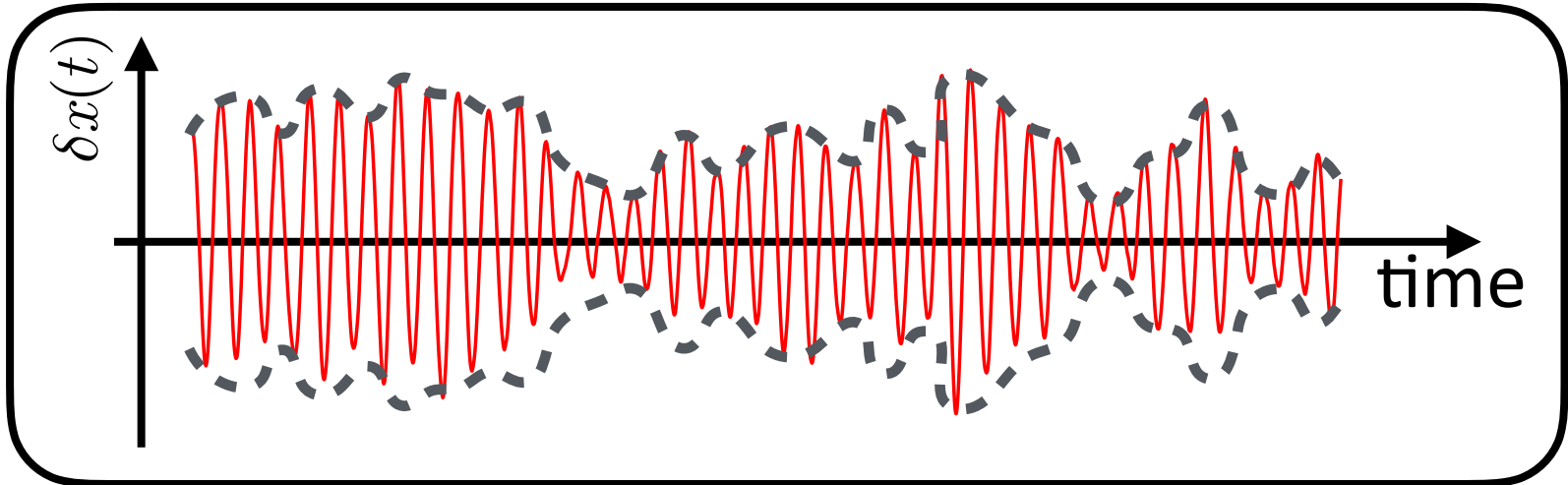
Dynamical Back-action



$\langle \delta x_T \rangle \approx 40 \text{ pm}$
 $g_{\text{om}} / 2\pi \approx 10^4 \text{ GHz/pm}$
 $\langle \delta \omega \rangle / 2\pi \approx 400 \text{ MHz}$

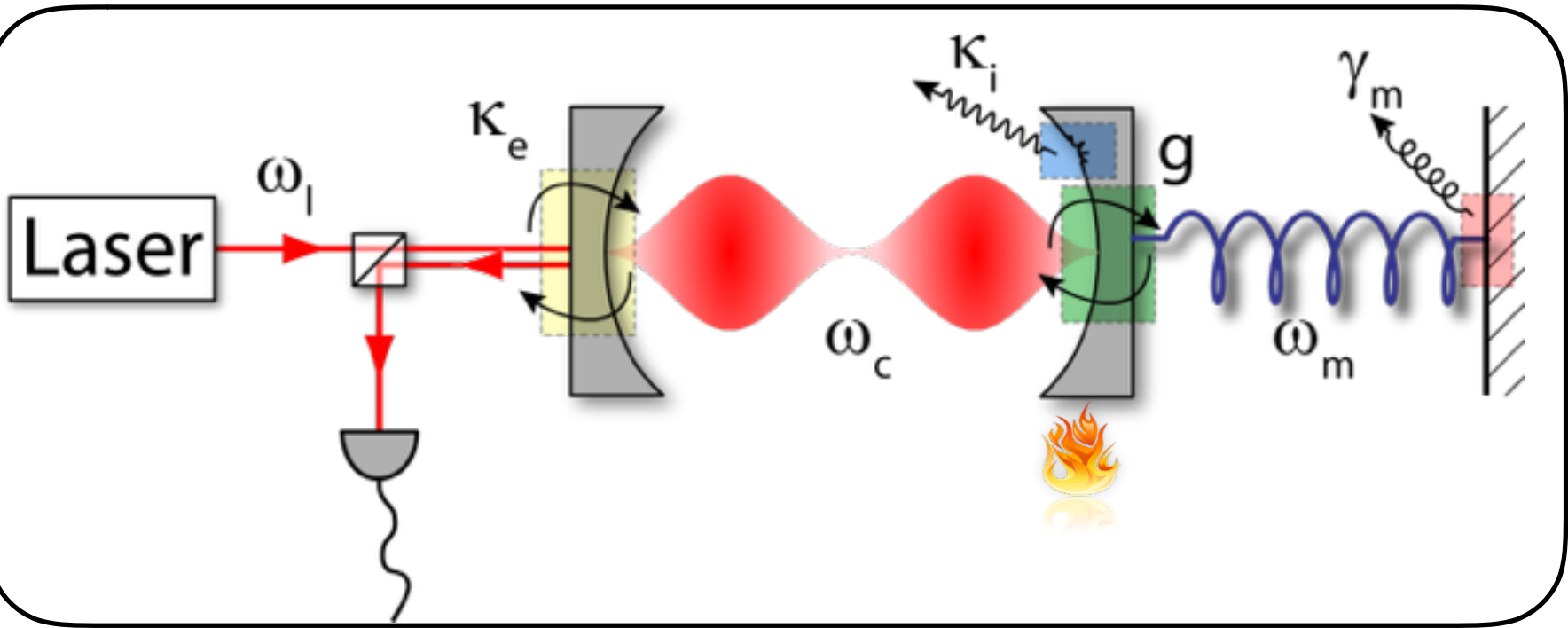


Thermal noise

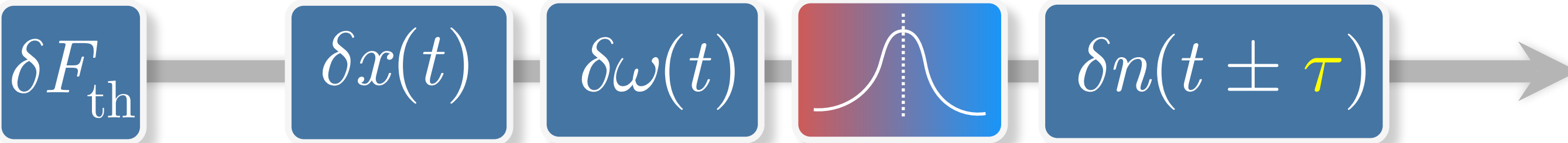




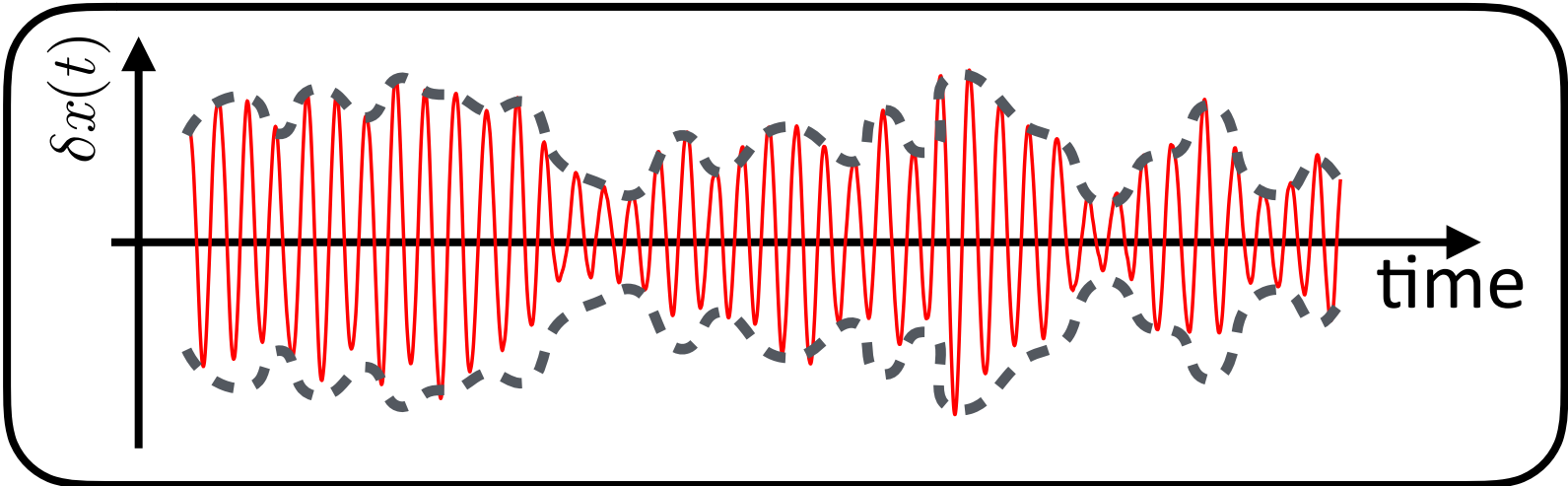
Dynamical Back-action



$\langle \delta x_T \rangle \approx 40 \text{ pm}$
 $g_{\text{om}} / 2\pi \approx 10^4 \text{ GHz/pm}$
 $\langle \delta \omega \rangle / 2\pi \approx 400 \text{ MHz}$

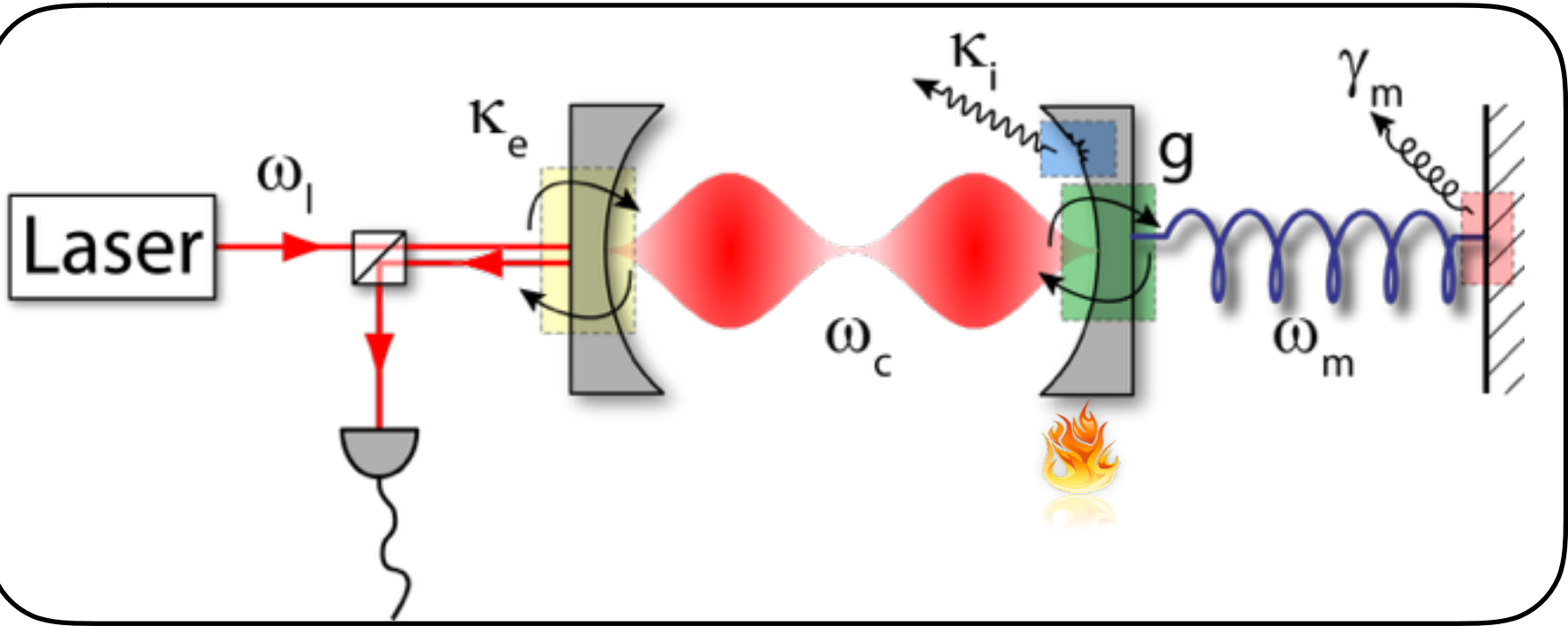


Thermal noise

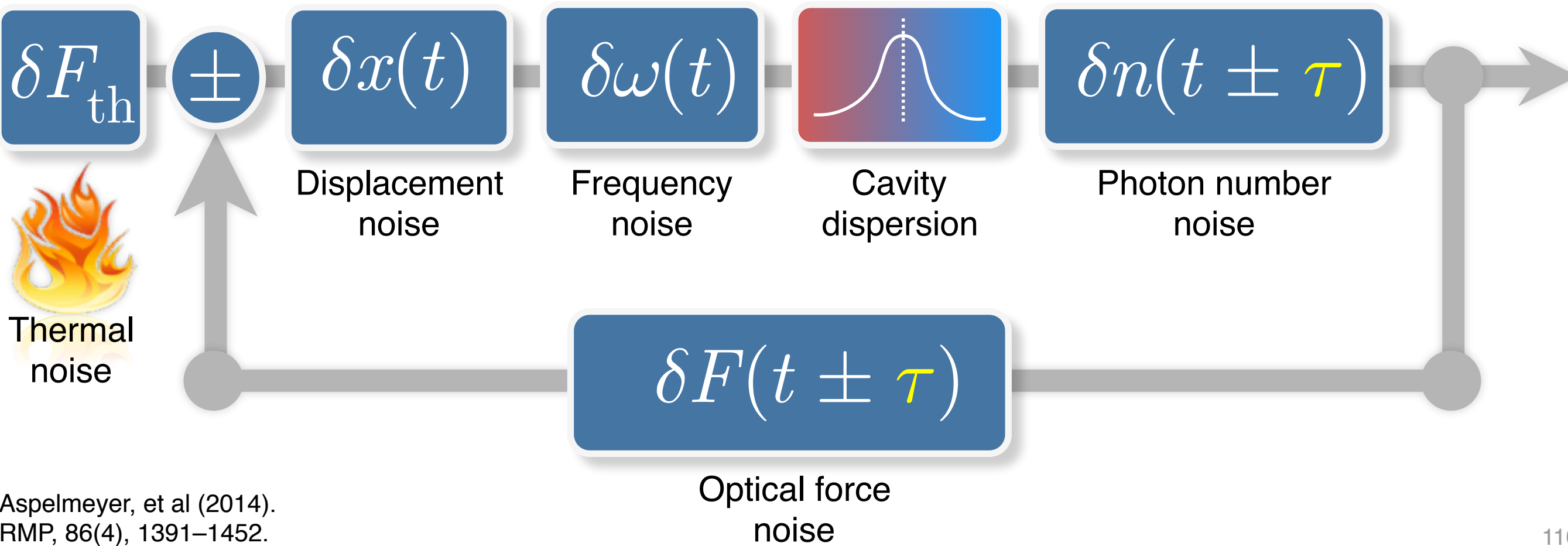




Dynamical Back-action

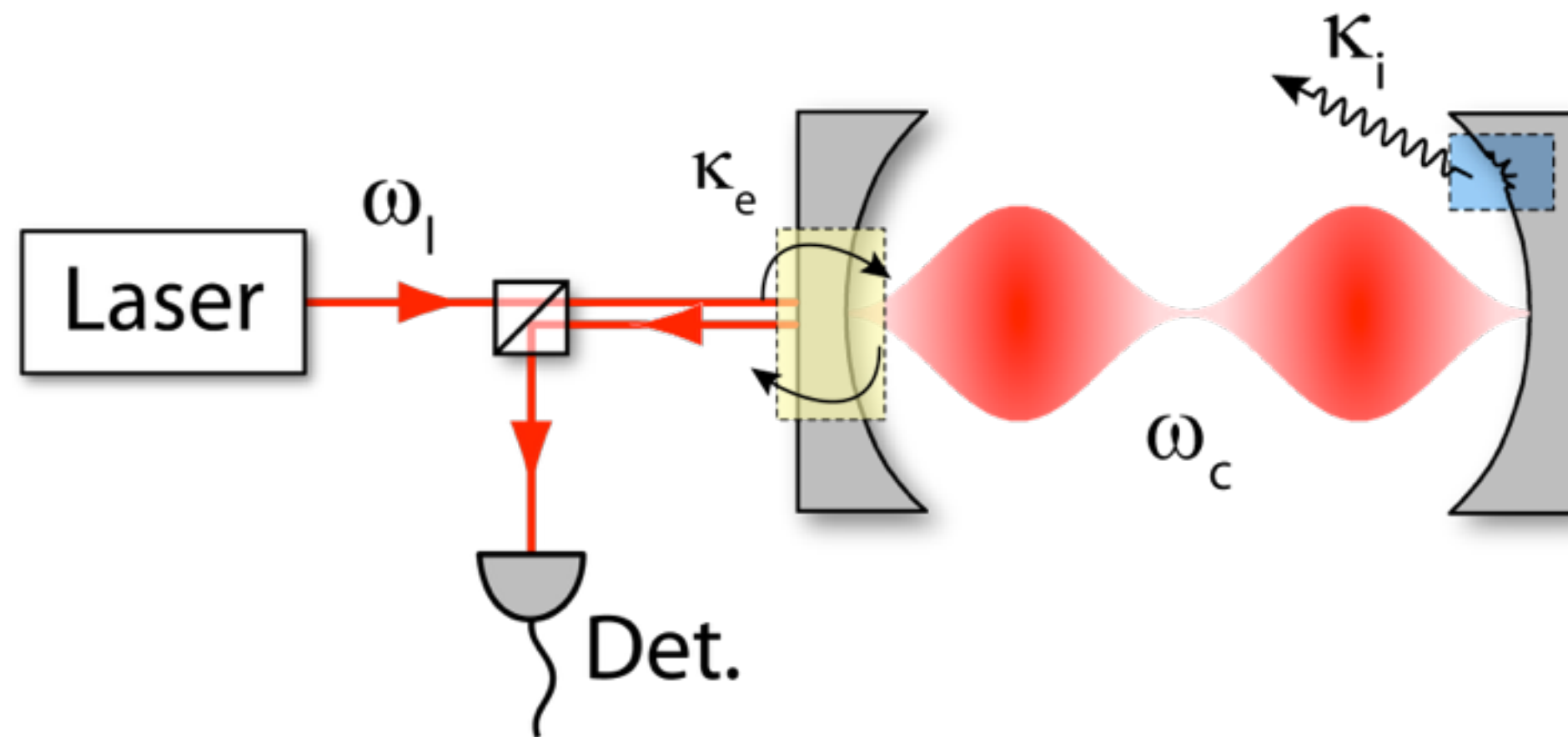


$\langle \delta x_T \rangle \approx 40 \text{ pm}$
 $g_{\text{om}} / 2\pi \approx 10^4 \text{ GHz/pm}$
 $\langle \delta \omega \rangle / 2\pi \approx 400 \text{ MHz}$





Optical Time Domain Equation

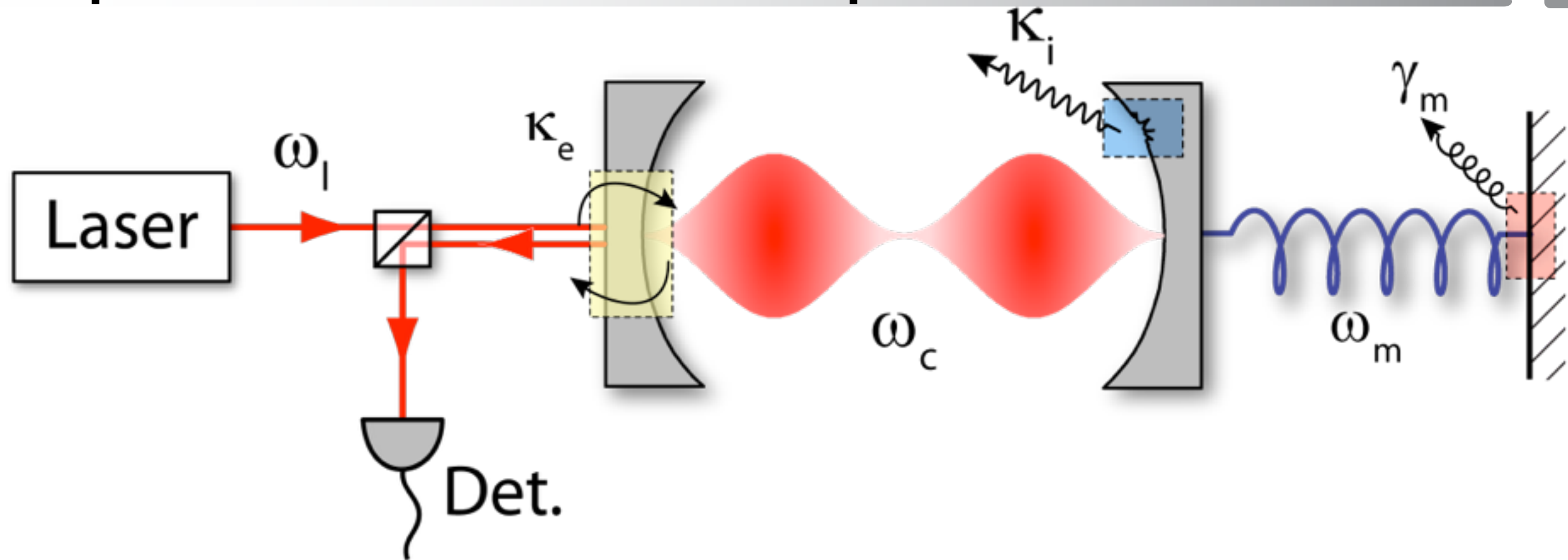


$$\dot{a}(t) = i(\omega_l - \omega_c)a - \frac{\kappa}{2}a + \sqrt{\kappa_e}a_{\text{in}}$$

Optical Amplitude Equation



Optical Time Domain Equation

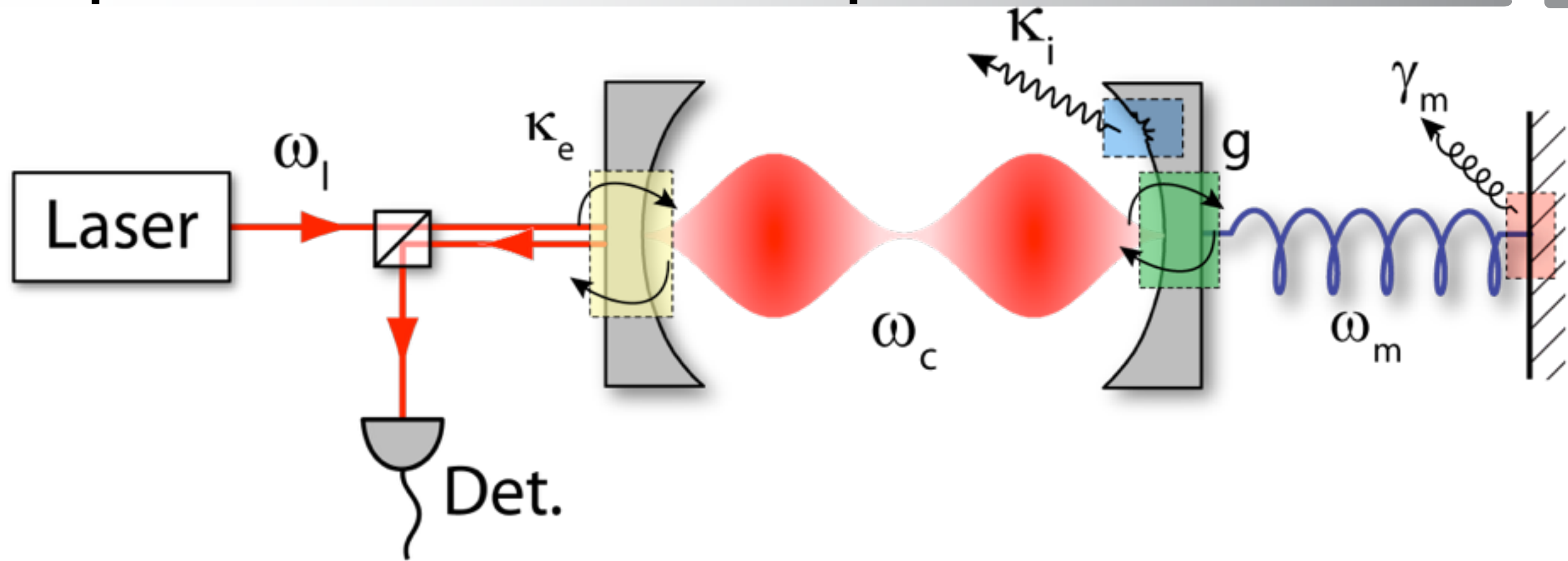


$$\dot{a}(t) = \left(i(\omega_l - \omega_c) - \frac{\kappa}{2} \right) a + \sqrt{\kappa_e} \alpha_{\text{in}}$$

$$\ddot{x} + \gamma_m \dot{x} + \omega_m^2 x = \frac{F_{\text{opt}}}{m}$$



Optical Time Domain Equation

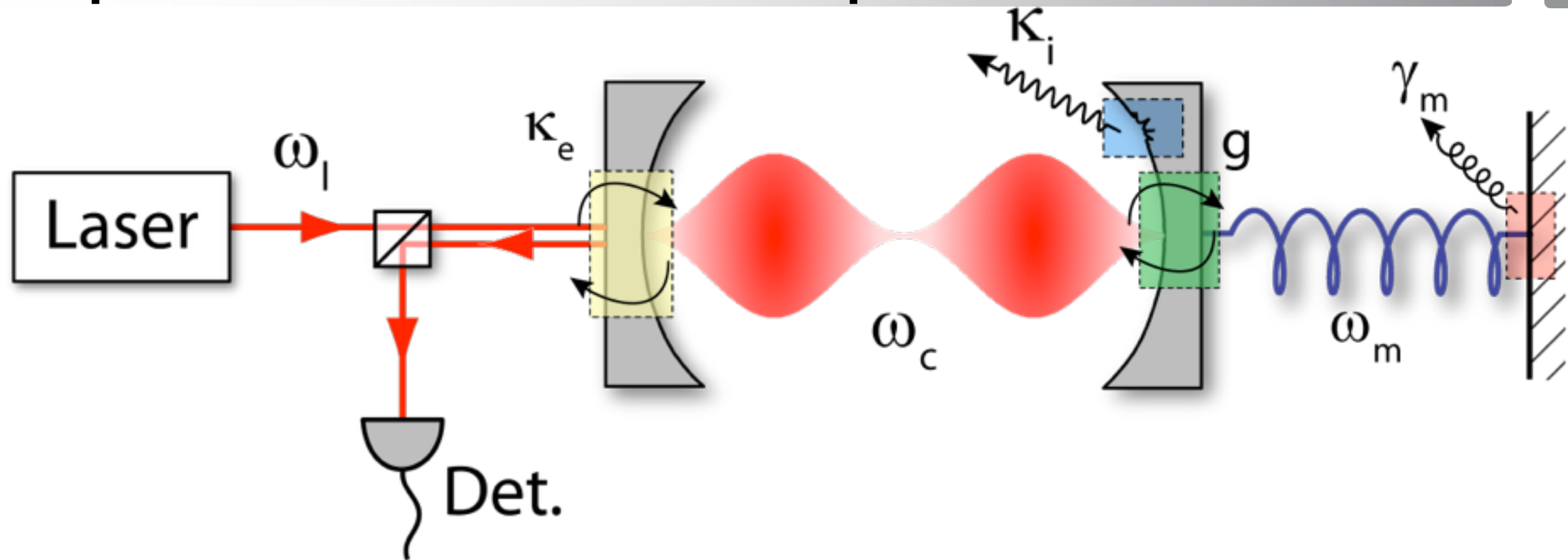


$$\dot{a}(t) = \left(i(\omega_l - \omega_c(x)) - \frac{\kappa}{2} \right) a + \sqrt{\kappa_e} \alpha_{\text{in}}$$

$$\ddot{x} + \gamma_m \dot{x} + \omega_m^2 x = \frac{F_{\text{opt}}(a)}{m}$$



Optical Time Domain Equation



$$\dot{a}(t) = \left(i(\Delta + g_{\text{OM}}x(t)) - \frac{\kappa}{2} \right) a + \sqrt{\kappa_e} \alpha_{\text{in}}$$

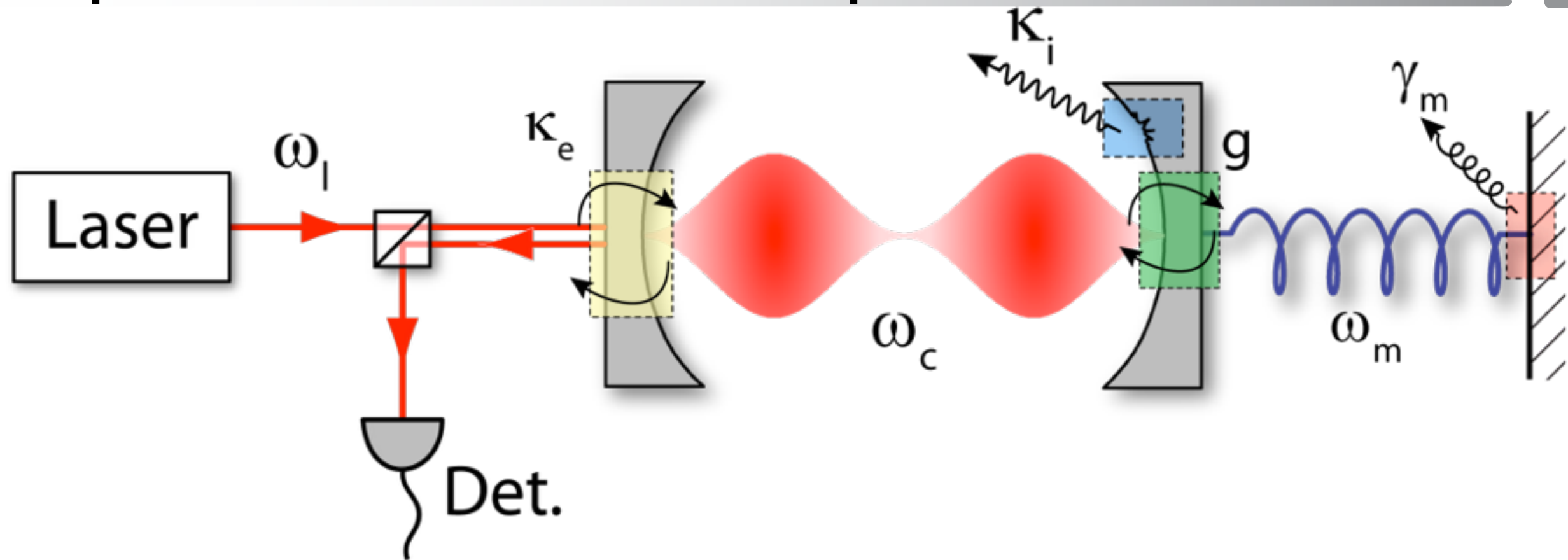
$$\ddot{x} + \gamma_m \dot{x} + \omega_m^2 x = \frac{g_{\text{OM}}}{m\omega_0} |a|^2$$

$$\omega_c = \omega_o(1 - g_{\text{OM}}x(t))$$

$$\Delta = \omega_l - \omega_o$$



Optical Time Domain Equation



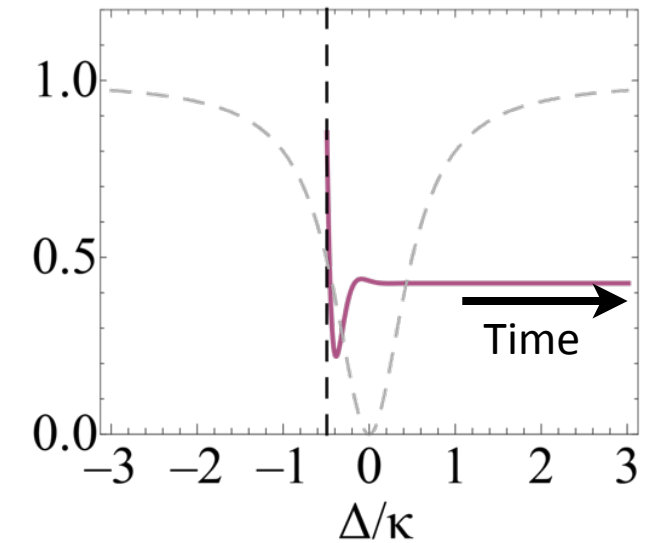
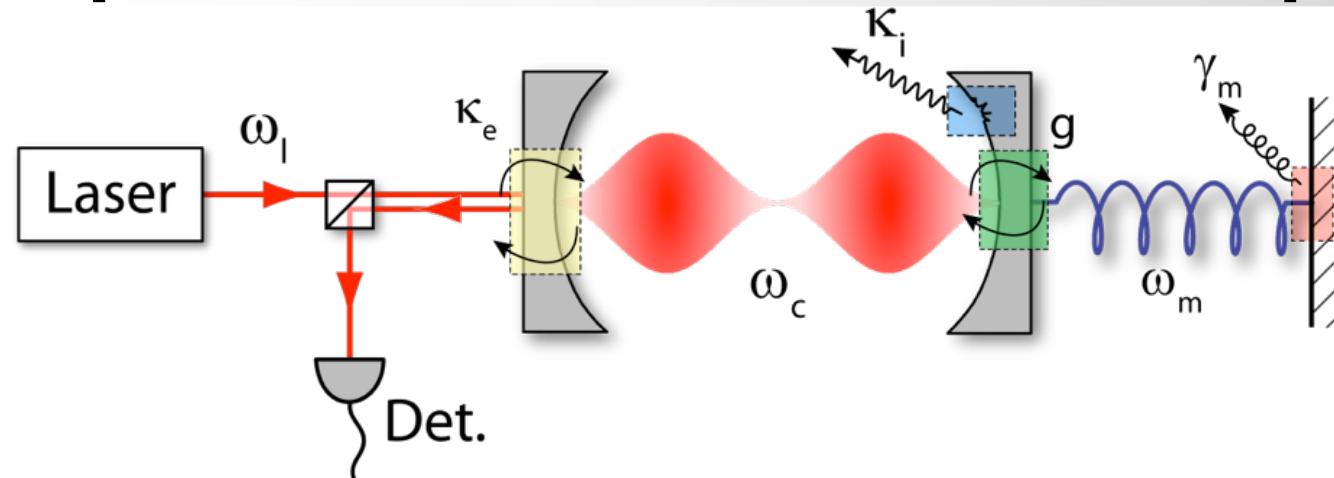
Coupled
Equations

$$\dot{a}(t) = \left(i\Delta - \frac{\kappa}{2} \right) a + i g_{\text{OM}} a x(t) + \sqrt{\kappa_e} \alpha_{\text{in}}$$

$$\ddot{x} + \gamma_m \dot{x} + \omega_m^2 x = \frac{|a(t)|^2 g_{\text{OM}}}{\omega_0 m}$$



Optical Time Domain Equation



$$\dot{a}(t) = \left(i\Delta - \frac{\kappa}{2} \right) a + ig_{\text{OM}} ax(t) + \sqrt{\kappa_e} \alpha_{\text{in}}$$

$$\ddot{x} + \gamma_m \dot{x} + \omega_m^2 x = \frac{g_{\text{OM}}}{\omega_0 m} |a(t)|^2$$

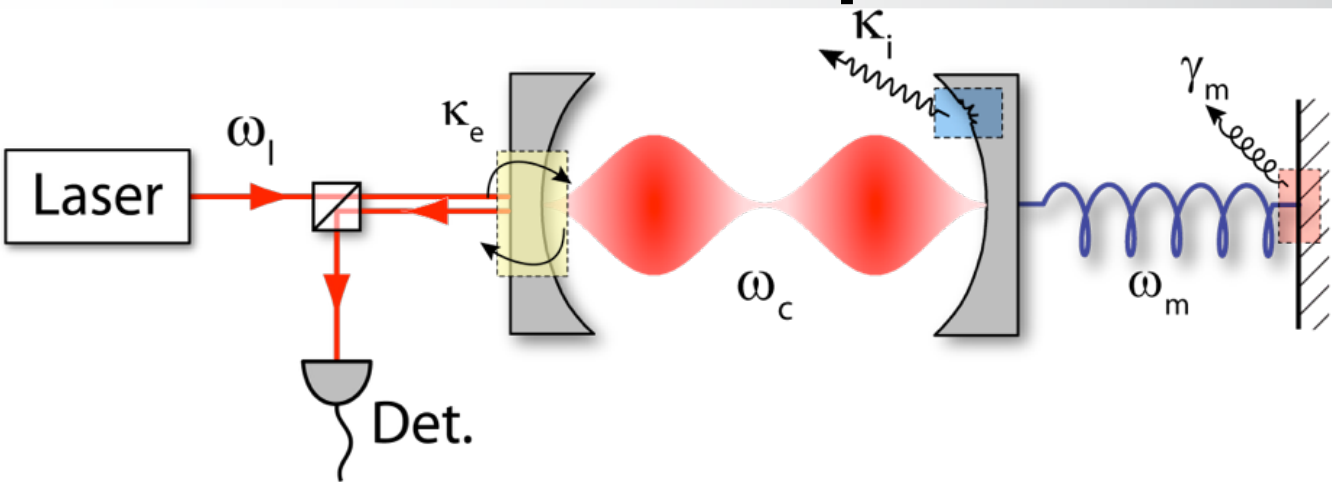
Linearized dynamics:

$$a(t) \approx a_0 + \delta a(t)$$

$$x(t) \approx x_0 + \delta x(t)$$



Zeroth order equation: static shift



$$a(t) \approx a_0 + \delta a(t)$$
$$x(t) \approx x_0 + \delta x(t)$$

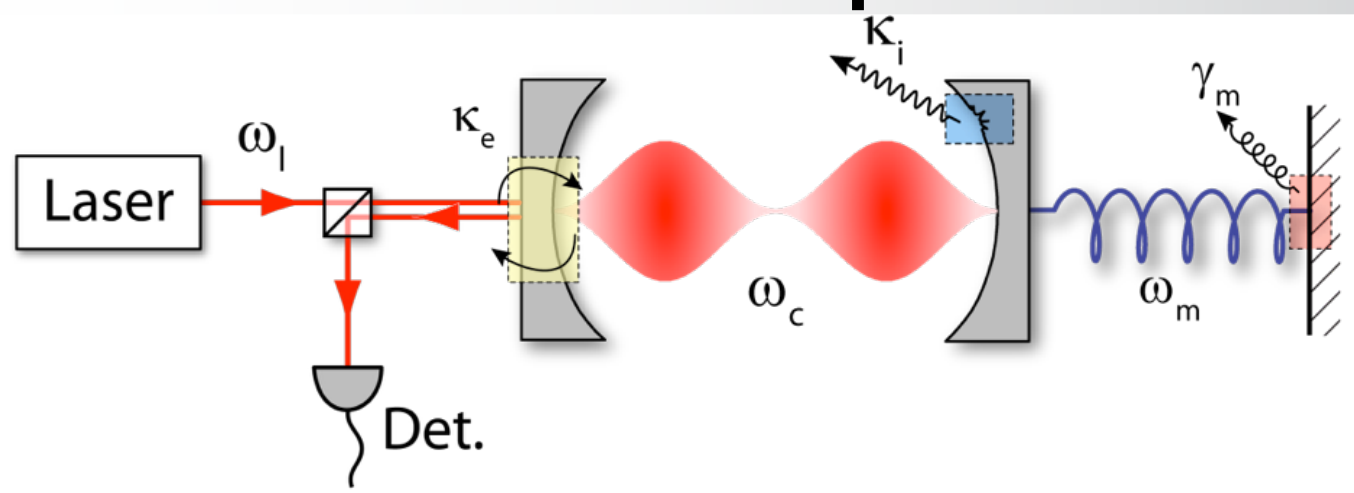
$$\dot{a}(t) = \left(i\Delta - \frac{\kappa}{2} \right) a + ig_{\text{OM}} ax(t) + \sqrt{\kappa_e} \alpha_{\text{in}}$$

$$\ddot{x} + \gamma_m \dot{x} + \omega_m^2 x = \frac{\hbar g_{\text{OM}}}{m_{\text{eff}}} |a(t)|^2$$

The zeroth order solution is basically the exact static solution of the problem!



Zeroth order equation: static shift



$$a(t) \approx a_0 + \delta a(t)$$
$$x(t) \approx x_0 + \delta x(t)$$

$$\dot{a}(t) = \left(i\Delta - \frac{\kappa}{2} \right) a + ig_{OM}ax(t) + \sqrt{\kappa_e}\alpha_{in}$$
$$\ddot{x} + \gamma_m\dot{x} + \omega_m^2x = \frac{\hbar g_{OM}}{m_{eff}}|a(t)|^2$$

The zeroth order solution is basically the exact static solution of the problem!

Average number of photons

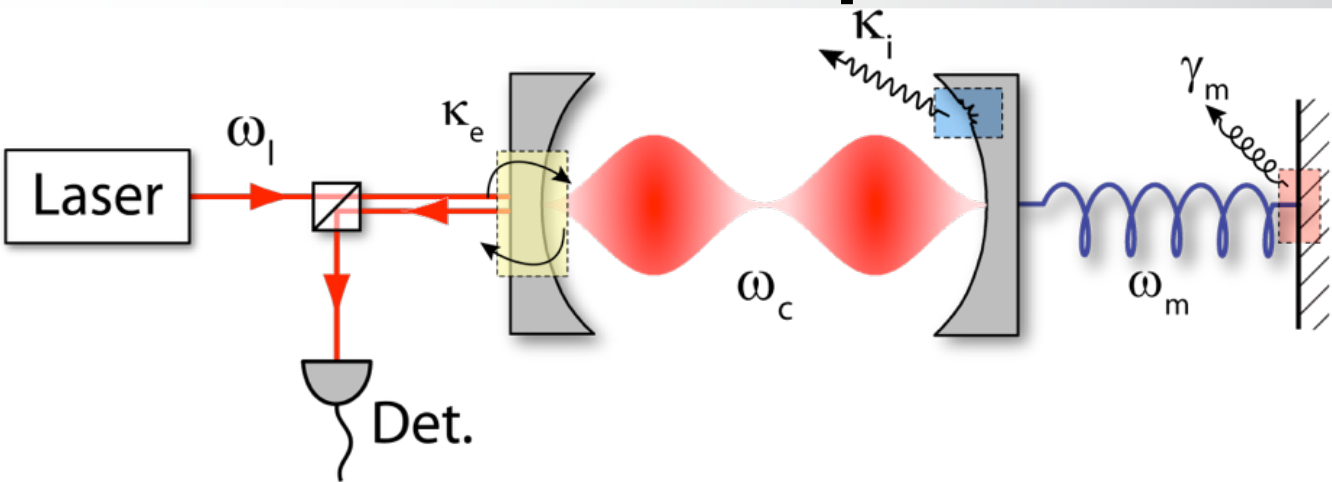
$$U \equiv |a_0|^2 = \frac{\alpha_{in}^2 \kappa_e}{\Delta'^2 + \kappa^2 / 4}$$

Static displacement

$$x_0 = \frac{g_{OM}}{\omega_0 m_{eff}} U$$



Zeroth order equation: static shift



$$\begin{aligned} a(t) &\approx a_0 + \delta a(t) \\ x(t) &\approx x_0 + \delta x(t) \end{aligned}$$

Average number of photons

$$U \equiv |a_0|^2 = \frac{\alpha_{in}^2 \kappa_e}{\Delta'^2 + \kappa^2 / 4}$$

Static displacement

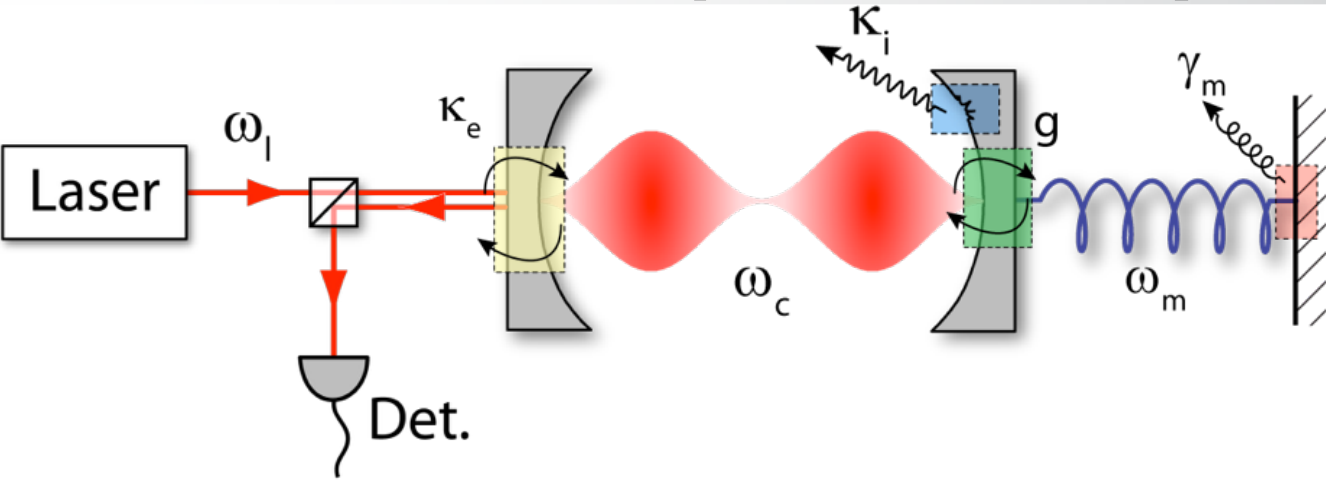
$$x_0 = \frac{g_{OM}}{\omega_0 m_{eff}} U$$

Not so easy buddy!
The detuning (Δ) is a function of position (x) !

$$\Delta' = \Delta - g_{OM} x_0$$



First order equation: dynamics

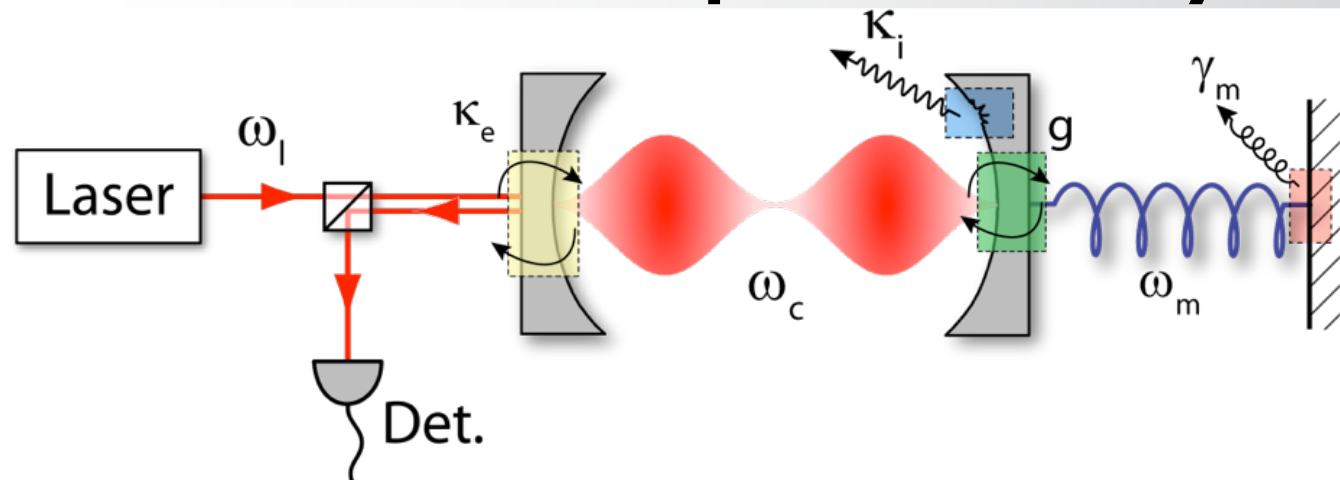


$$\begin{aligned} a(t) &\approx a_0 + \delta a(t) \\ x(t) &\approx x_0 + \delta x(t) \end{aligned}$$

$$F(t) \propto |a_0 + \delta a|^2 = \frac{g_{\text{OM}}}{\omega_0} \left(a_0^2 + \underbrace{a_0 \delta a^* + a_0^* \delta a}_{\delta F(t)} + \delta a^2 \right)$$



First order equation: dynamics



$$a(t) \approx a_0 + \delta a(t)$$

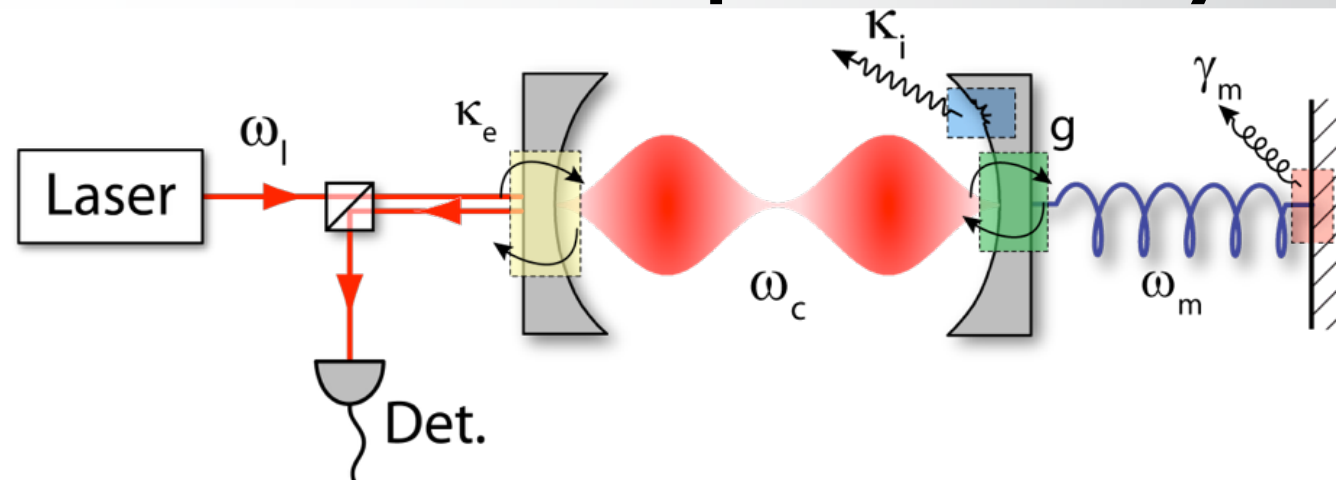
$$x(t) \approx x_0 + \delta x(t)$$

$$F(t) \propto |a_0 + \delta a|^2 = \frac{g_{OM}}{\omega_0} (a_0^2 + \underbrace{a_0 \delta a^* + a_0^* \delta a}_{\delta F(t)} + \delta a^2)$$

$$\delta \ddot{x} = -\Omega^2 \delta x - \gamma \delta \dot{x} - \frac{g_{OM}}{\omega_0 m_{\text{eff}}} (a_0 \delta a^* + a_0^* \delta a)$$



First order equation: dynamics



$$a(t) \approx a_0 + \delta a(t)$$

$$x(t) \approx x_0 + \delta x(t)$$

$$F(t) \propto |a_0 + \delta a|^2 = \frac{g_{OM}}{\omega_0} (a_0^2 + \underbrace{a_0 \delta a^* + a_0^* \delta a}_{\delta F(t)} + \delta a^2)$$

$$\delta \ddot{x} = -\Omega^2 \delta x - \gamma \delta \dot{x} - \frac{g_{OM}}{\omega_0 m_{\text{eff}}} (a_0 \delta a^* + a_0^* \delta a)$$

$$\delta \ddot{x} = -\Omega^2 \delta x - \gamma \delta \dot{x} + \delta F(t) / m_{\text{eff}}$$



First order equation: dynamics

$$\delta\ddot{x} = -\Omega^2\delta x - \gamma\delta\dot{x} + \delta F(t)$$

$$\delta\dot{a} = (i\Delta' - \frac{\kappa}{2})\delta a - ig_{\text{OM}}a_0\delta x$$

Fourier transform: $\tilde{f}[\omega] = \int_{-\infty}^{\infty} f(t)e^{i\omega t} dt$

$$m_{\text{eff}}(\Omega^2 - \omega^2 + i\gamma\omega)\delta\tilde{x}[\omega] = \delta\tilde{F}[\omega]$$

$$\delta\tilde{a}(\omega) = \frac{-ig_{\text{OM}}a_0}{-\omega - (i\Delta - \frac{\kappa}{2})}\delta\tilde{x}(\omega)$$



First order equation: dynamics

$$\delta \tilde{a}(\omega) = \frac{-i g_{\text{OM}} a_0}{-\omega - (i\Delta - \frac{\kappa}{2})} \delta \tilde{x}(\omega)$$

$$m_{\text{eff}}(\Omega^2 - \omega^2 + i\gamma\omega)\delta \tilde{x}[\omega] = \delta \tilde{F}[\omega]$$

$$\delta \tilde{F}[\omega] = \delta \tilde{x} \frac{g_{\text{OM}}^2 U}{\omega_0} \left(\frac{1}{(\Delta' - \omega) - i\kappa / 2} + \frac{1}{(\Delta' + \omega) + i\kappa / 2} \right) \equiv \Sigma(\omega)$$



First order equation: dynamics

$$\delta \tilde{F}[\omega] = \delta \tilde{x} \frac{g_{\text{OM}}^2 U}{\omega_0} \left(\frac{1}{(\Delta' - \omega) - i\kappa / 2} + \frac{1}{(\Delta' + \omega) + i\kappa / 2} \right) \equiv \Sigma(\omega)$$

$$m_{\text{eff}}(\Omega^2 - \omega^2 + i\gamma\omega)\delta \tilde{x}[\omega] = 2g^2 m_{\text{eff}} \Omega \Sigma(\omega) \delta \tilde{x}[\omega]$$
$$\equiv \chi_{xx}^{-1}(\omega) \qquad g \equiv (g_{\text{OM}} x_{\text{zpf}}) \sqrt{n_c}$$

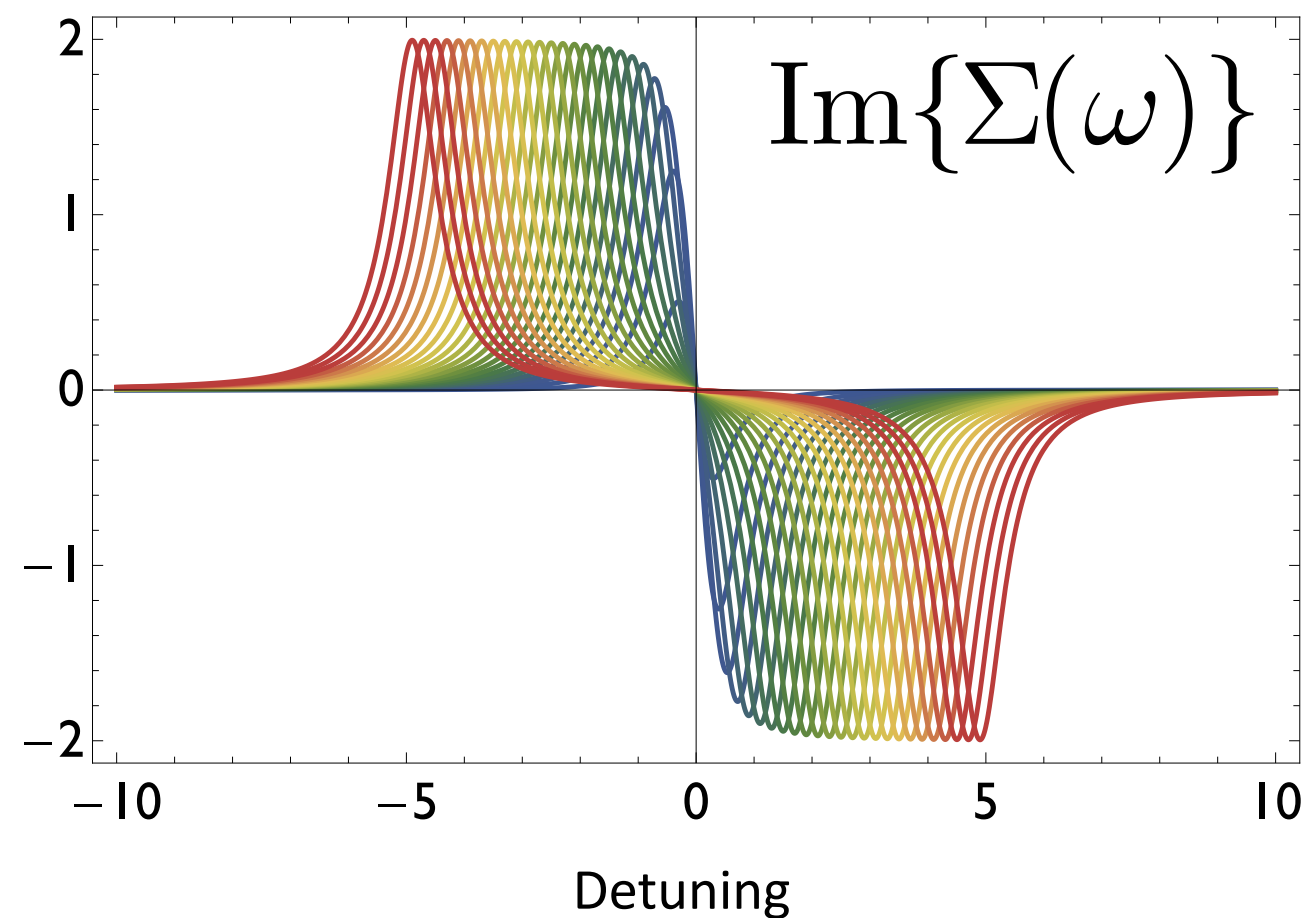
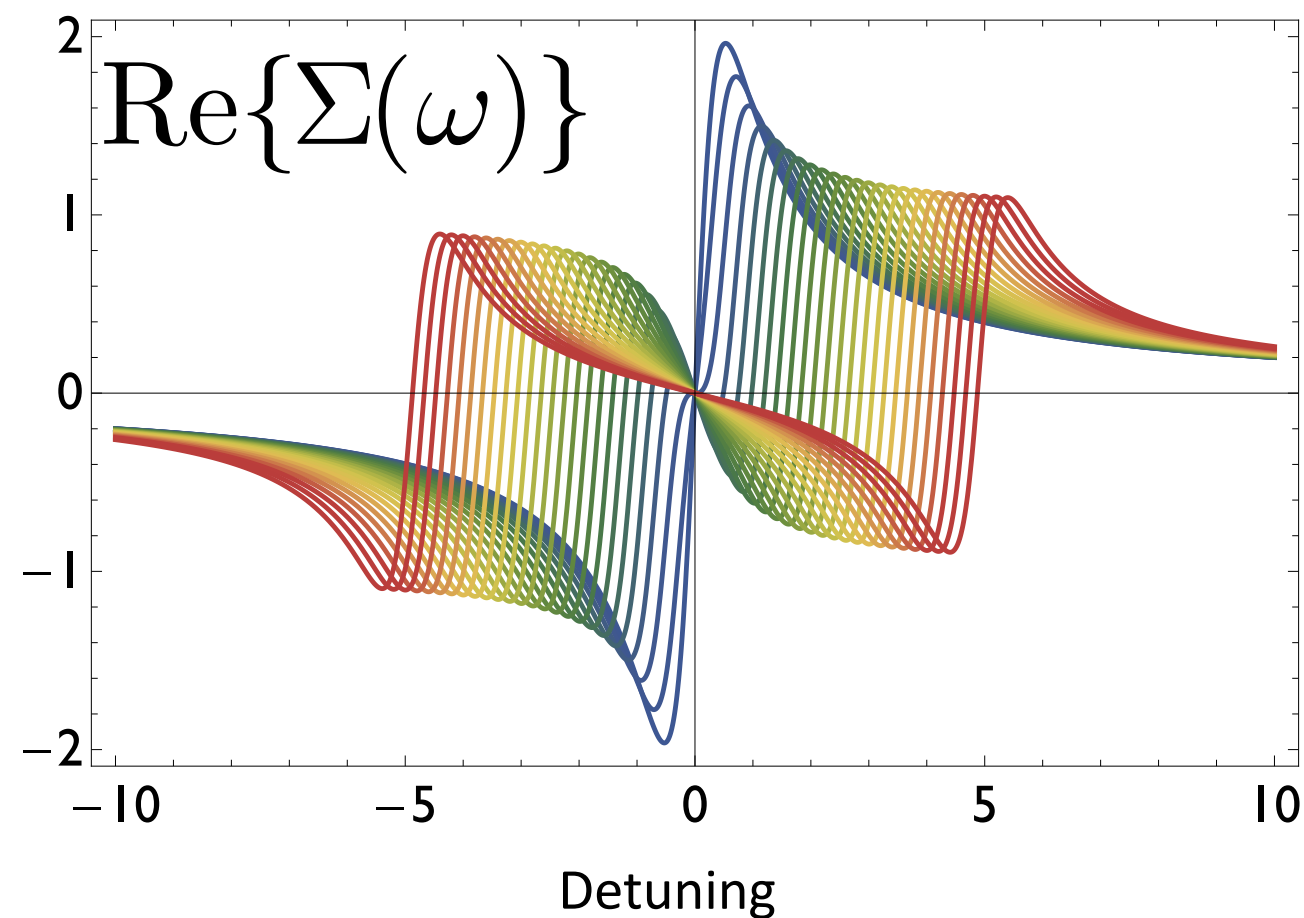
$$g_0 \hat{a}^\dagger \hat{a} (\hat{b} + \hat{b}^\dagger) \mapsto \underbrace{g_0 \alpha}_g (\delta \hat{a} + \delta \hat{a}^\dagger) (\hat{b} + \hat{b}^\dagger)$$

bilinear interaction
tunable coupling!



First order equation: dynamics

$$\delta \tilde{F}[\omega] = \delta \tilde{x} \frac{g_{\text{OM}}^2 U}{\omega_0} \left(\frac{1}{(\Delta' - \omega) - i\kappa / 2} + \frac{1}{(\Delta' + \omega) + i\kappa / 2} \right) \equiv \Sigma(\omega)$$





Dynamical Back-action

$$\ddot{x} + \overset{\text{Optical damping}}{(\Gamma_m + \boxed{\Gamma_{\text{opt}}})} \dot{x} + \overset{\text{Optical spring effect}}{(\Omega_m + \boxed{\delta\Omega_{\text{opt}}})^2} x = 0$$



Dynamical Back-action

$$\ddot{x} + (\Gamma_m + \Gamma_{\text{opt}})\dot{x} + (\Omega_m + \delta\Omega_{\text{opt}})^2 x = 0$$

Optical damping

$$\Gamma_{\text{opt}} = 4g^2 n_c \left(\frac{\kappa}{\kappa^2 + 4(\Delta + \Omega_m)^2} - \frac{\kappa}{\kappa^2 + 4(\Delta - \Omega_m)^2} \right)$$

Optical spring effect

$$\delta\Omega_{\text{opt}} = 4g^2 n_c \left(\frac{\Delta - \Omega_m}{\kappa^2 + 4(\Delta + \Omega_m)^2} - \frac{\Delta + \Omega_m}{\kappa^2 + 4(\Delta - \Omega_m)^2} \right)$$

$g = g_{\text{om}} x_{\text{zpf}}$
Optomechanical
coupling rate

x_{zpf} Zero-point fluctuation
 n_c Intracavity photon number



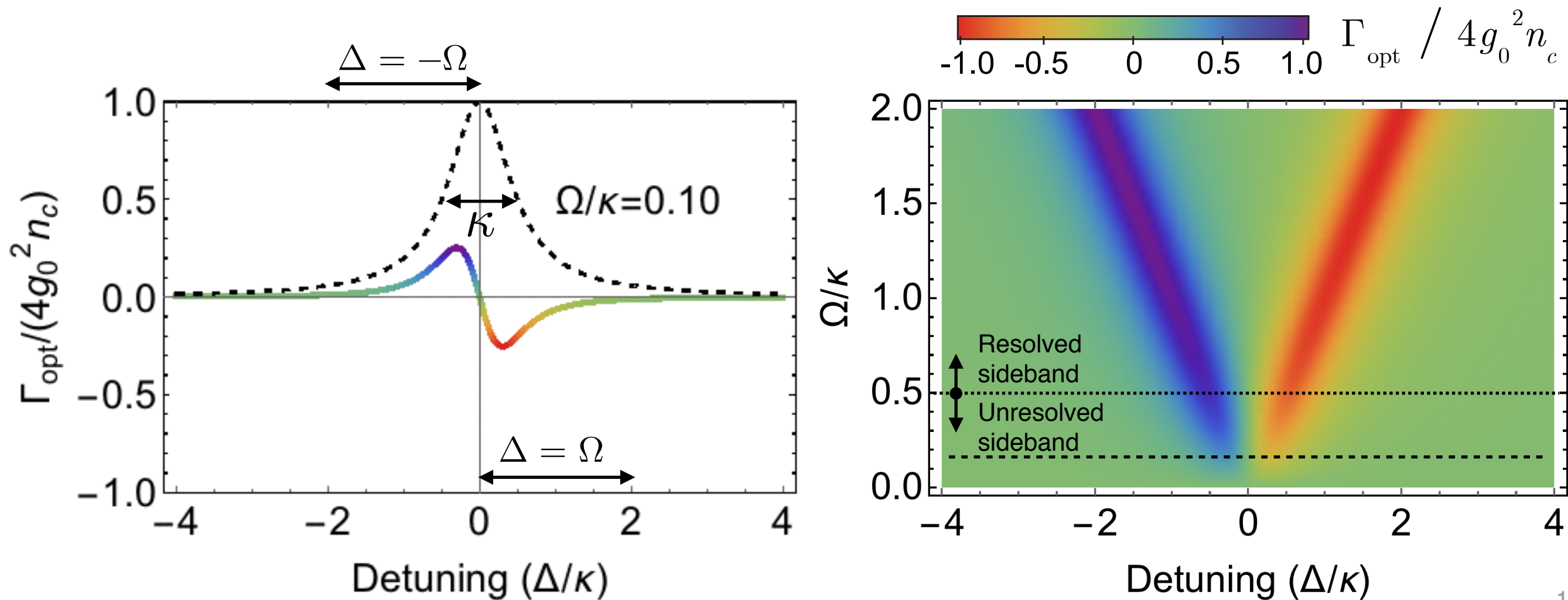
Dynamical Back-action

Effective damping

$$\Gamma_{\text{eff}} = (\Gamma_m + \Gamma_{\text{opt}})$$

Optical damping

$$\Gamma_{\text{opt}} = 4g_0^2 n_c \left(\frac{\kappa}{\kappa^2 + 4(\Delta + \Omega_m)^2} - \frac{\kappa}{\kappa^2 + 4(\Delta - \Omega_m)^2} \right)$$





Cooperativity

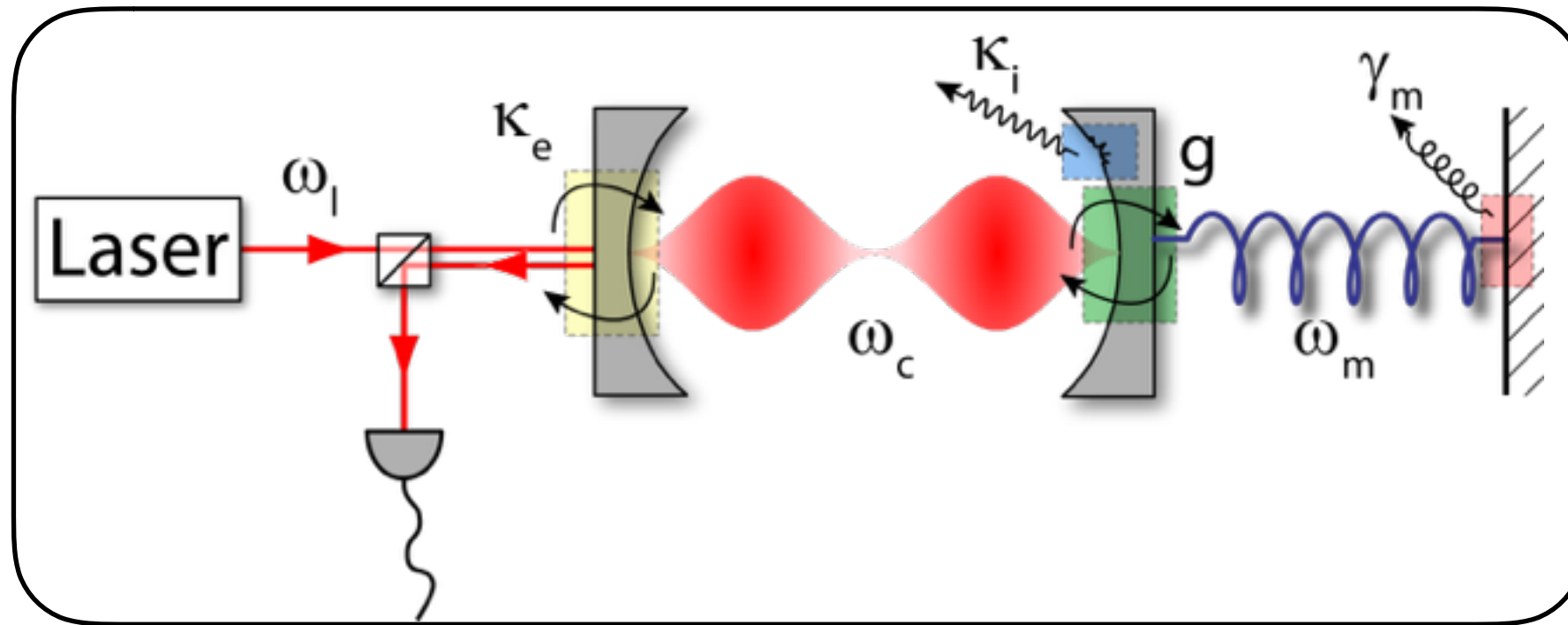


Figure of merit: cooperativity ($\Gamma_{\text{opt}} = \pm \Gamma_m$)

$$C = \frac{4g^2 n_c}{\kappa \Gamma_m}$$

n_c Intracavity photon number

κ Optical mode linewidth

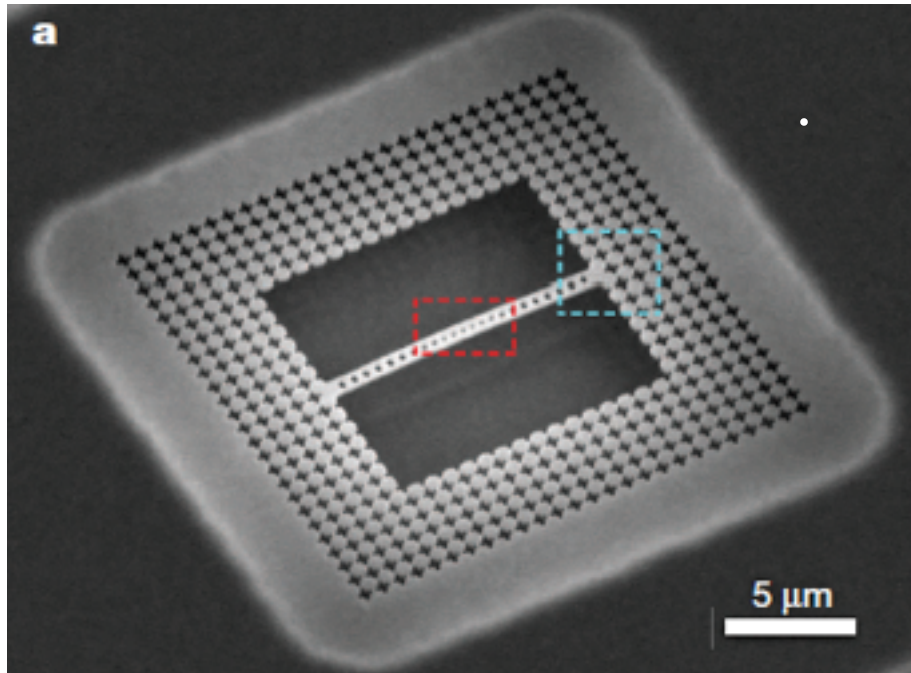
Γ_m Mechanical linewidth

$$g = g_{\text{om}} x_{\text{zpf}}$$

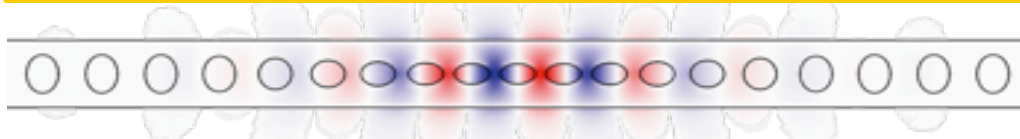
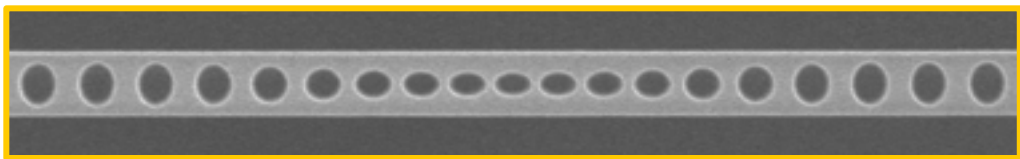


Real world devices

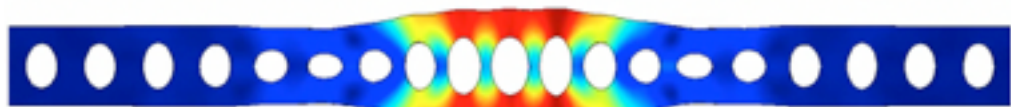
Non-linear optics and mechanical Optical Memories and Switches



Confined photonic crystal optical Mode



Mechanical breathing mode

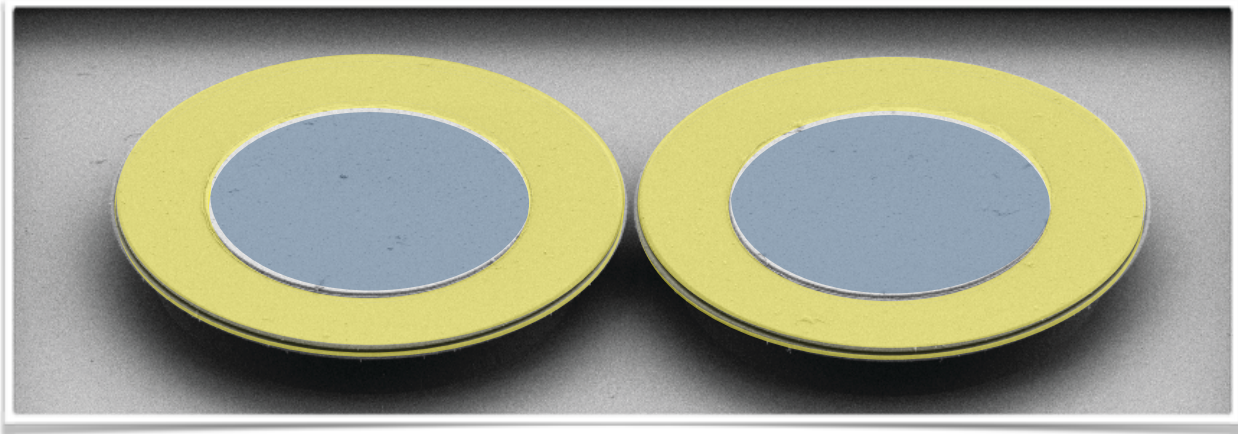


4 GHz

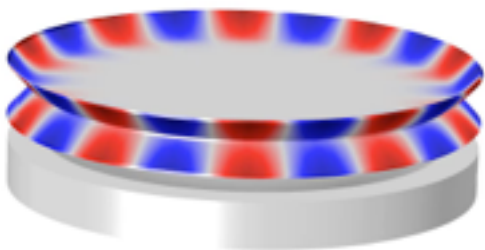
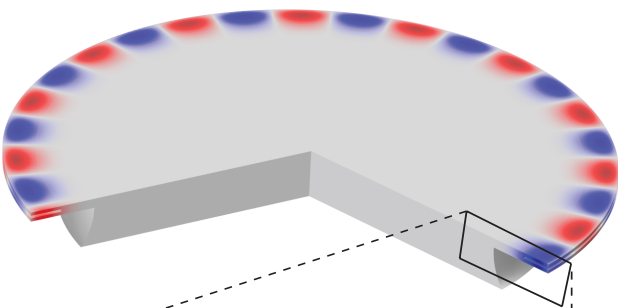
Alegre et. al, Nature **472**, 69 (2011)

J. Chan et. al, Nature **478**, 89 (2011)

Coupled Optomechanical Oscillators



Whispering Gallery
Optical mode



Anti-symmetric
mechanical mode

50 MHz

M. Zhang, G. S. Wiederhecker et al, PRL **109**, 233906 (2012)

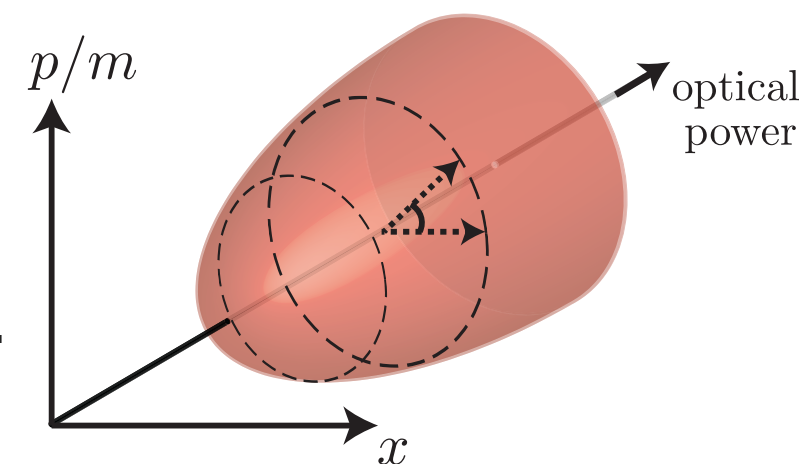
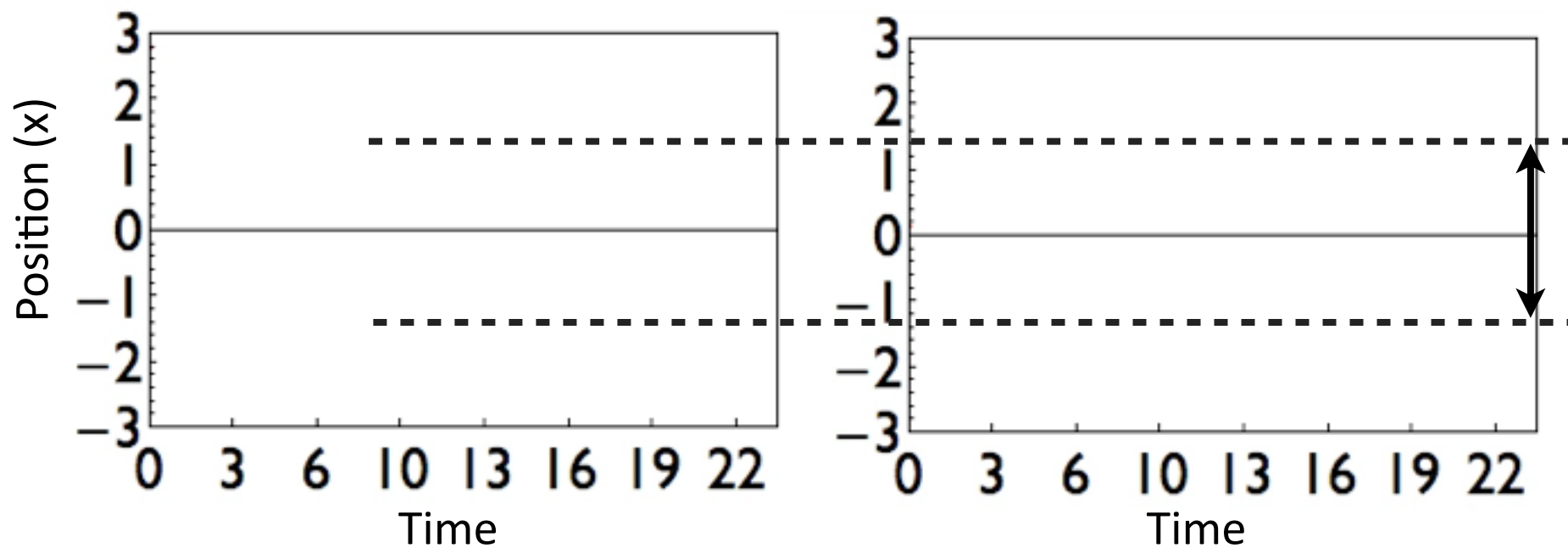
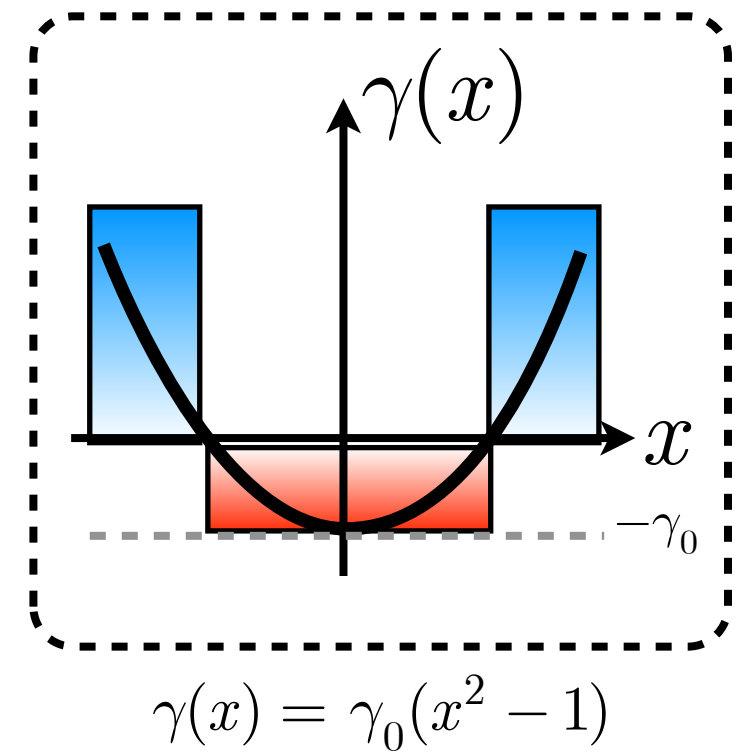
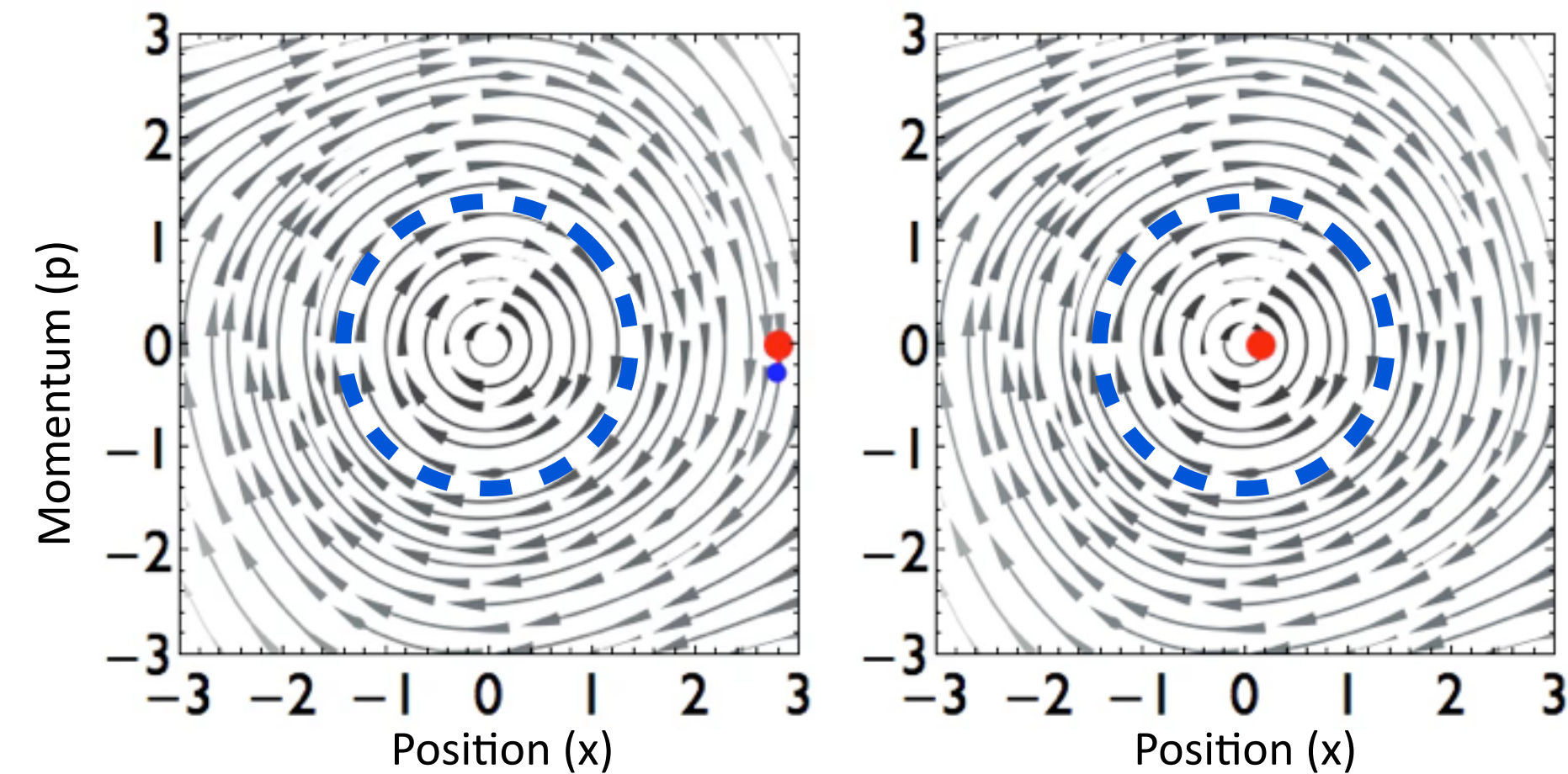


Outline

- ★ **Optical and acoustic mode interaction**
- ★ Optical force actuation
- ★ Dynamical back-action
- ★ **Optomechanical clocks**
- ★ Bullseye - a case study
- ★ Outlook

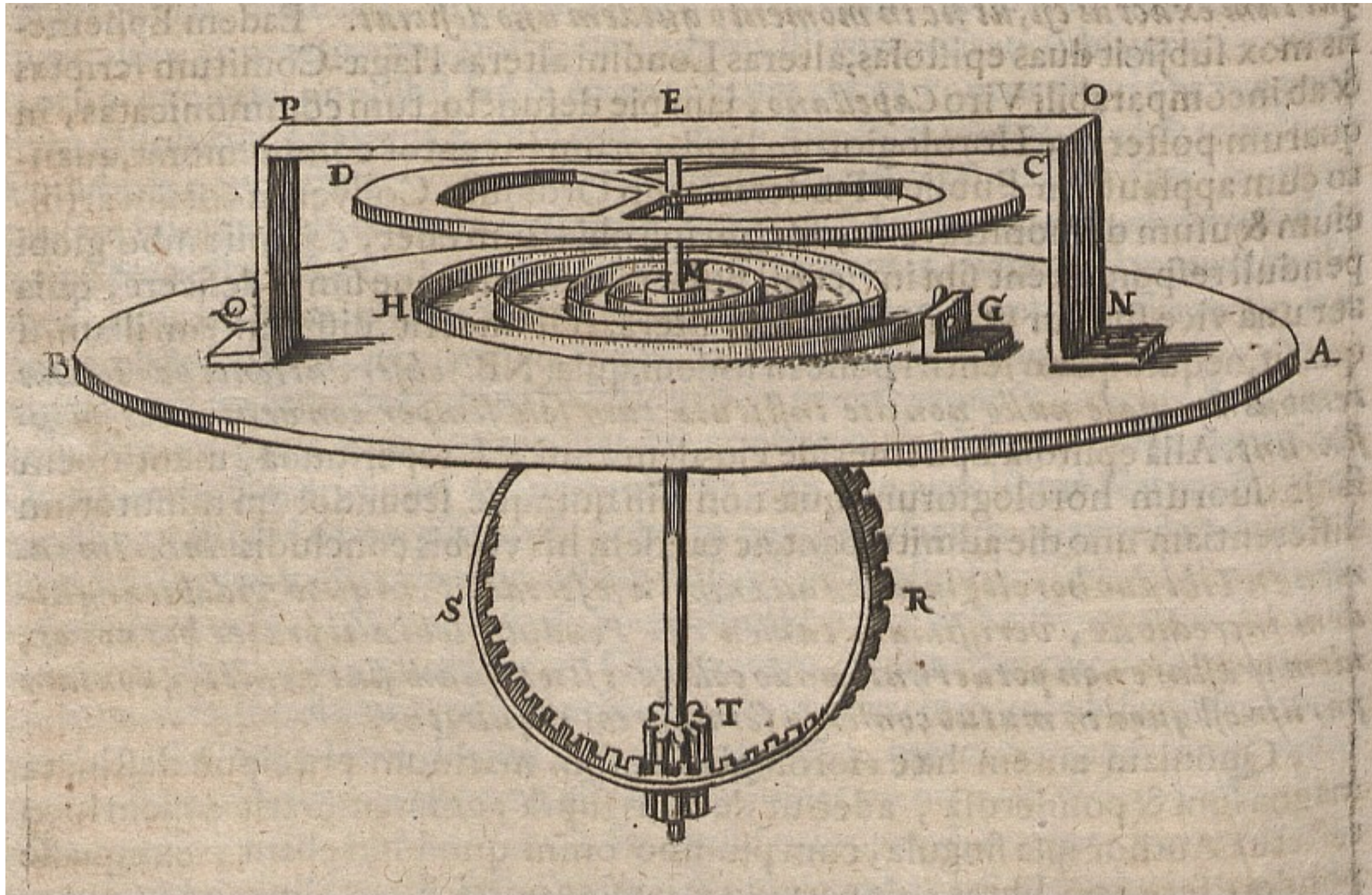


Self-sustaining oscillators





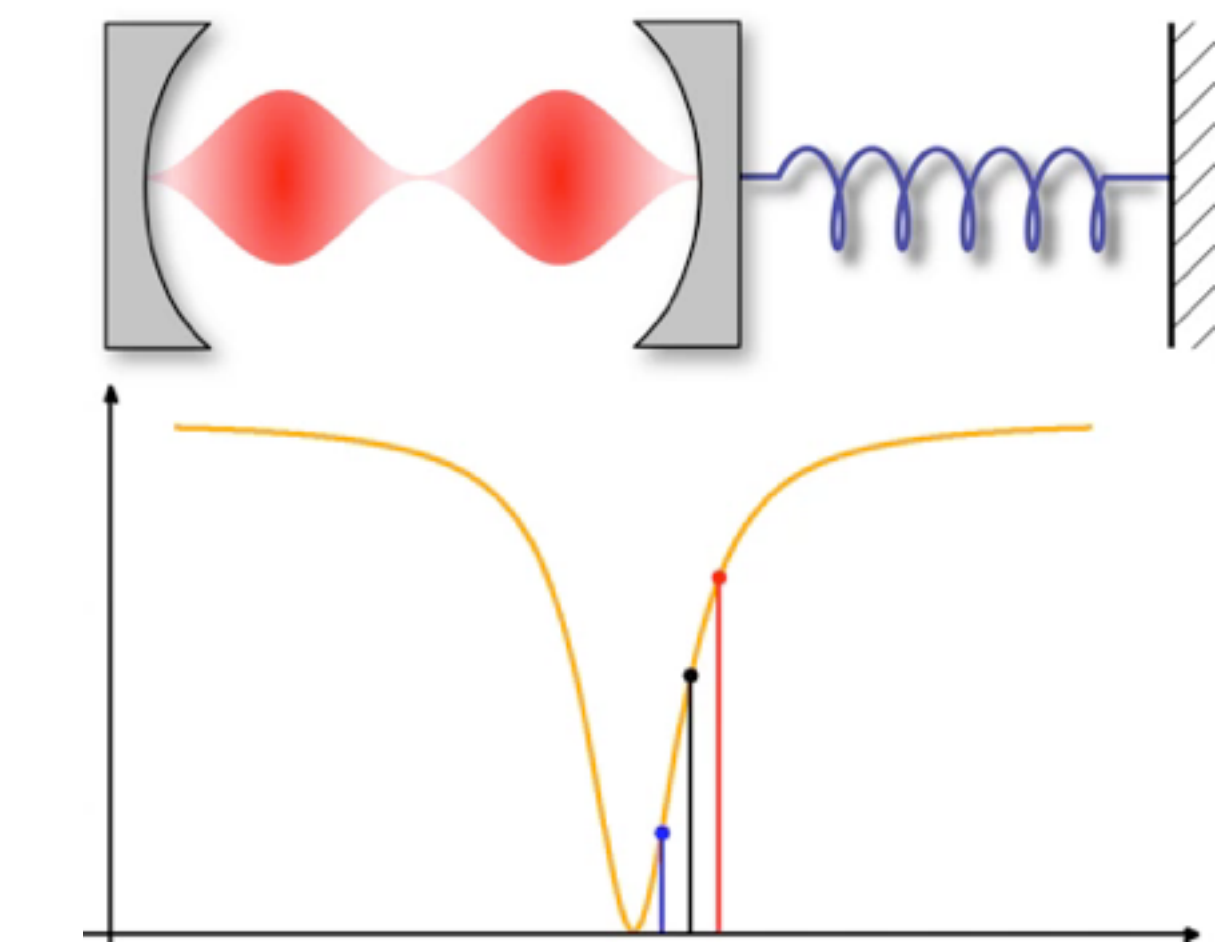
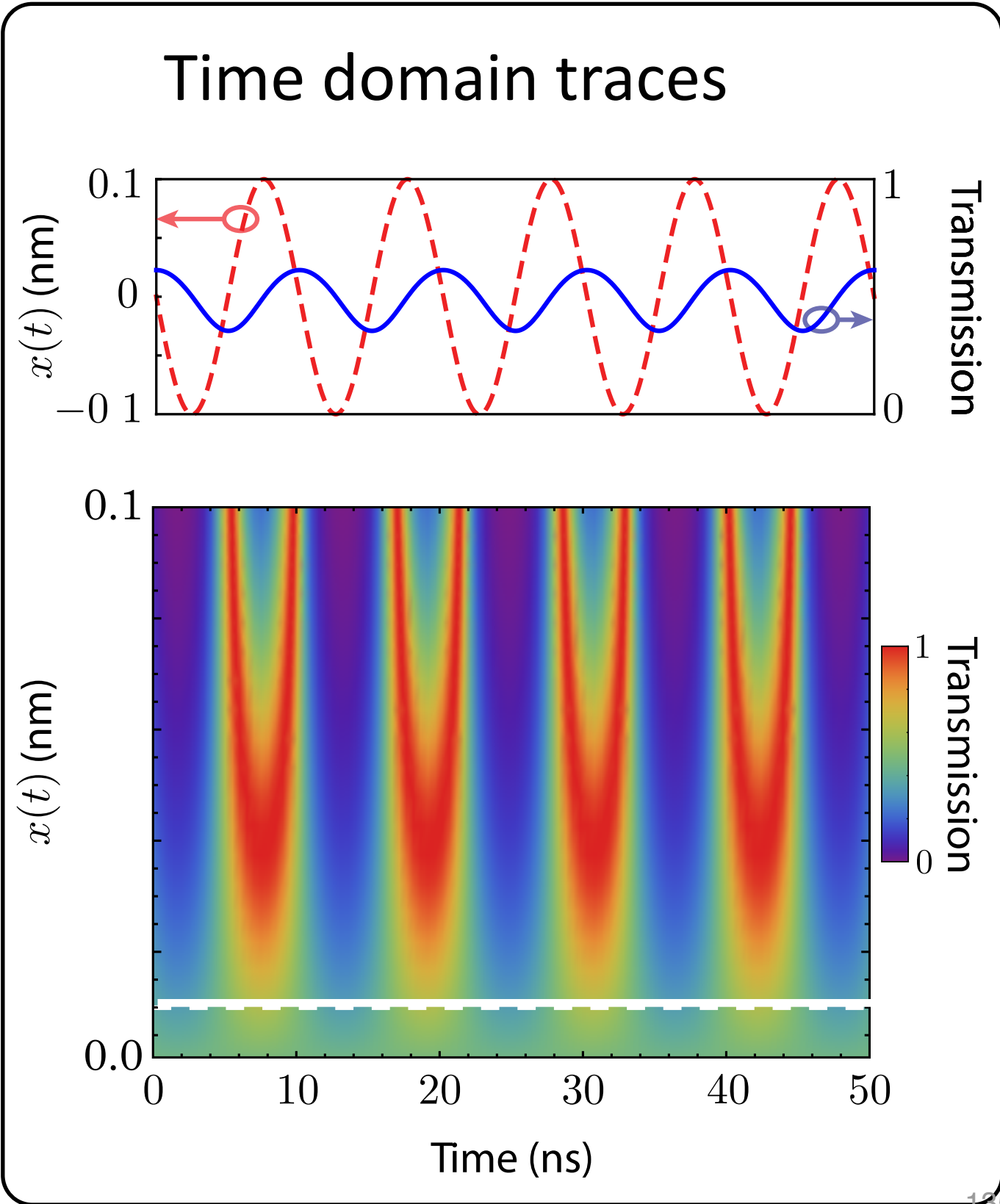
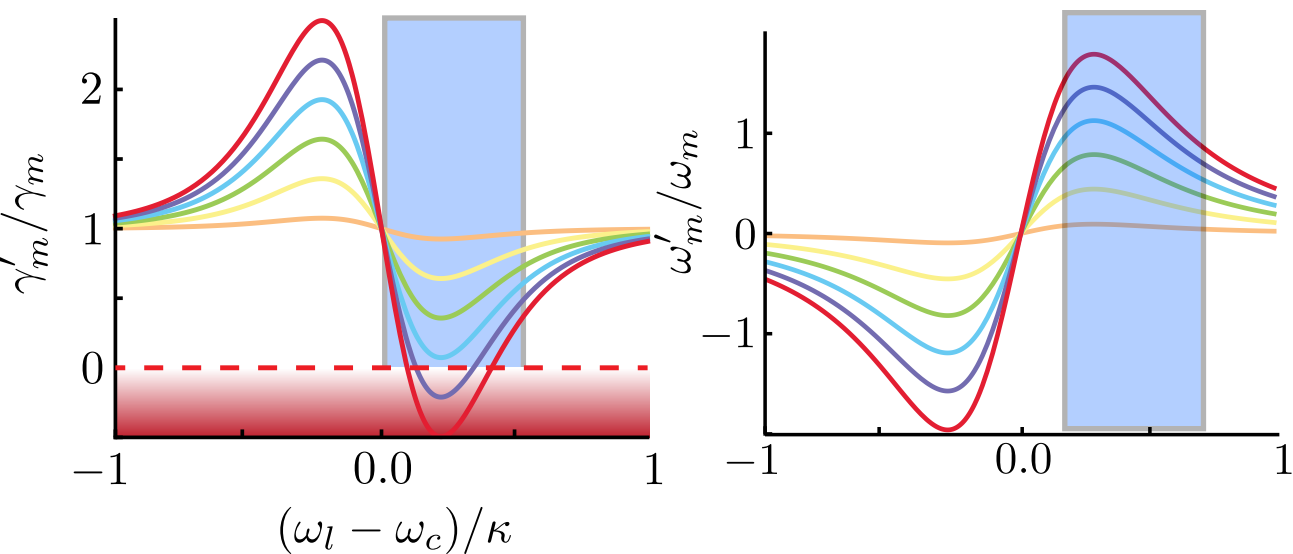
Building clocks



Quelle: Deutsche Fotothek



Optical Spring Effect, Cooling & Heating

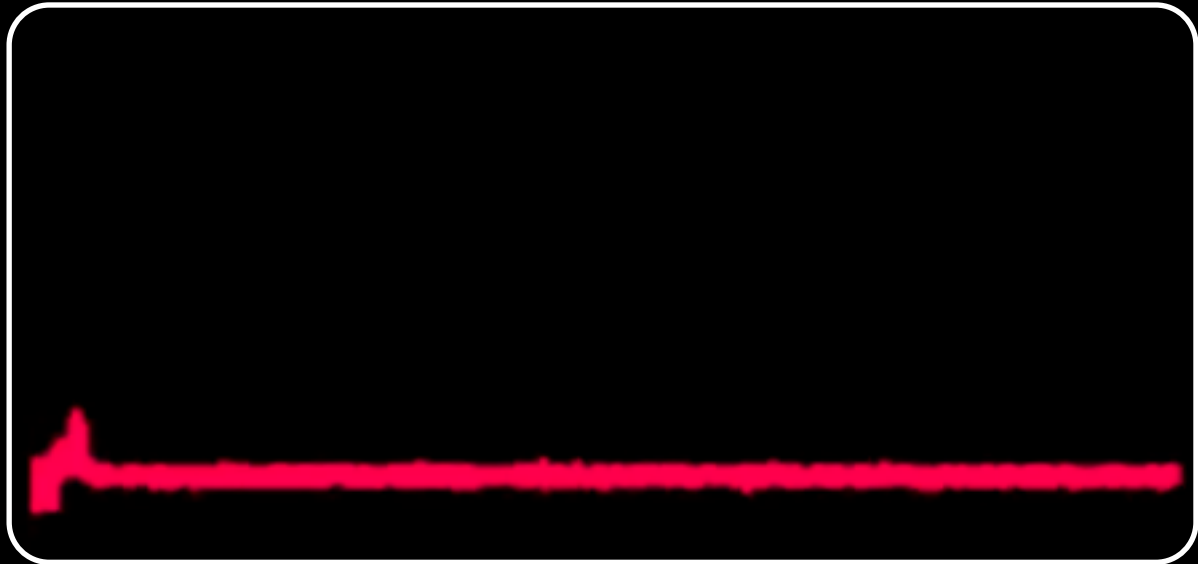




Limit cycle oscillations at the Lab.



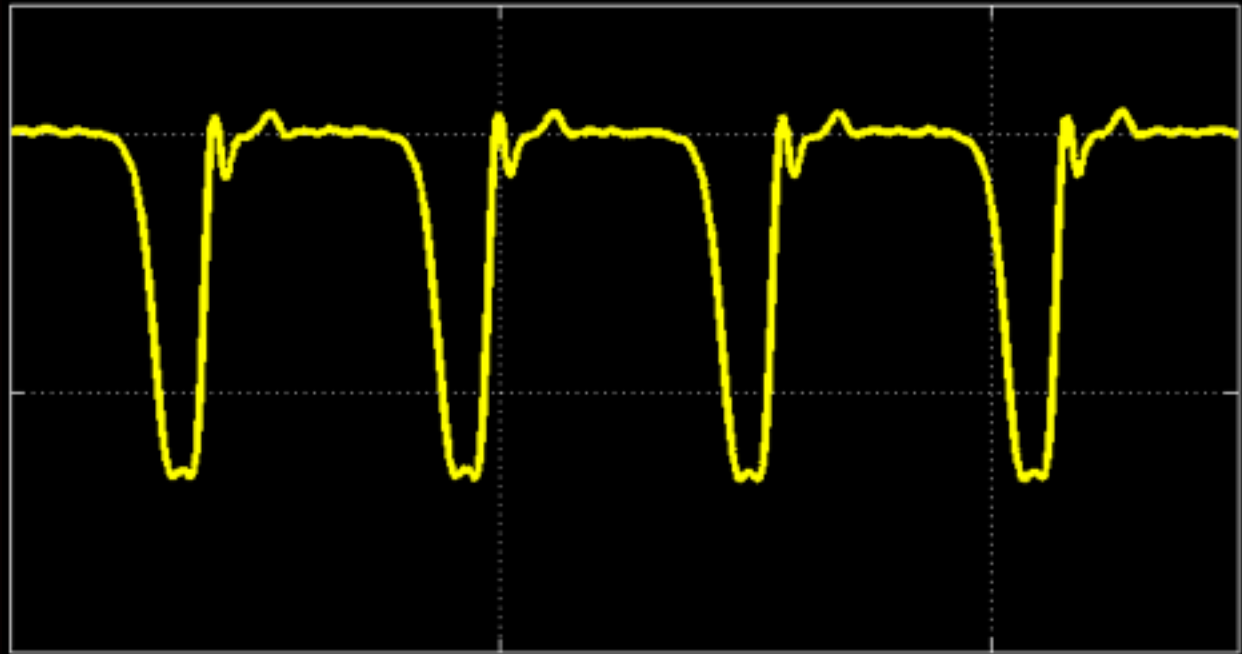
Relative RF Power (dB)



0 0.5 1.0 1.5 2.0

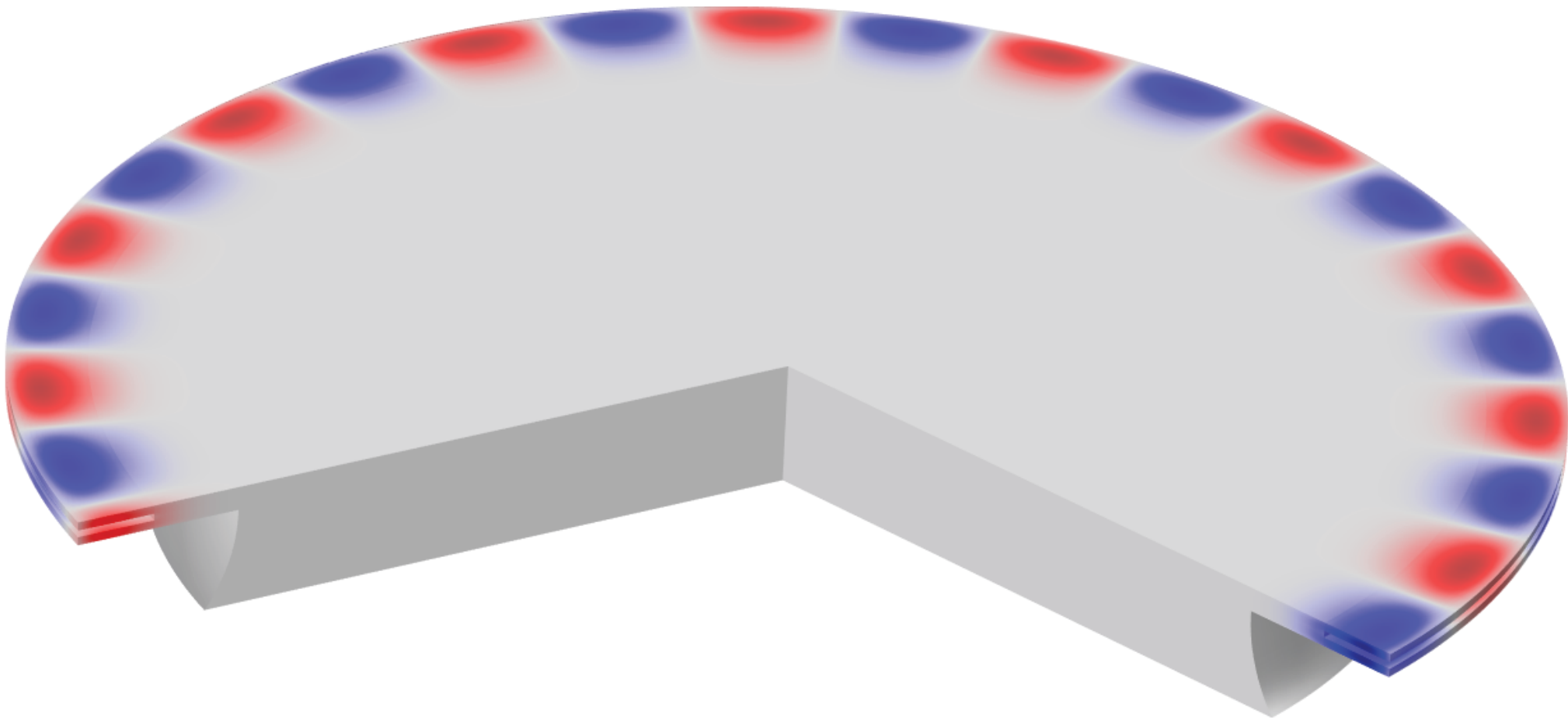
Frequency (GHz)

Amplitude (arb .units)



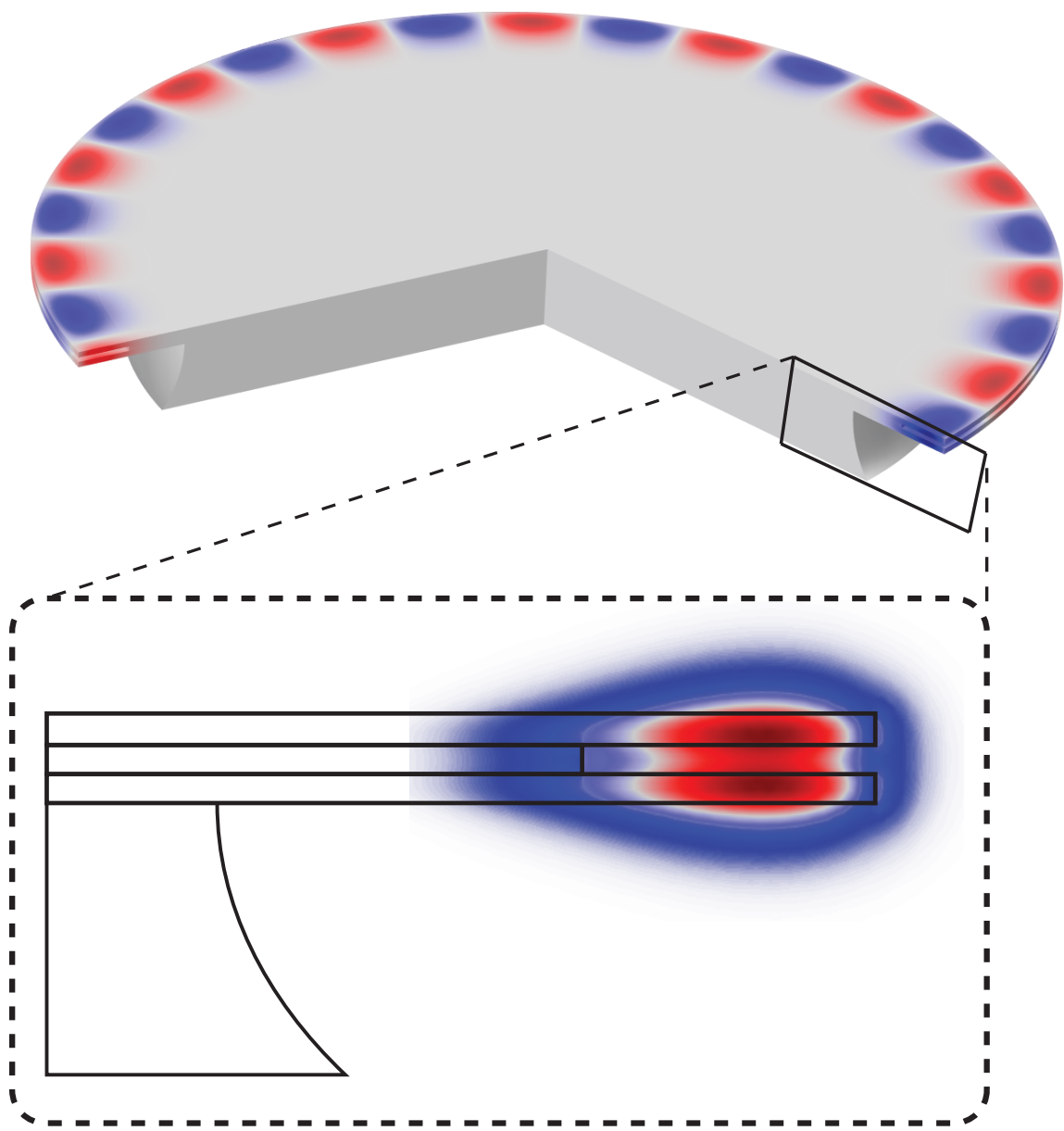


Double-disk Optomechanical Cavity

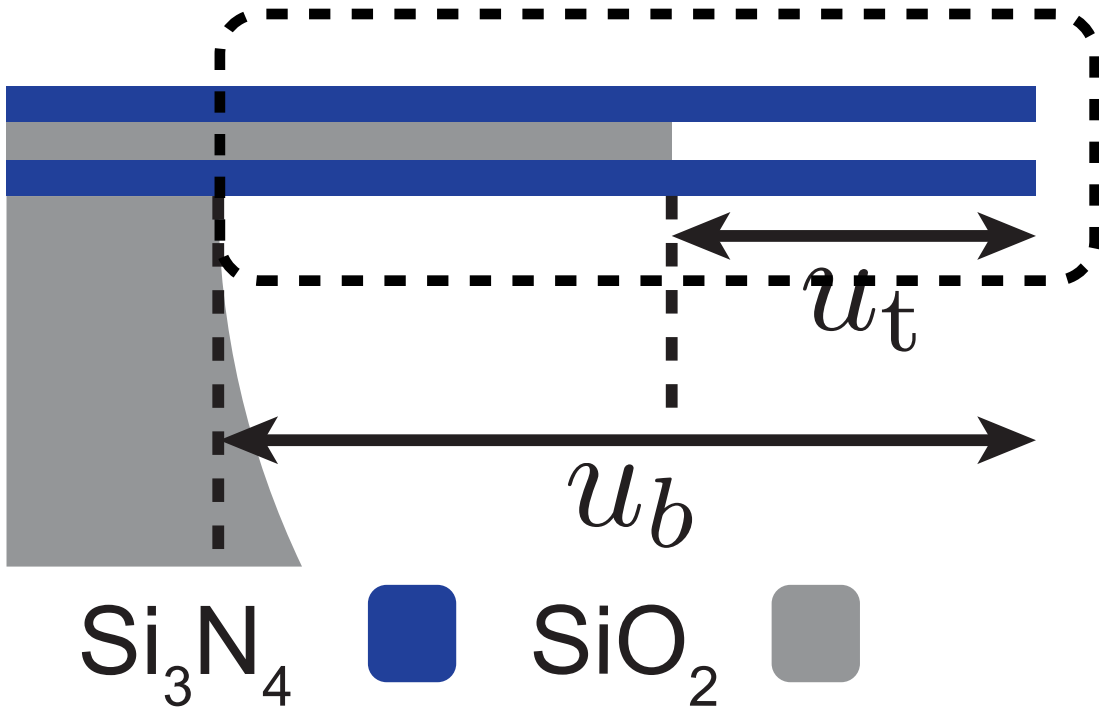




Single Optomechanical Cavity



Whispering Gallery Optical mode



Mechanical Modes



200 MHz



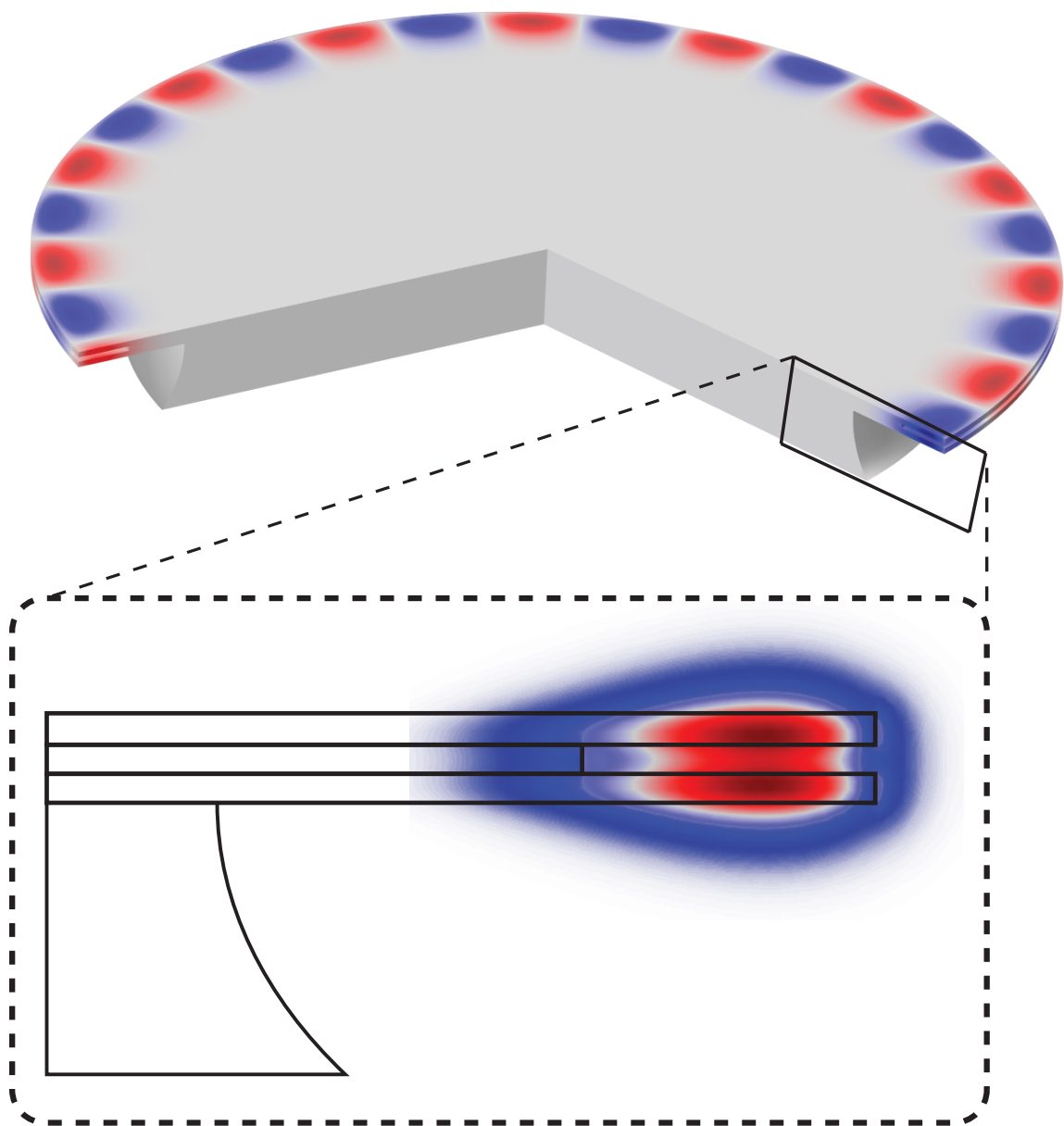
50 MHz



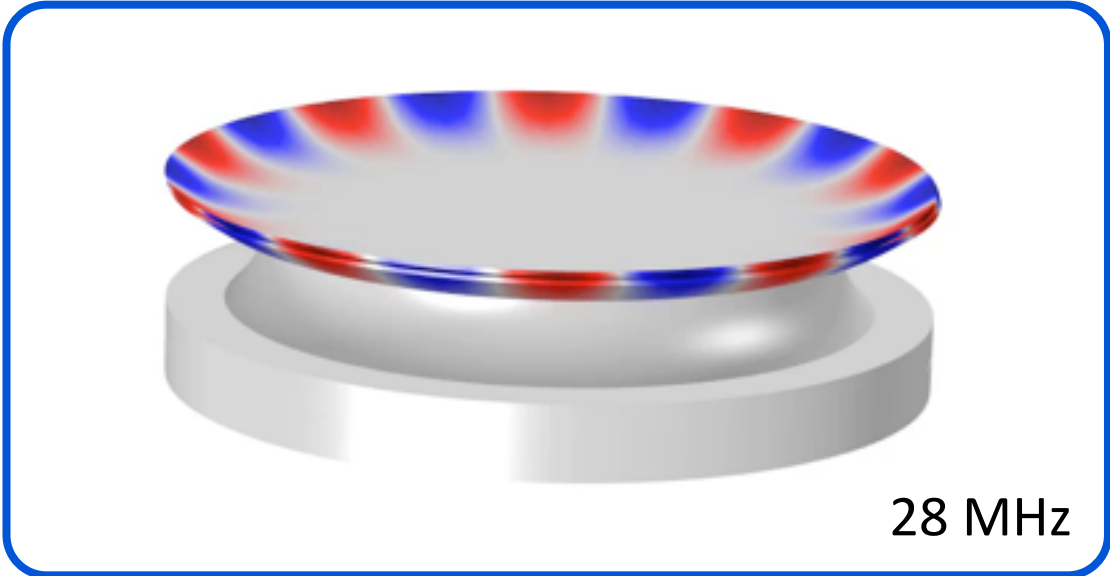
28 MHz



Double-disk Optomechanical Cavity

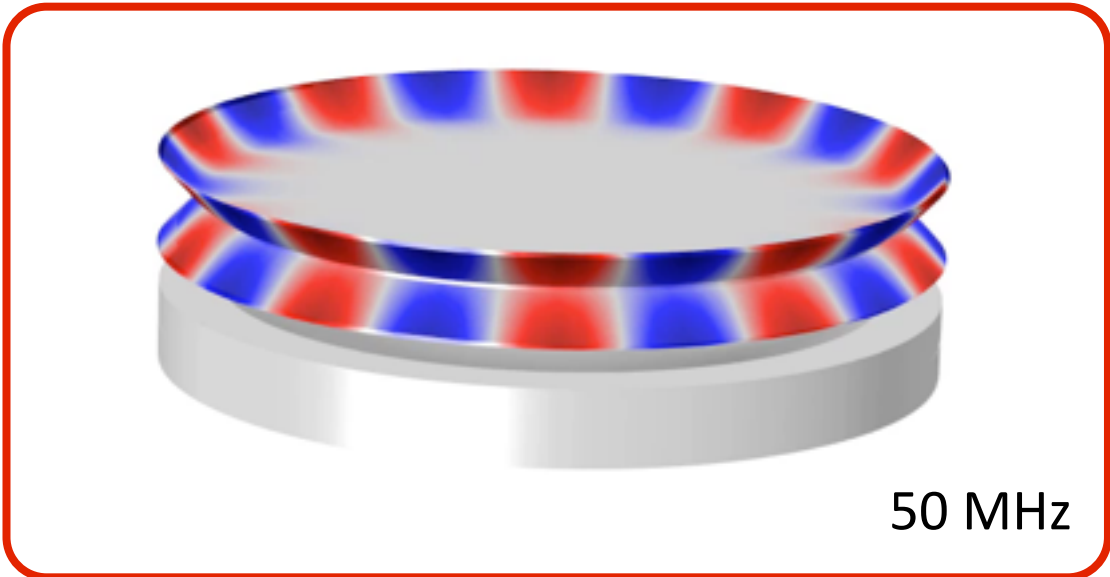


Whispering Gallery Optical mode



28 MHz

Symmetric mechanical mode

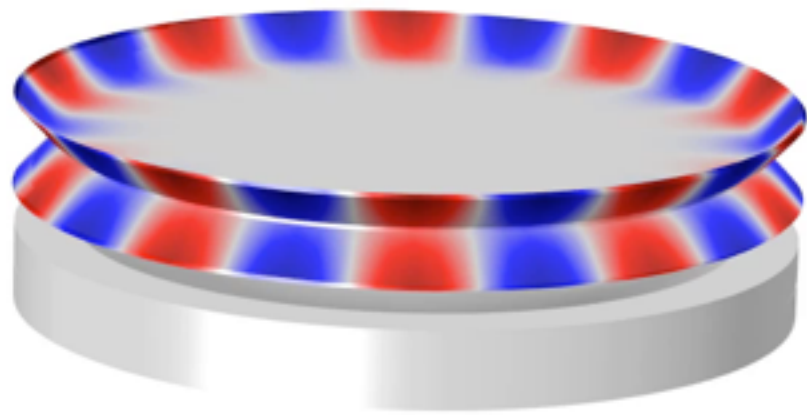
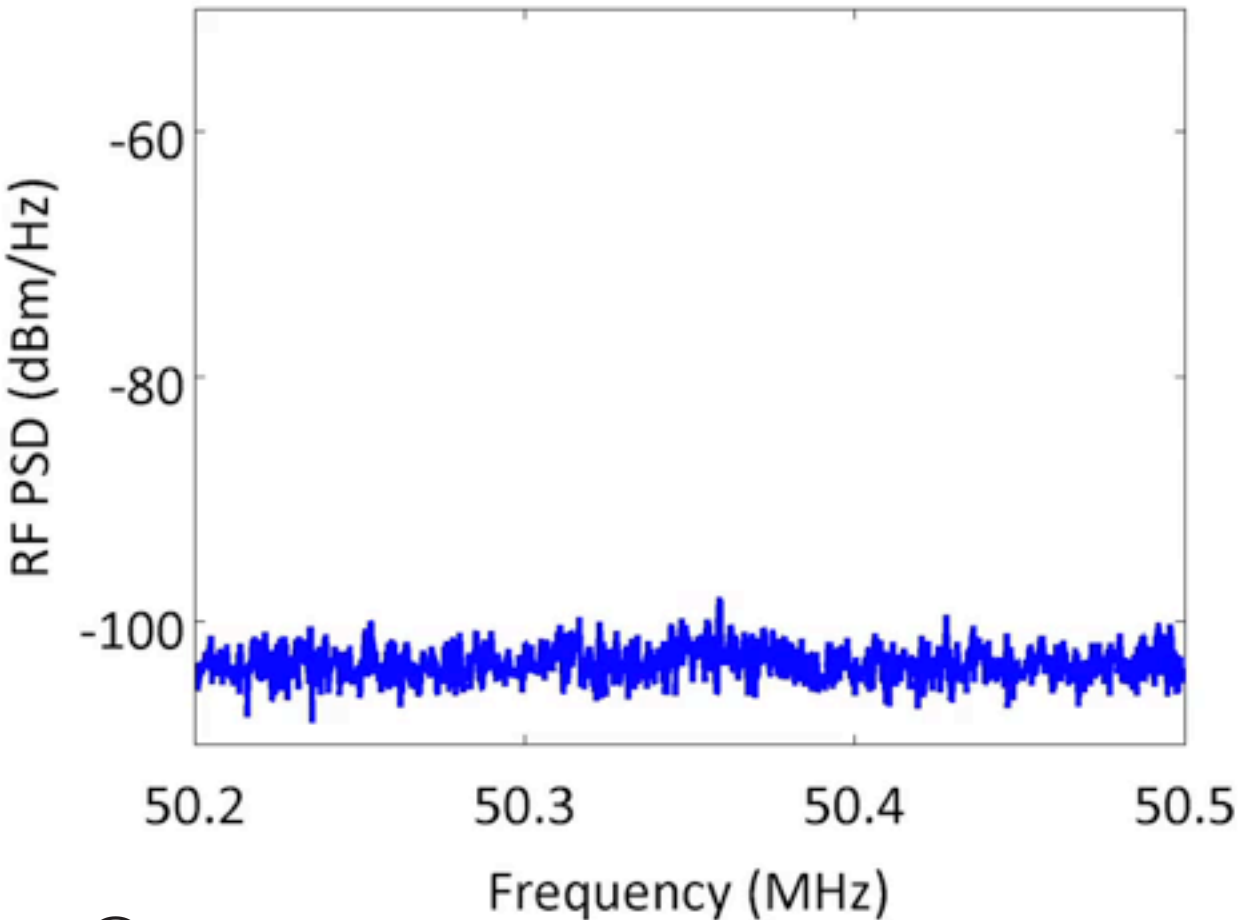
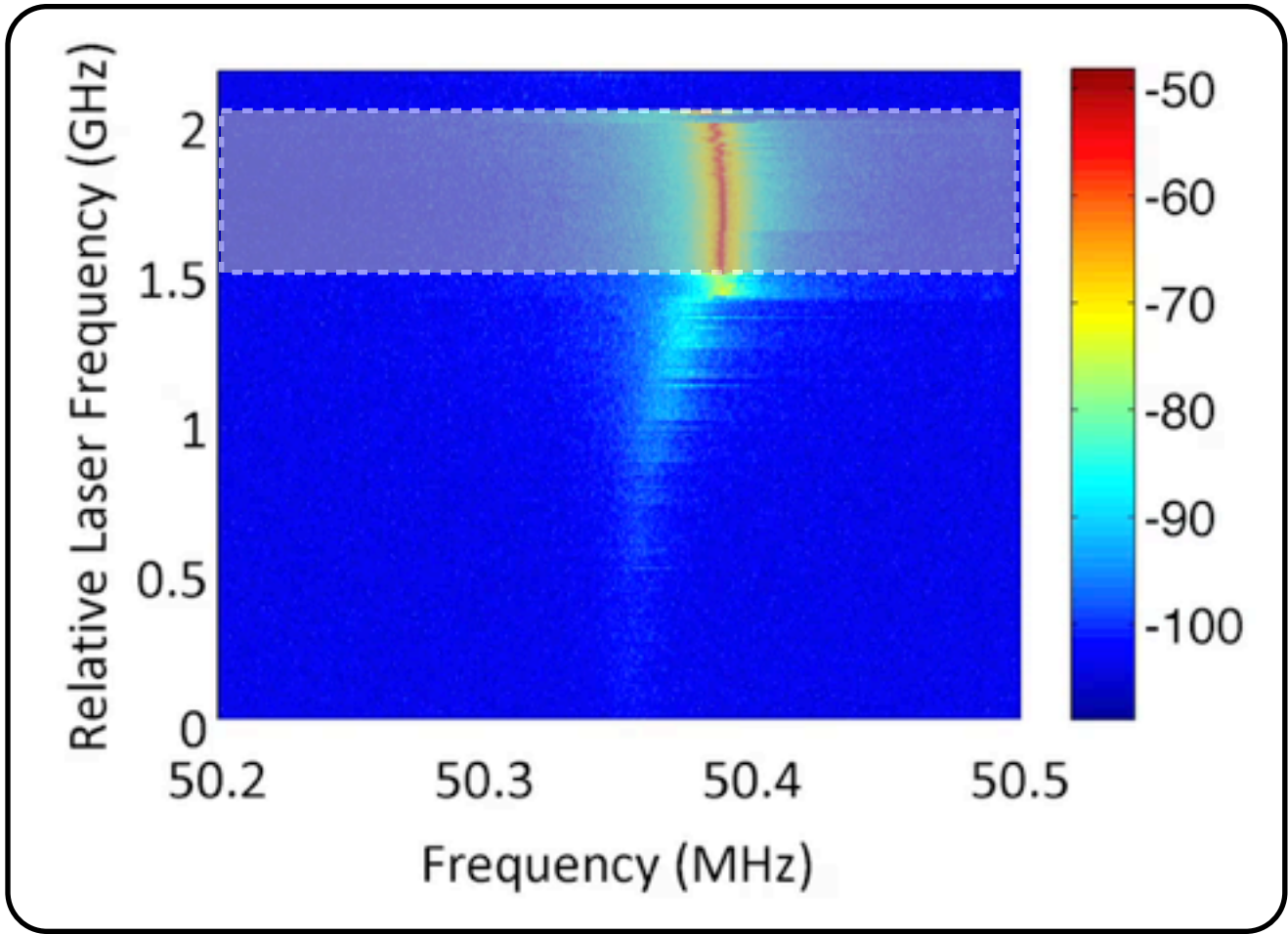


50 MHz

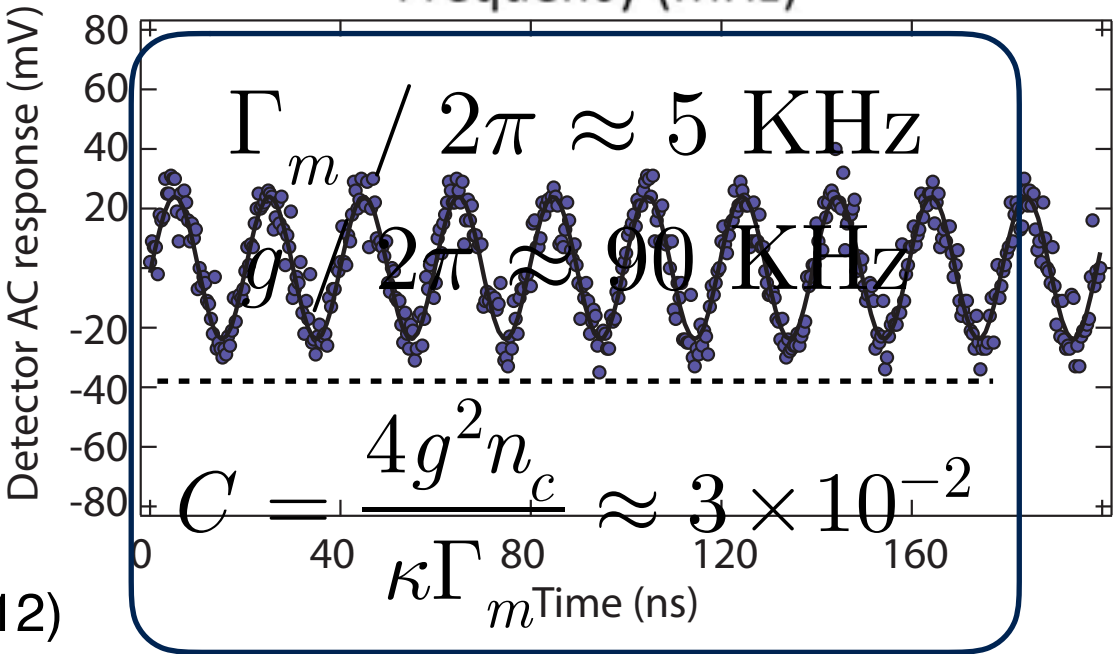
Anti-symmetric mechanical mode



Double-disk Limit cycle

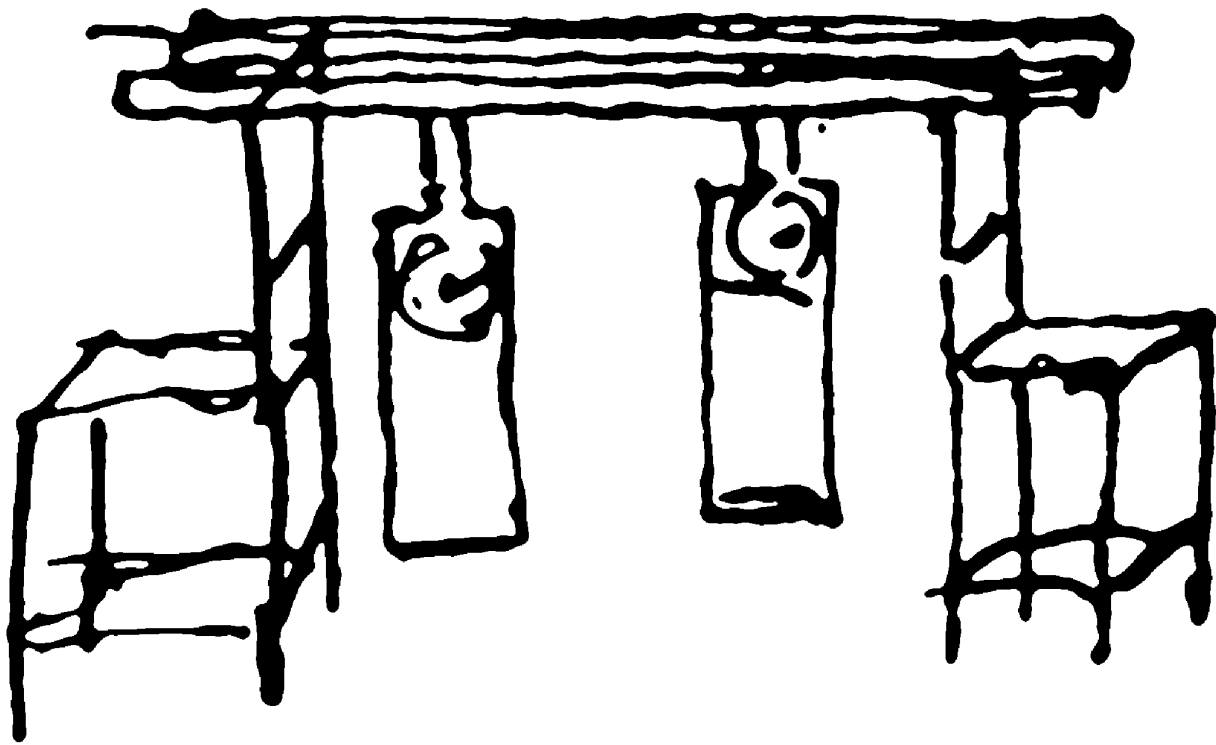


$$f = 50.35 \text{ MHz}$$





Synchronization of Oscillators



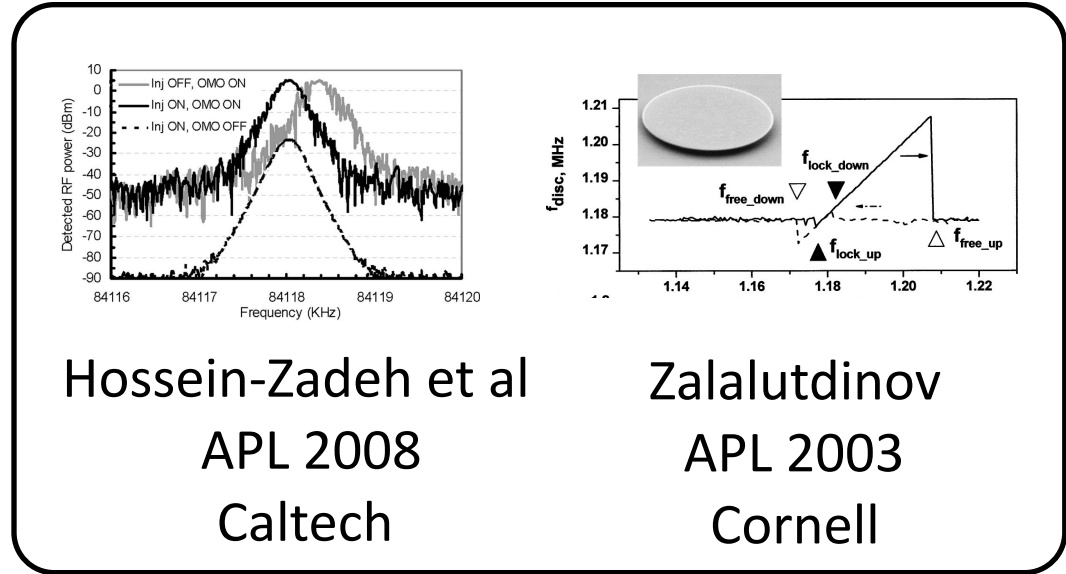
“It is quite worth noting that when we suspended two clocks so constructed from two hooks imbedded in the same wooden beam, the motions of each pendulum in opposite swings were so much in agreement that they never receded the least bit from each other and the sound of each was always heard simultaneously...”

“Further, if this agreement was disturbed by some interference, it reestablished itself in a short time.”

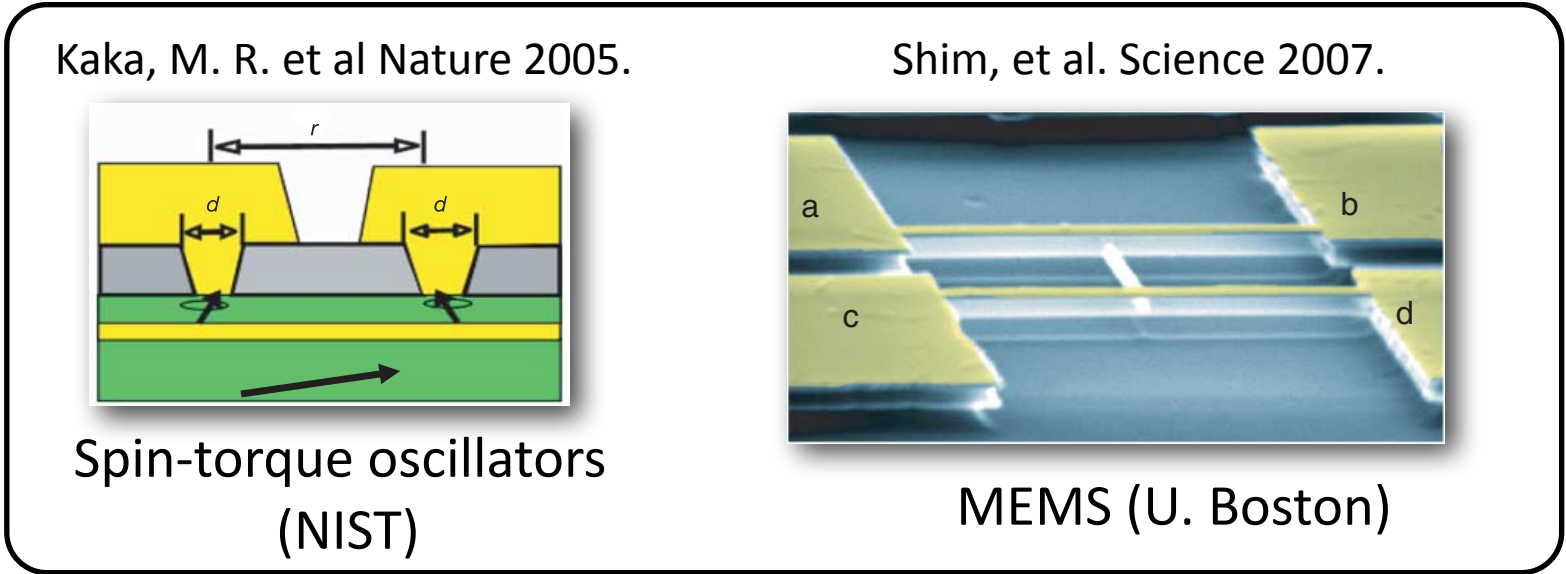
Christiaan Huygens, 26 February 1665.



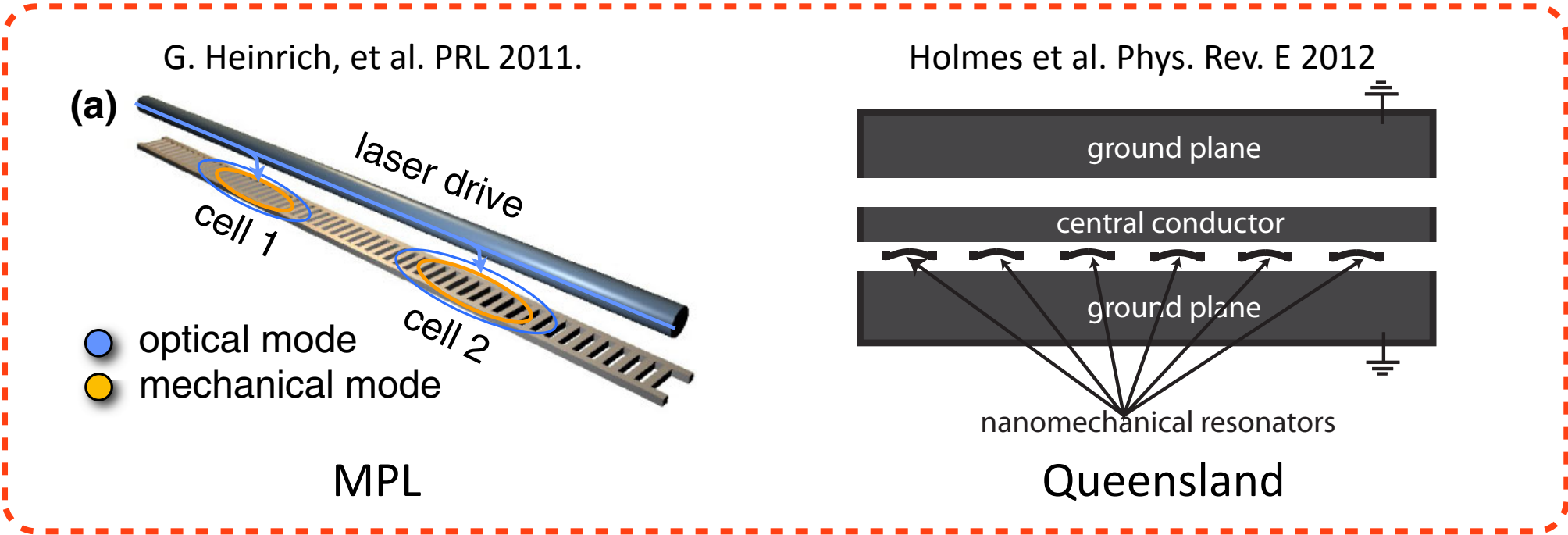
Synchronization at the nanoscale



Entrainment by external force



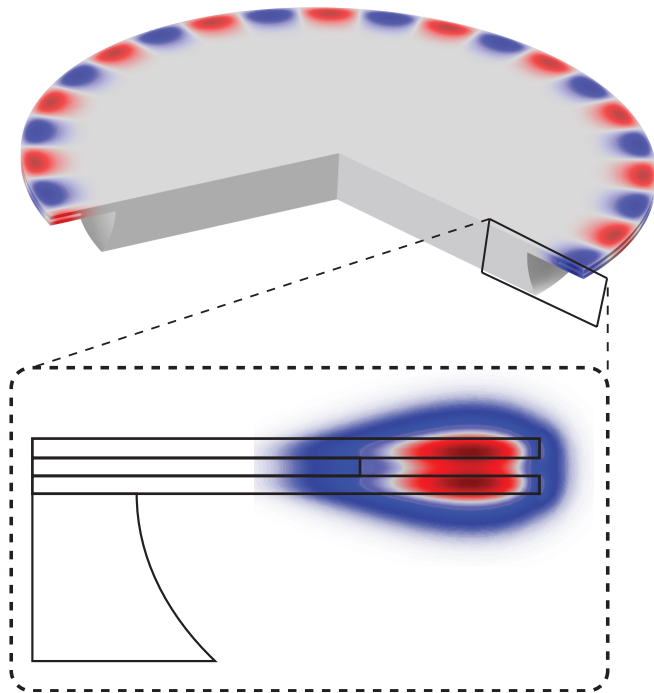
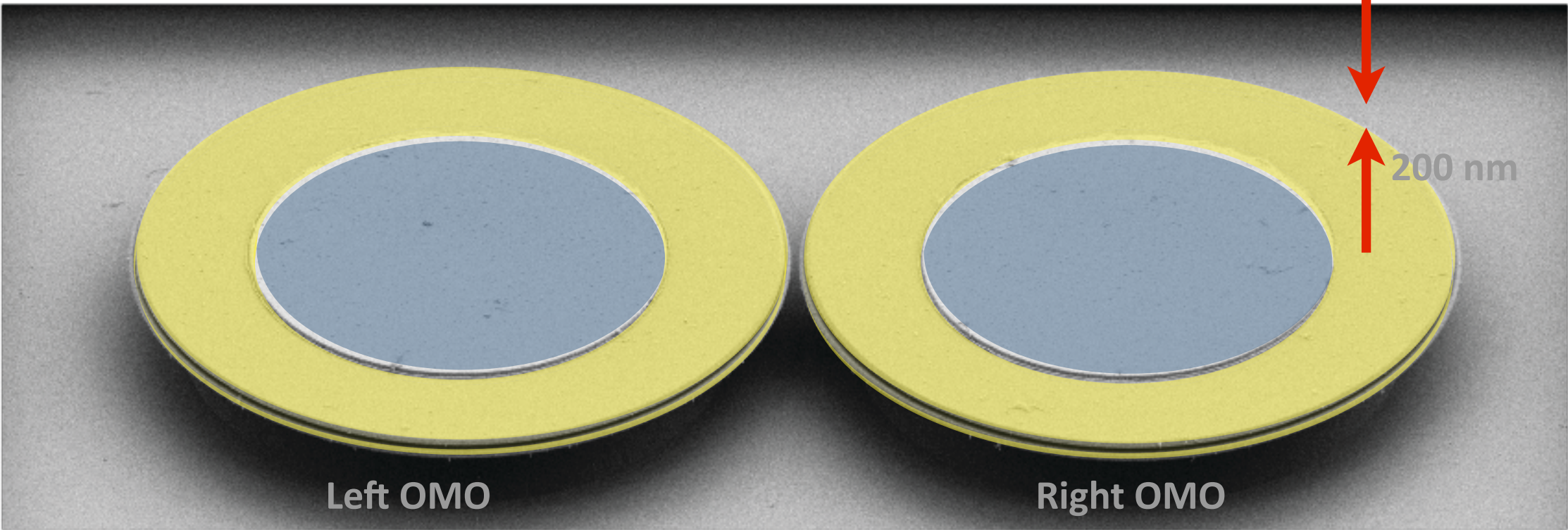
Mutual synchronization



Theory: Mutual synchronization with Nanomechanical oscillators



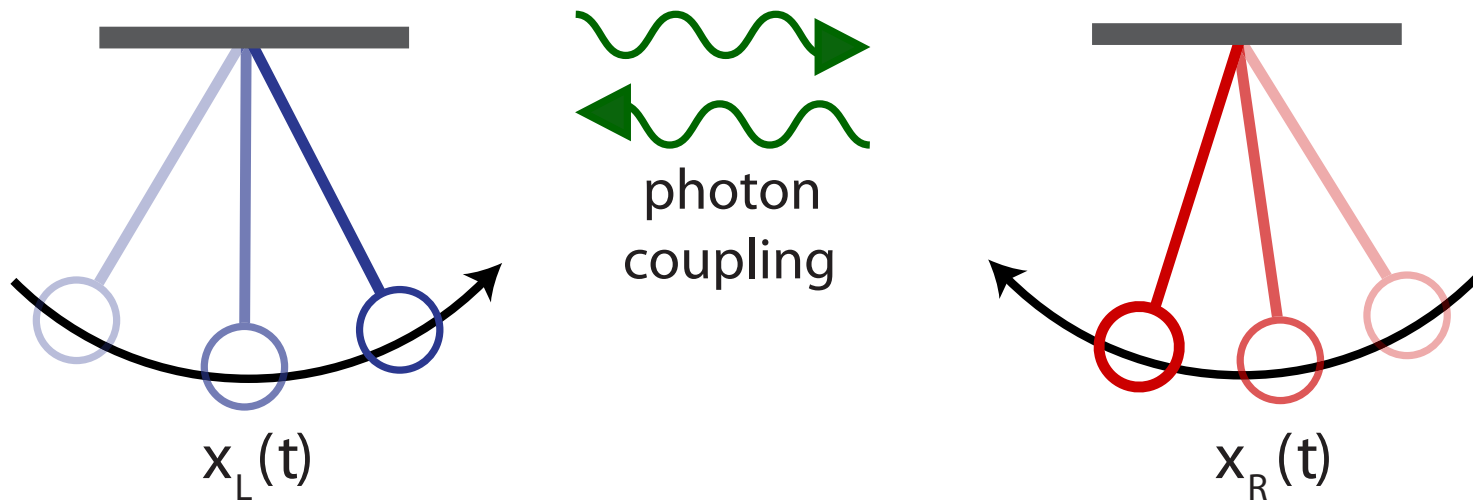
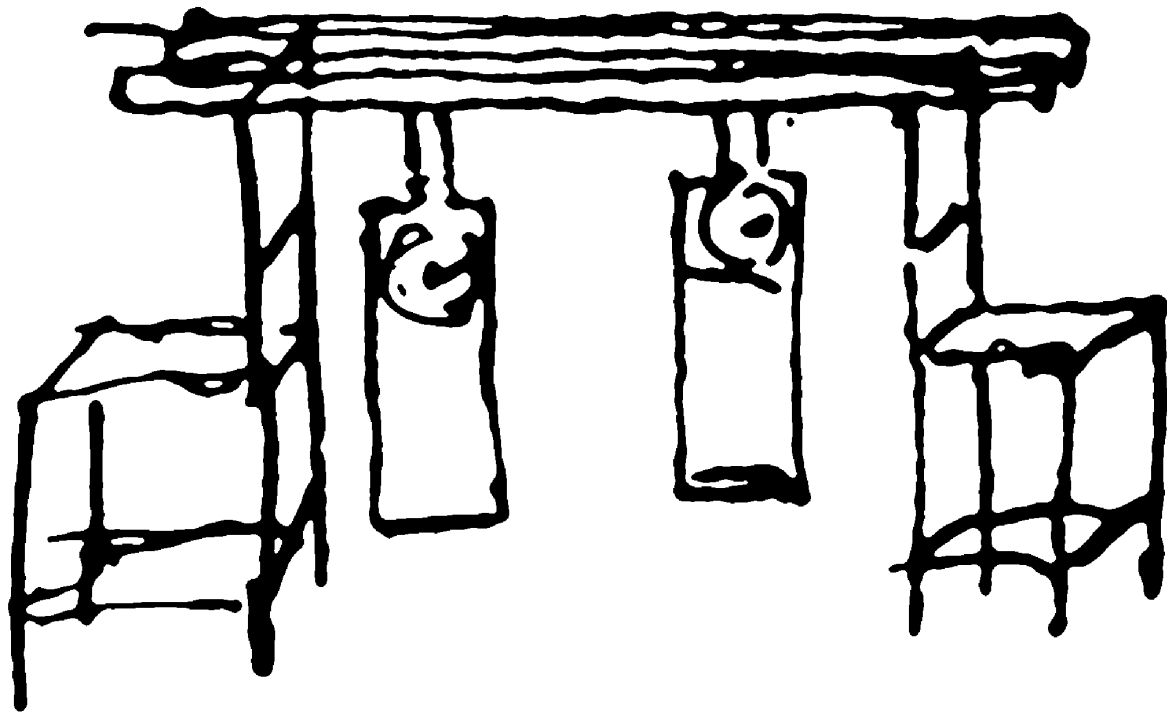
Coupled Optomechanical Oscillators





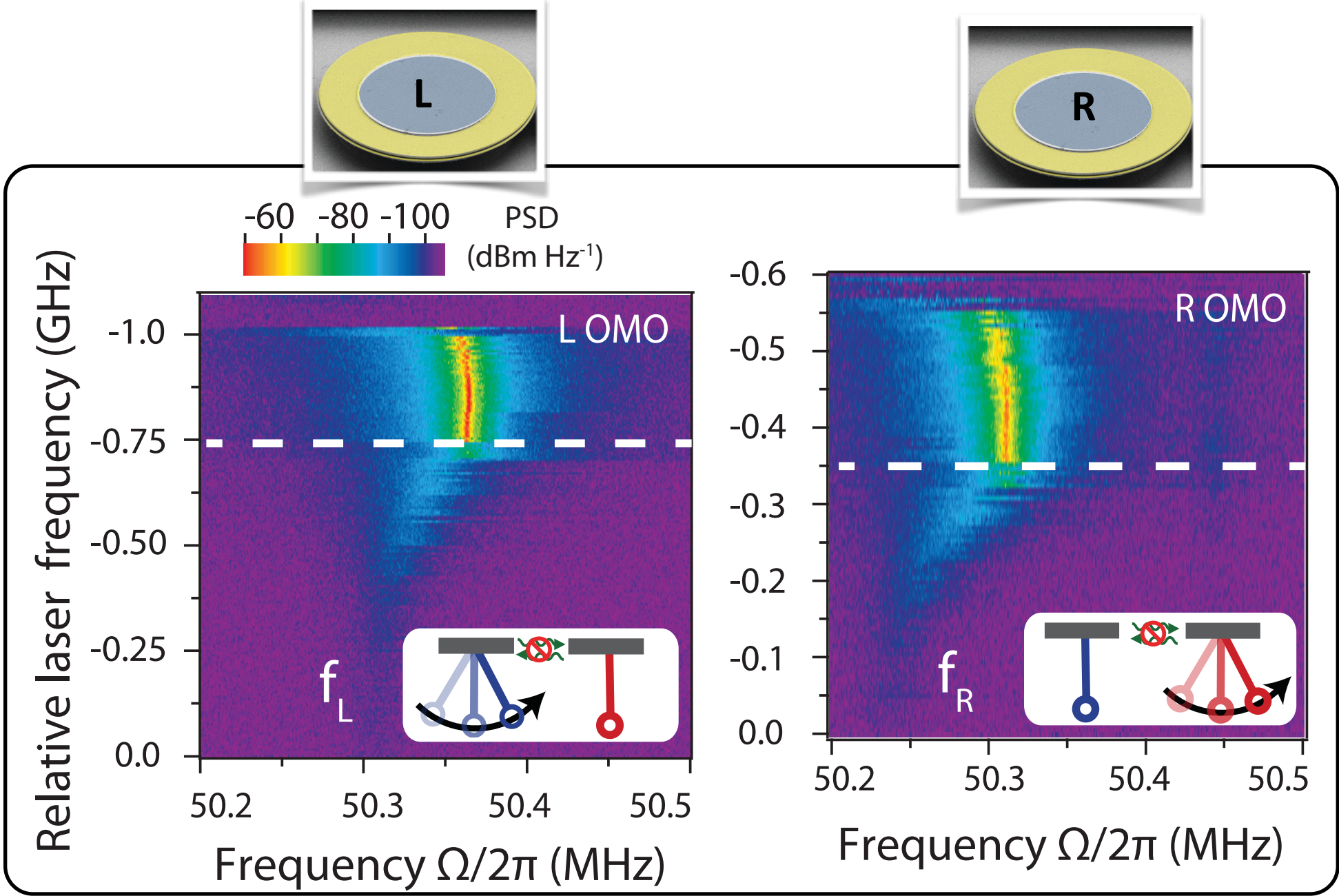
Synchronization of Oscillators

Huygens, 1665





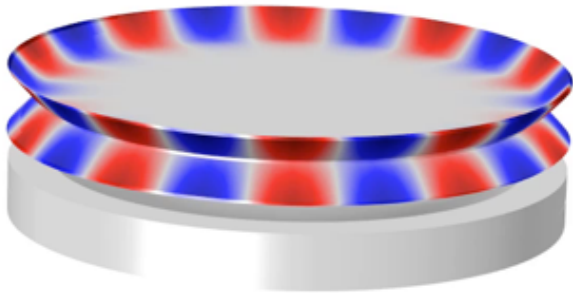
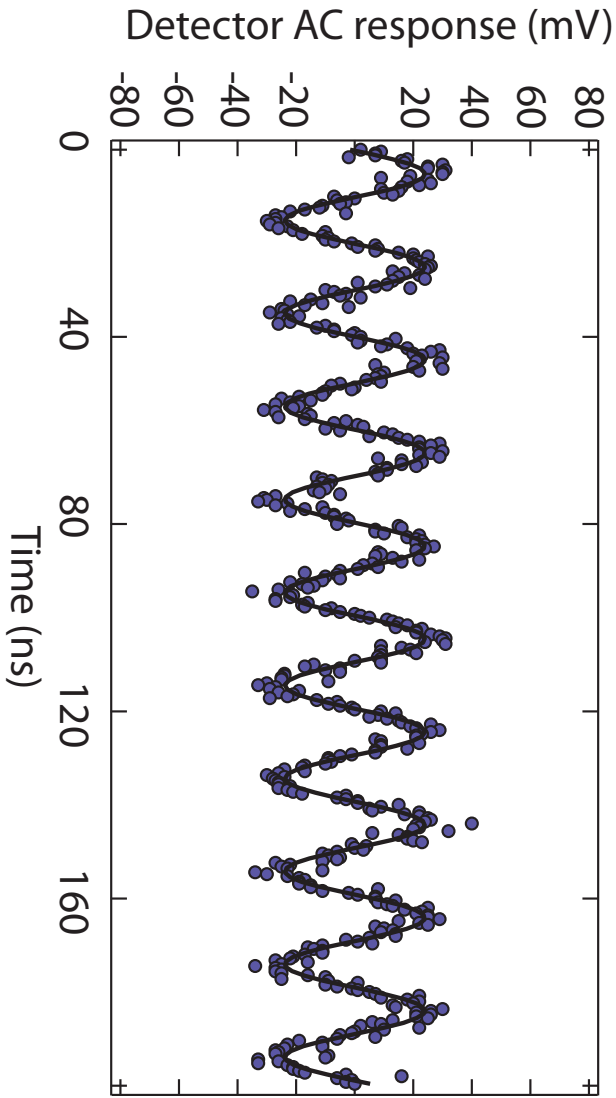
Individual Optomechanical Oscillations



$$f_L = 50.35 \text{ MHz}$$

$$f_L = 50.30 \text{ MHz}$$

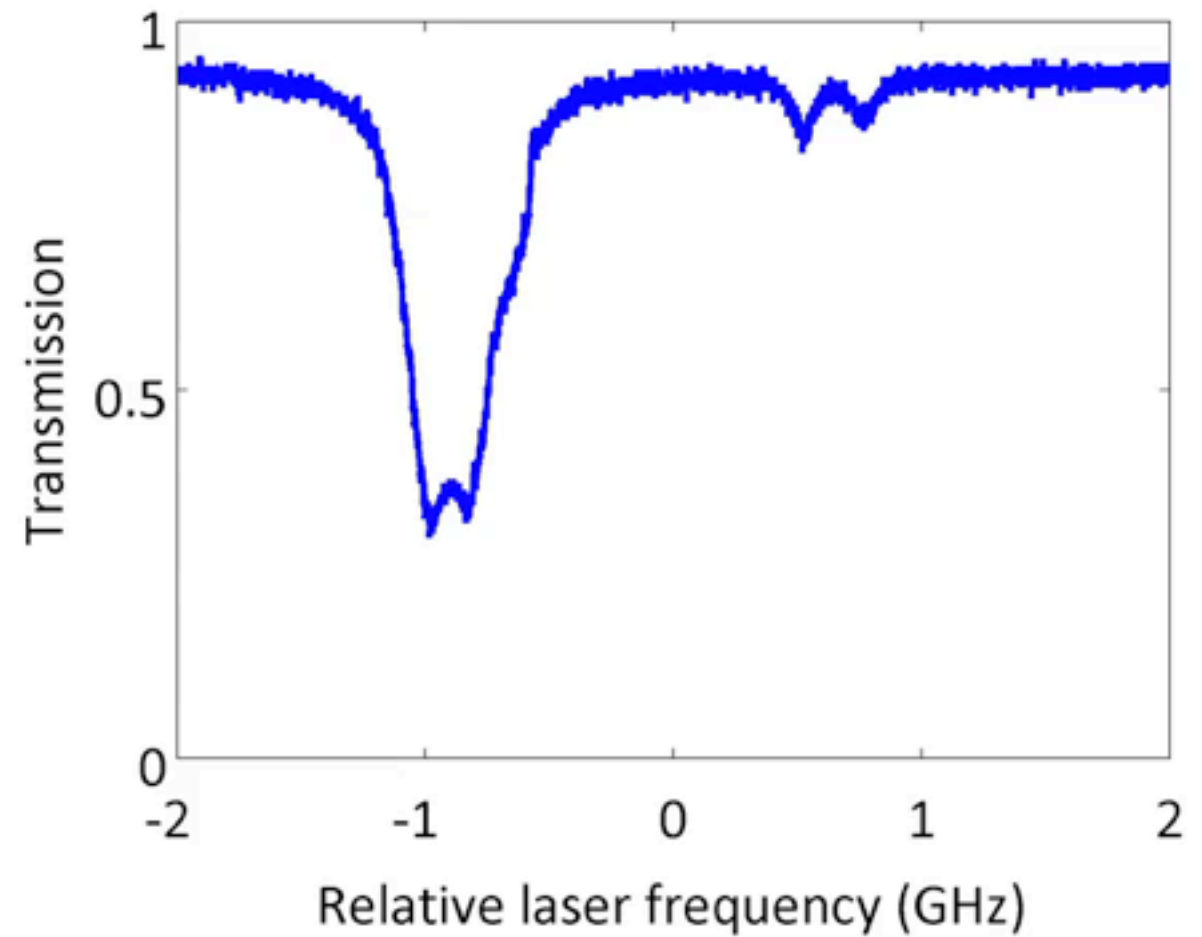
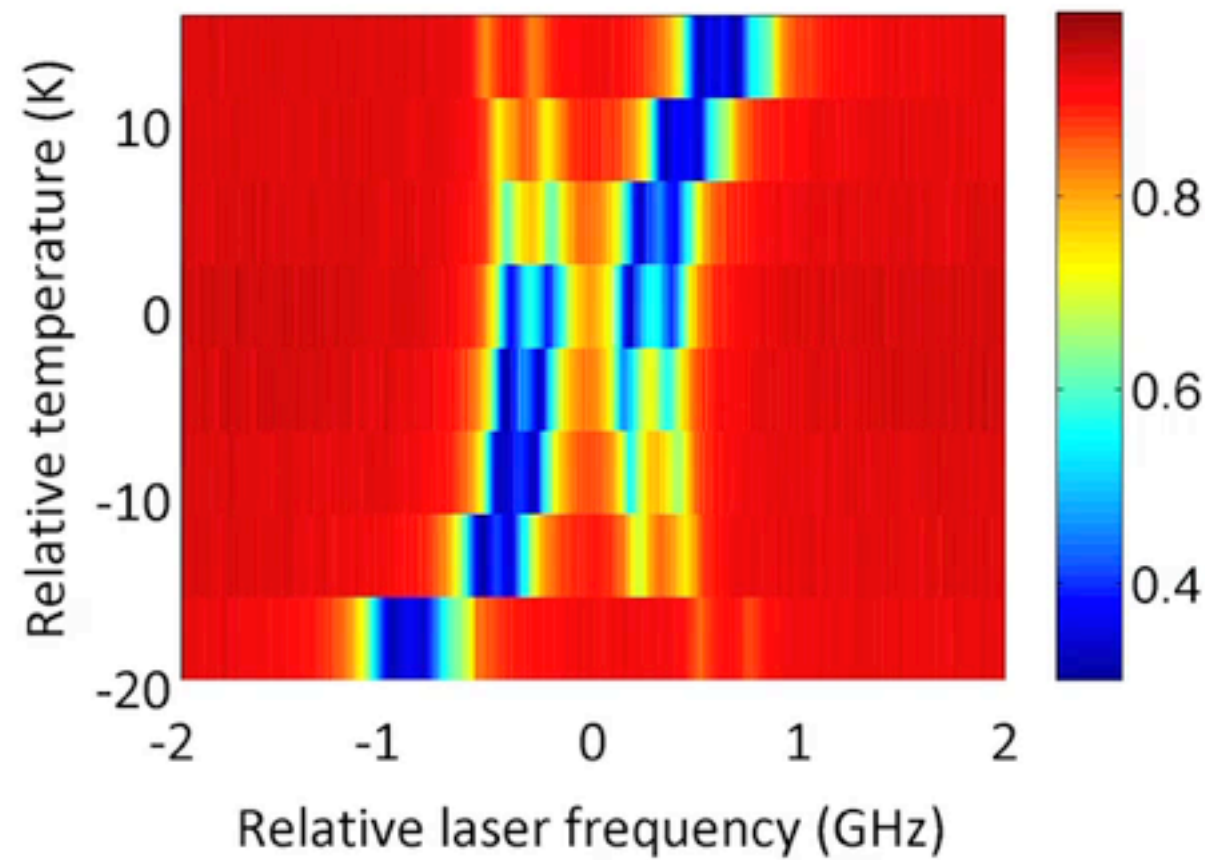
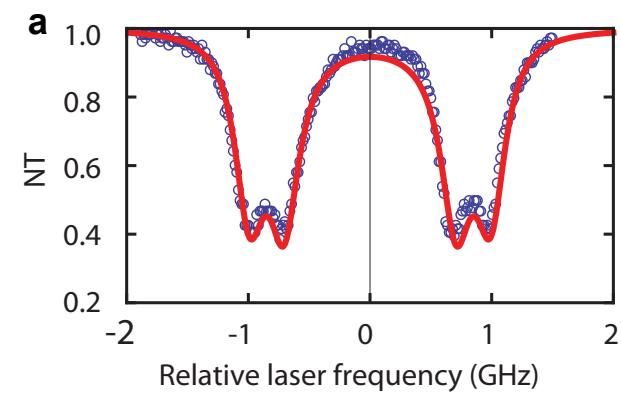
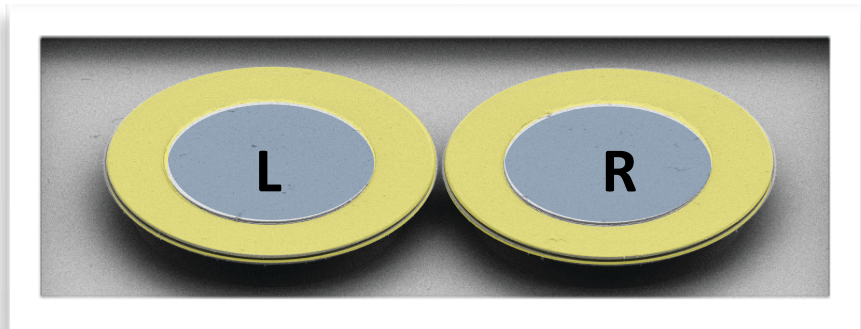
$$\delta f \approx 50 \text{ KHz}$$



50 MHz mechanical mode

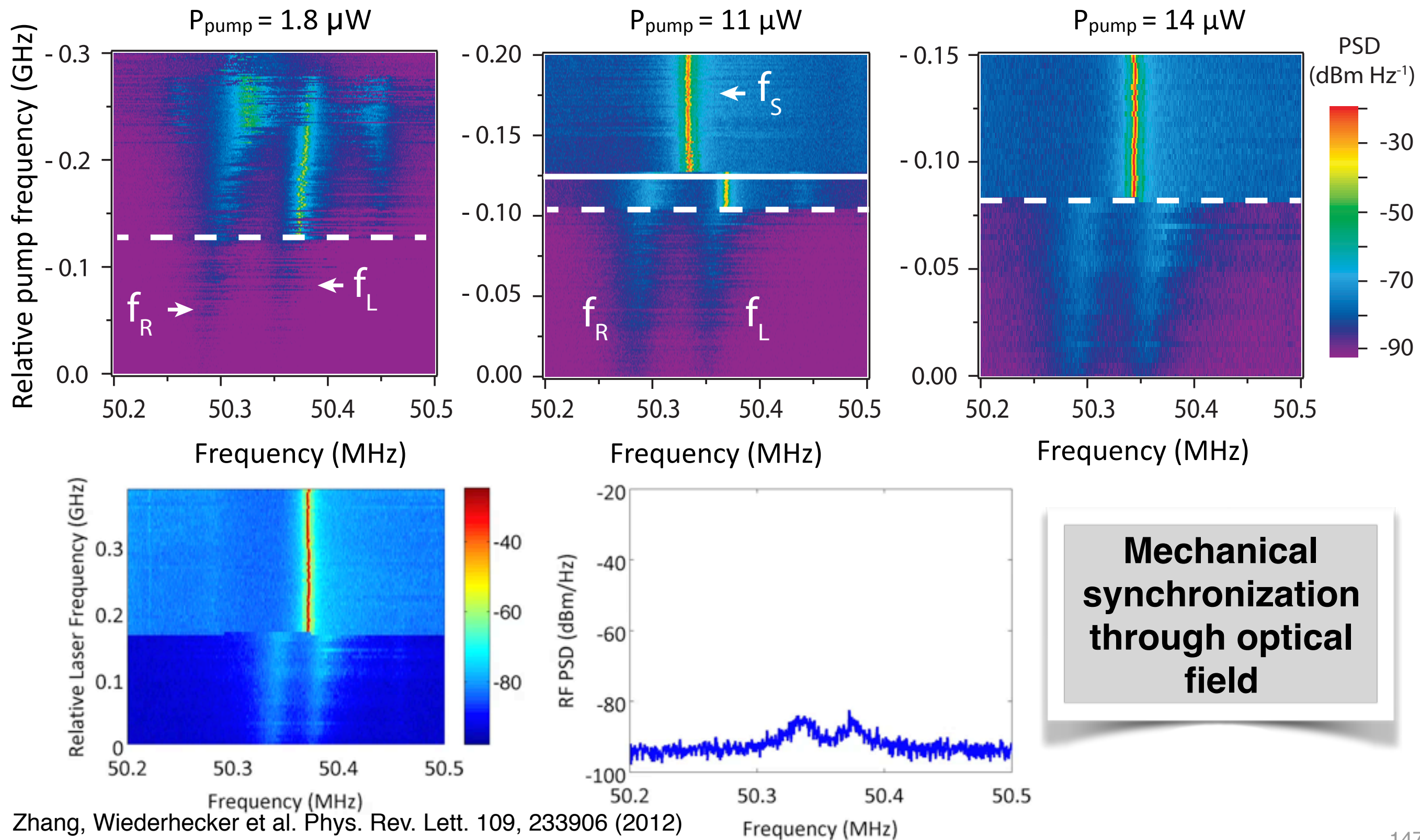


Tuning the Optical Coupling





Synchronization





Optically-induced mechanical coupling

$$\ddot{x}_1 + \underbrace{\gamma(n_p, \Delta)}_{\text{Optomechanical damping}} \dot{x}_1 + \underbrace{\Omega^2(n_p, \Delta)}_{\text{Optical Spring}} x_1 = \underbrace{-k_{\text{in}} x_2 + k_{\text{quad}} \dot{x}_2}_{\text{Linearized coupling}}$$
$$\ddot{x}_2 + \underbrace{\gamma(n_p, \Delta)}_{\text{Optomechanical damping}} \dot{x}_2 + \underbrace{\Omega^2(n_p, \Delta)}_{\text{Optical Spring}} x_2 = \underbrace{-k_{\text{in}} x_1 + k_{\text{quad}} \dot{x}_1}_{\text{Linearized coupling}}$$

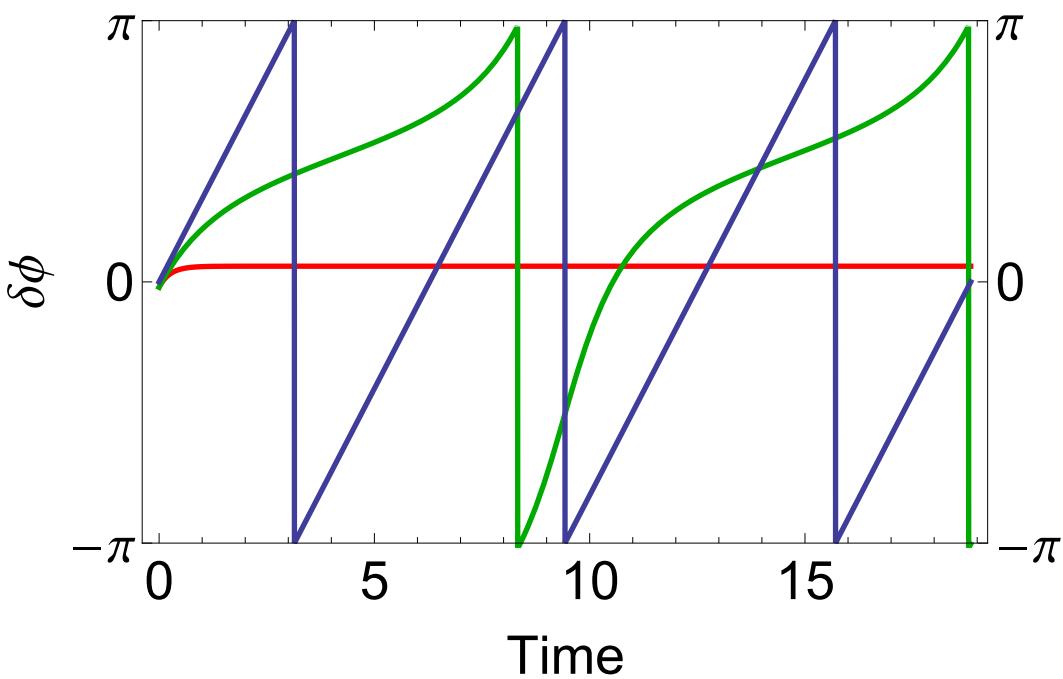
\Rightarrow

Kuramoto model

Periodic solutions $\delta\Omega > 2k_{\text{in}}$

Asymptotic solutions $\delta\Omega < 2k_{\text{in}}$

$$\delta\dot{\phi} = \delta\Omega - 2k_{\text{in}} \sin(\delta\phi)$$





Optically-induced mechanical coupling

Kuramoto model

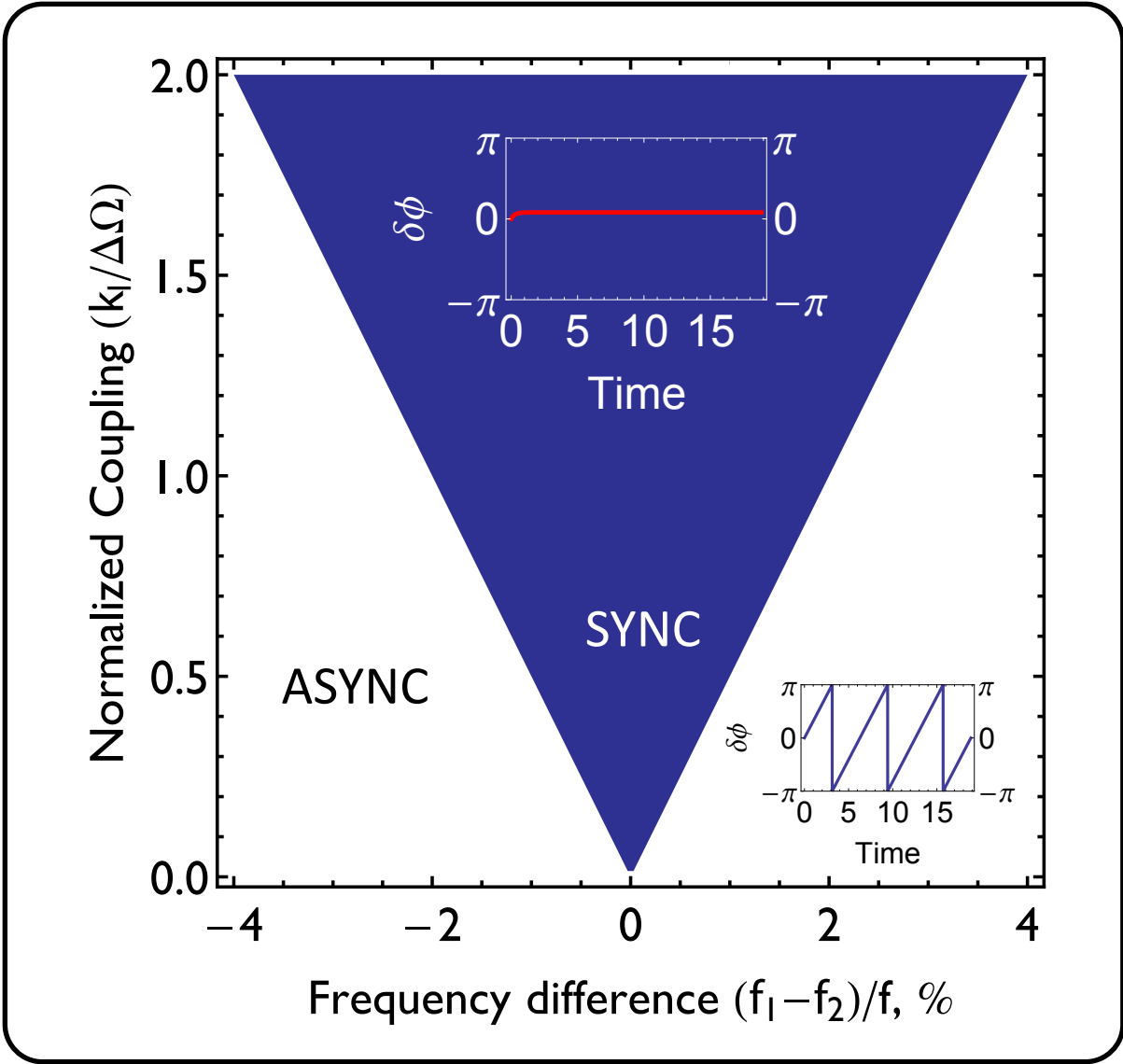
Periodic solutions

$\delta\Omega > 2k_{in}$

Asymptotic solutions

$\delta\Omega < 2k_{in}$

Arnold Tongue





Optically-induced mechanical coupling

Kuramoto model

Periodic solutions

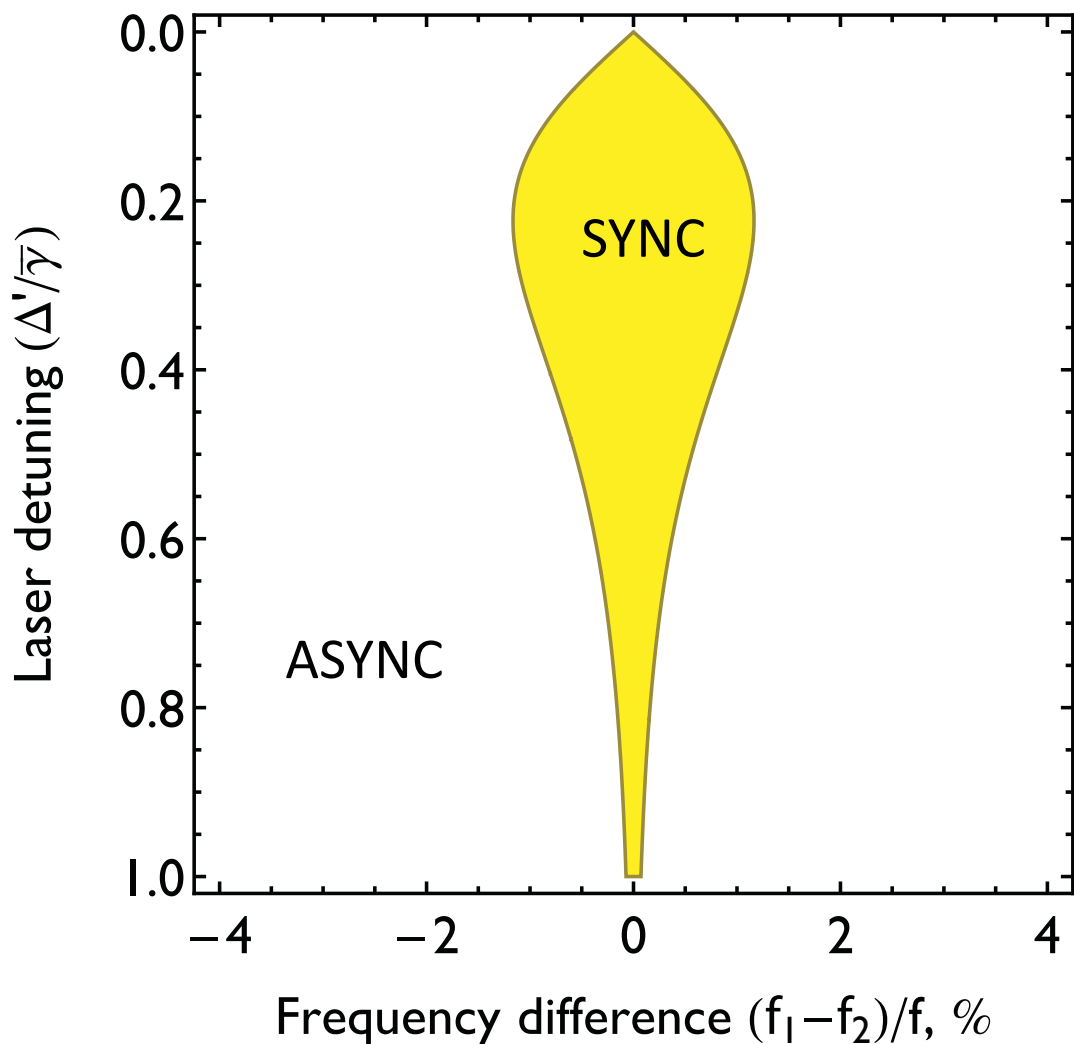
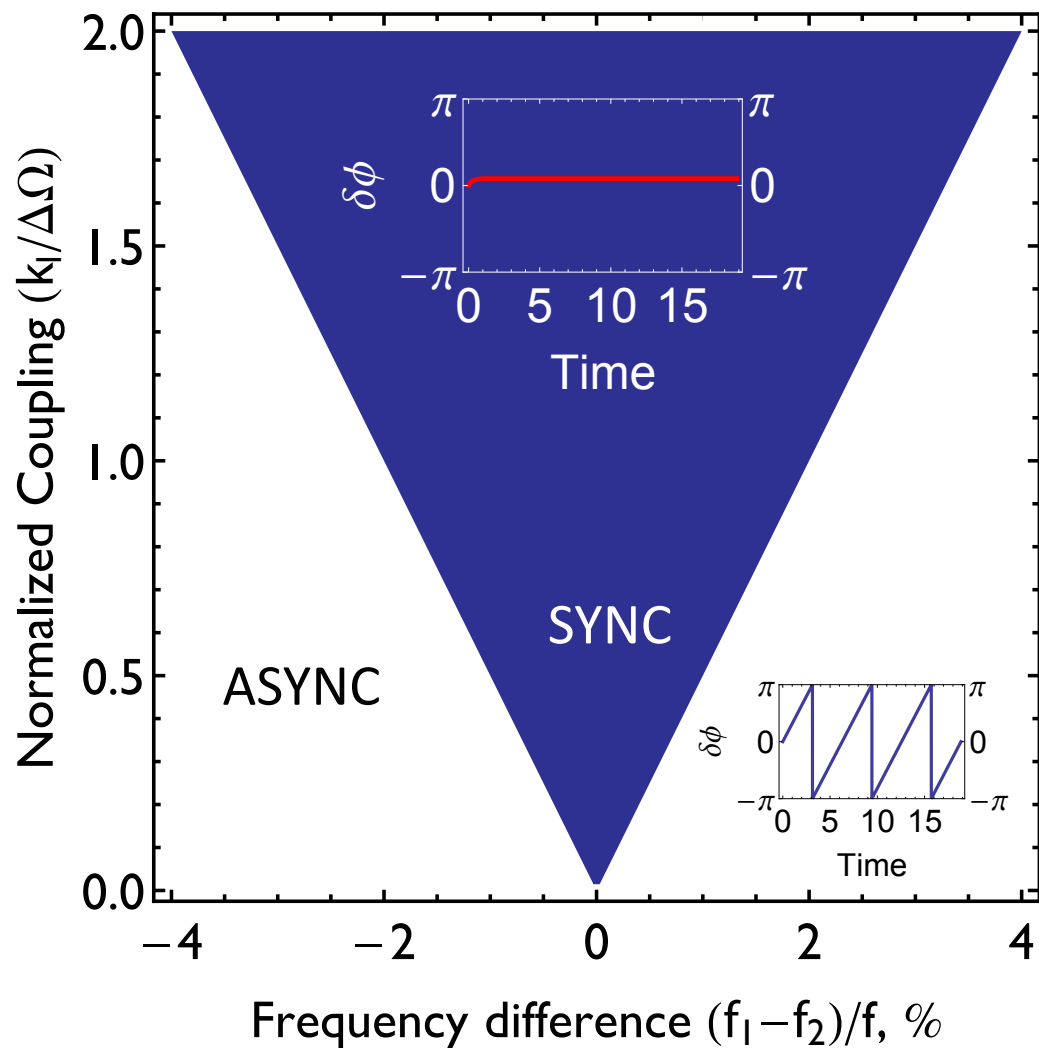
$\delta\Omega > 2k_{in}$

Asymptotic solutions

$\delta\Omega < 2k_{in}$

$$k_{in} \propto n_p \frac{\Delta}{((\kappa / 2)^2 + \Delta^2)^2}$$

Arnold Tongue





Optically-induced mechanical coupling

Kuramoto model

Periodic solutions

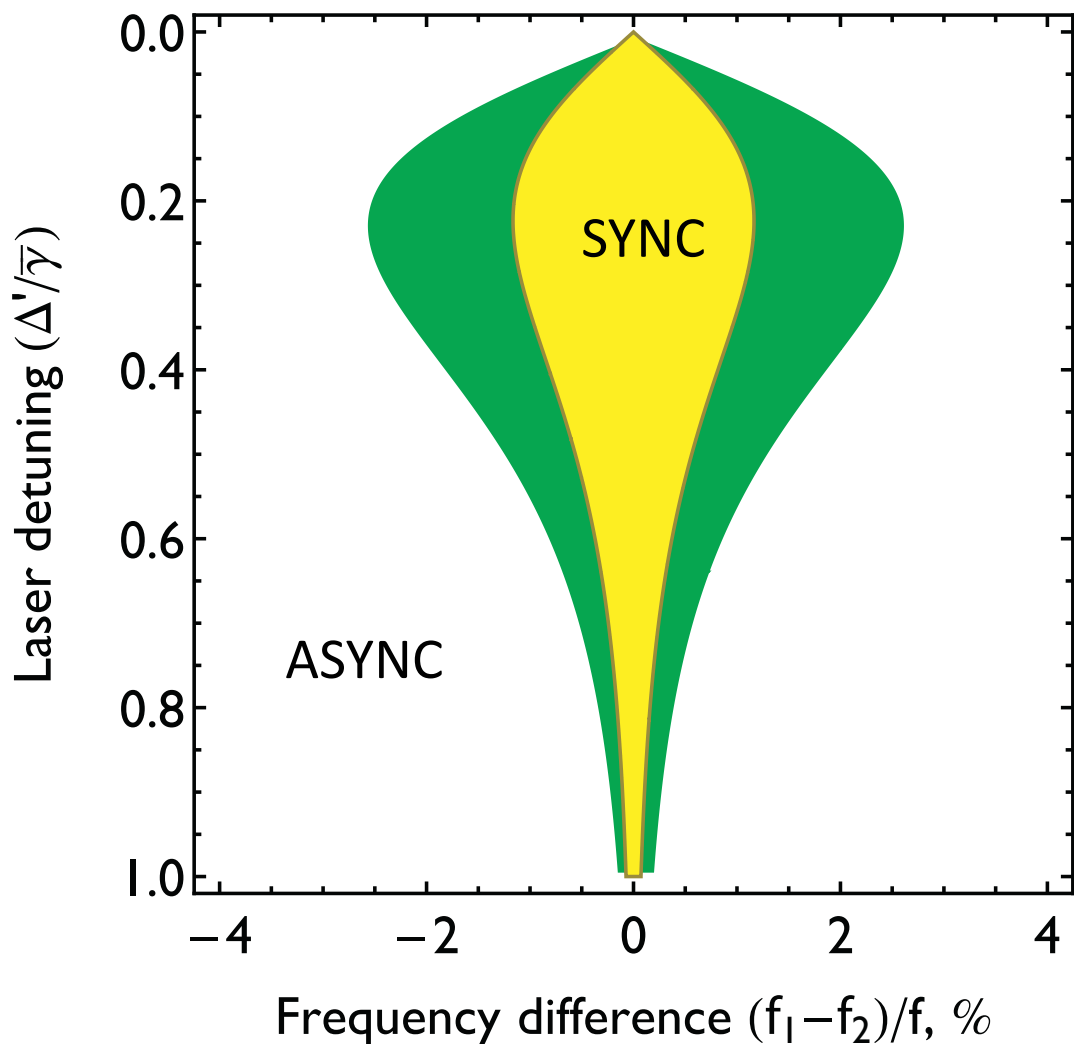
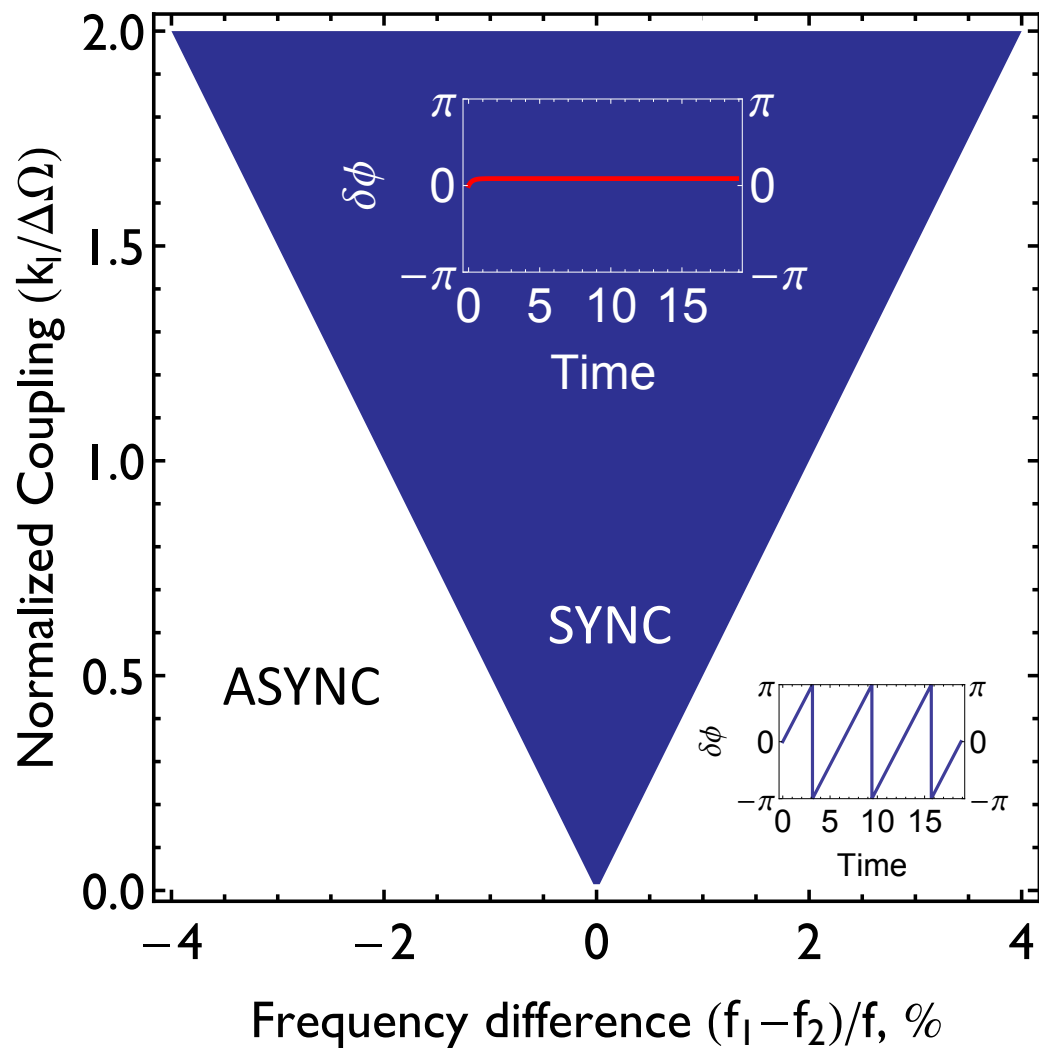
$\delta\Omega > 2k_{in}$

Asymptotic solutions

$\delta\Omega < 2k_{in}$

$$k_{in} \propto n_p \frac{\Delta}{((\kappa / 2)^2 + \Delta^2)^2}$$

Arnold Tongue





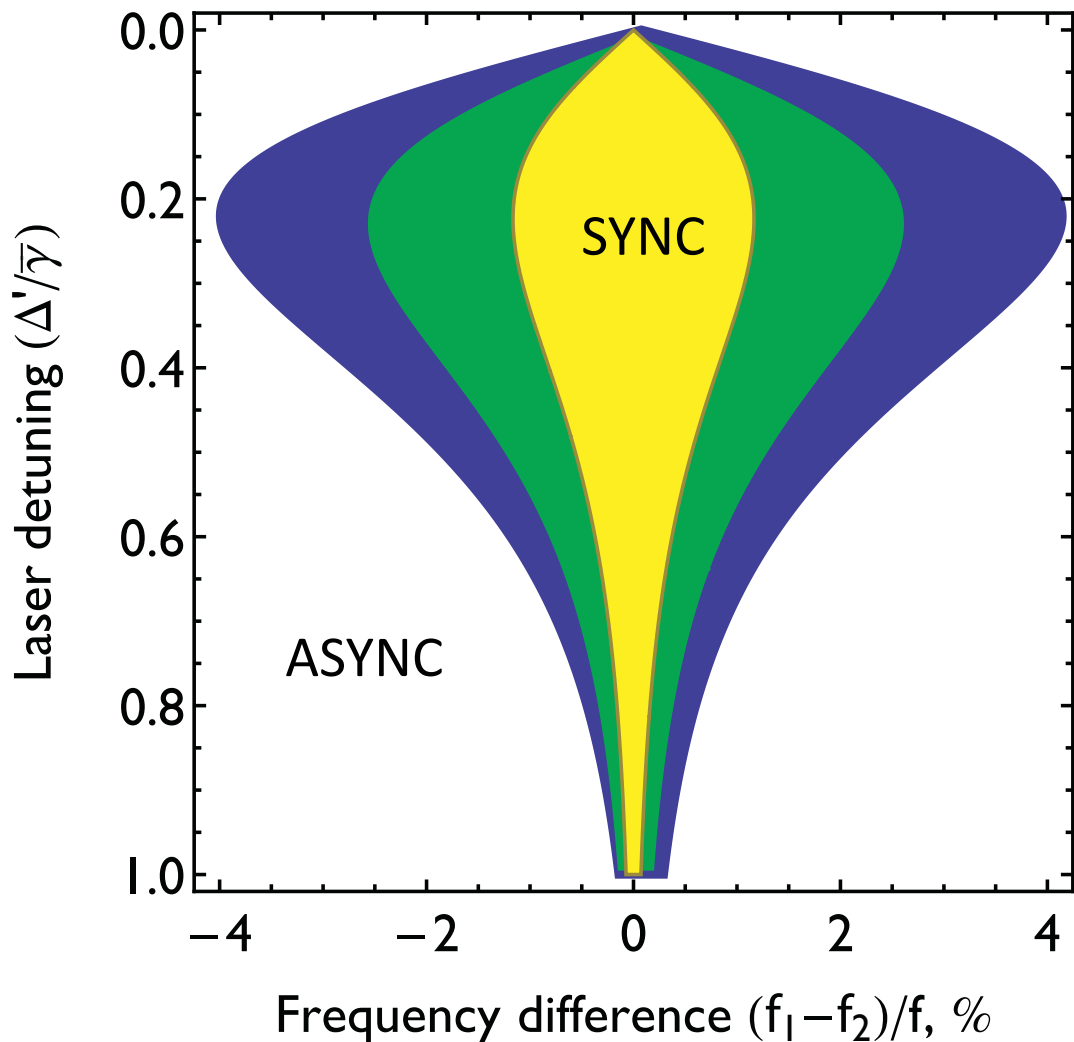
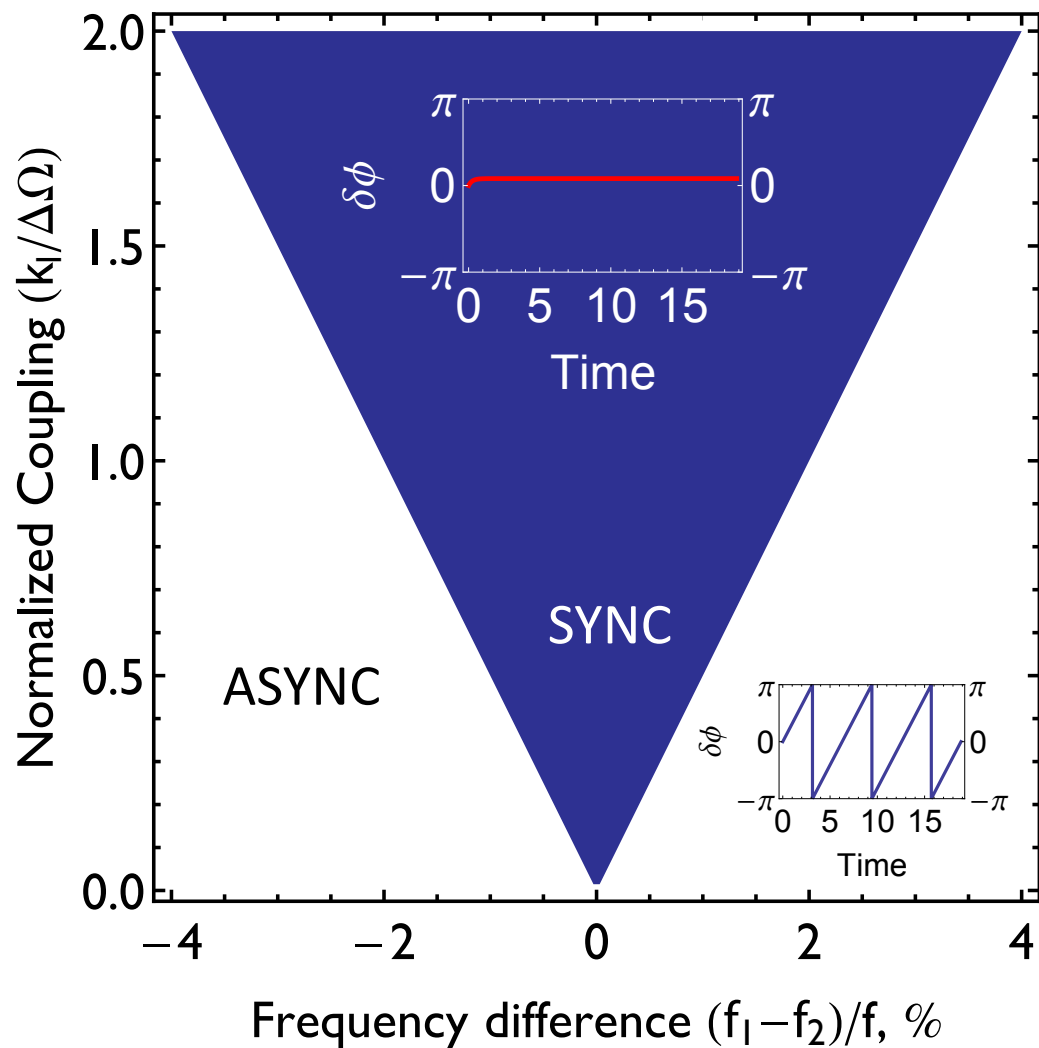
Optically-induced mechanical coupling

Kuramoto model

Periodic solutions	$\delta\Omega > 2k_{in}$
Asymptotic solutions	$\delta\Omega < 2k_{in}$

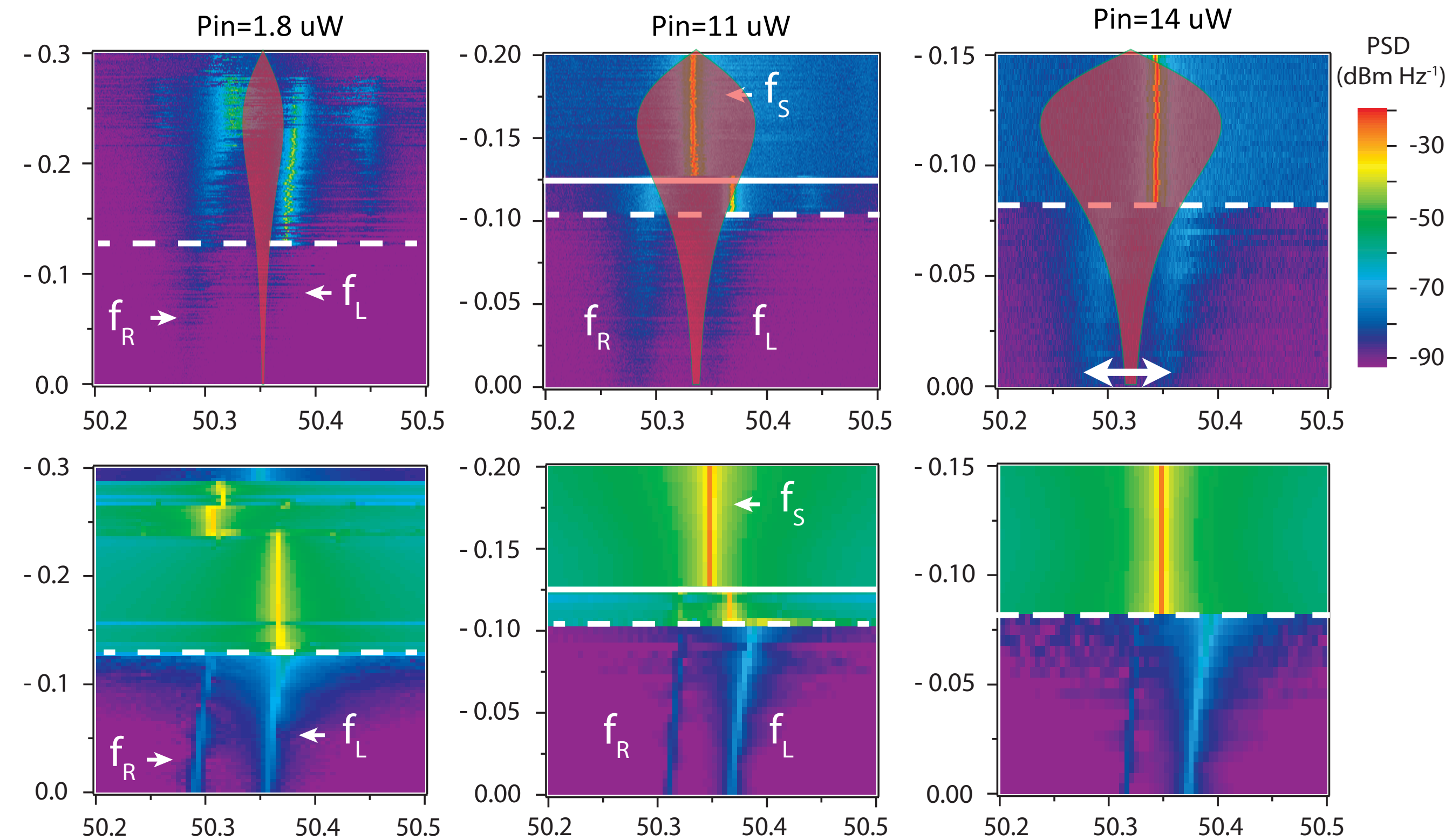
$$k_{in} \propto n_p \frac{\Delta}{((\kappa / 2)^2 + \Delta^2)^2}$$

Arnold Tongue





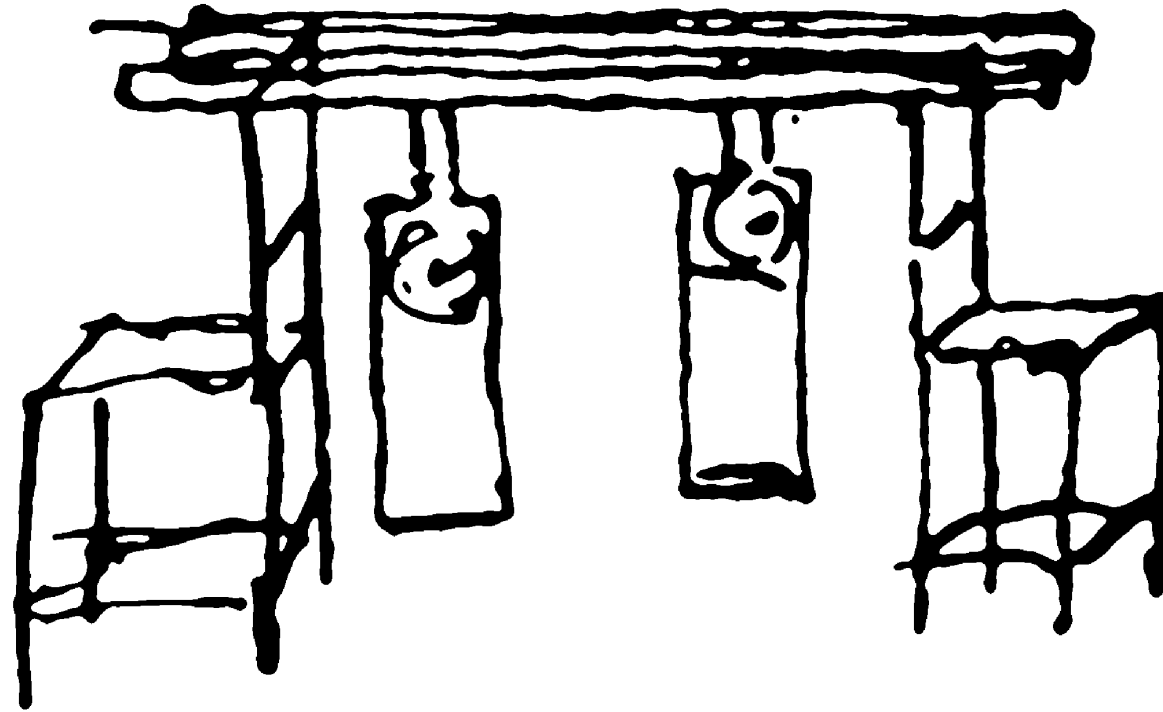
Synchronization





Synchronization of Oscillators

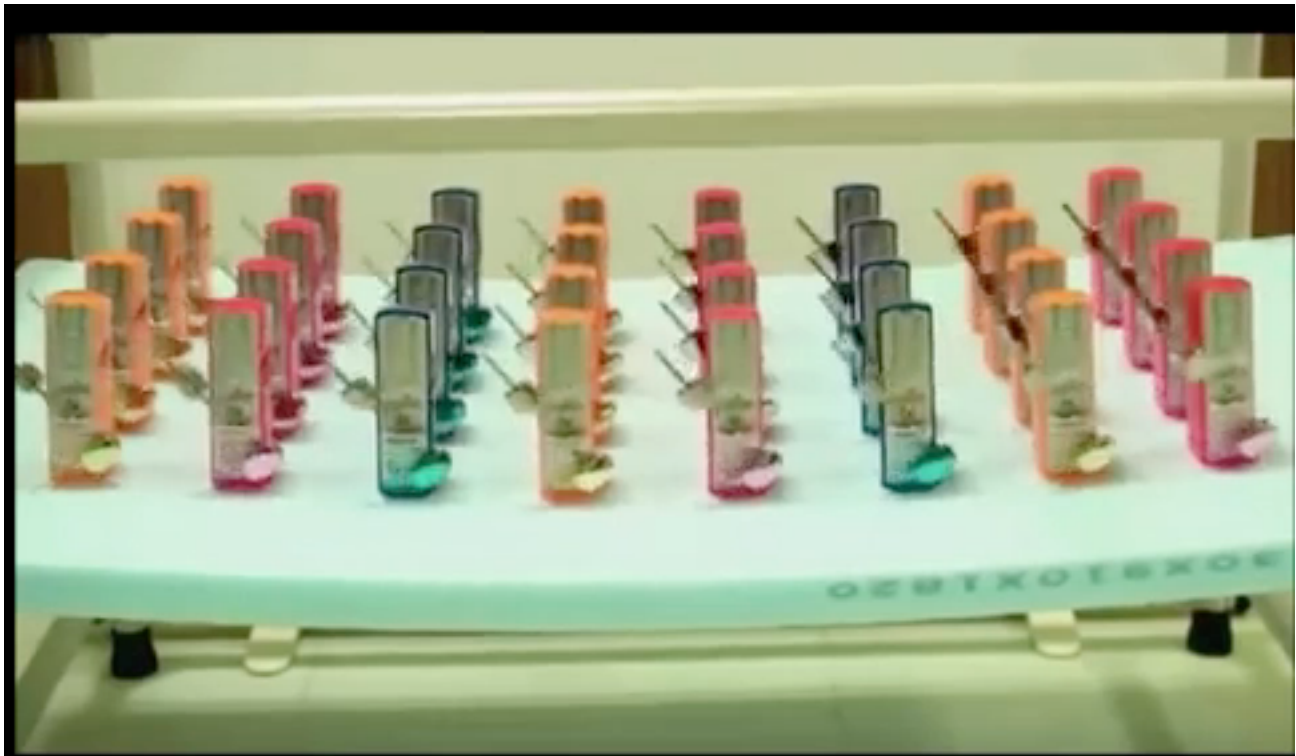
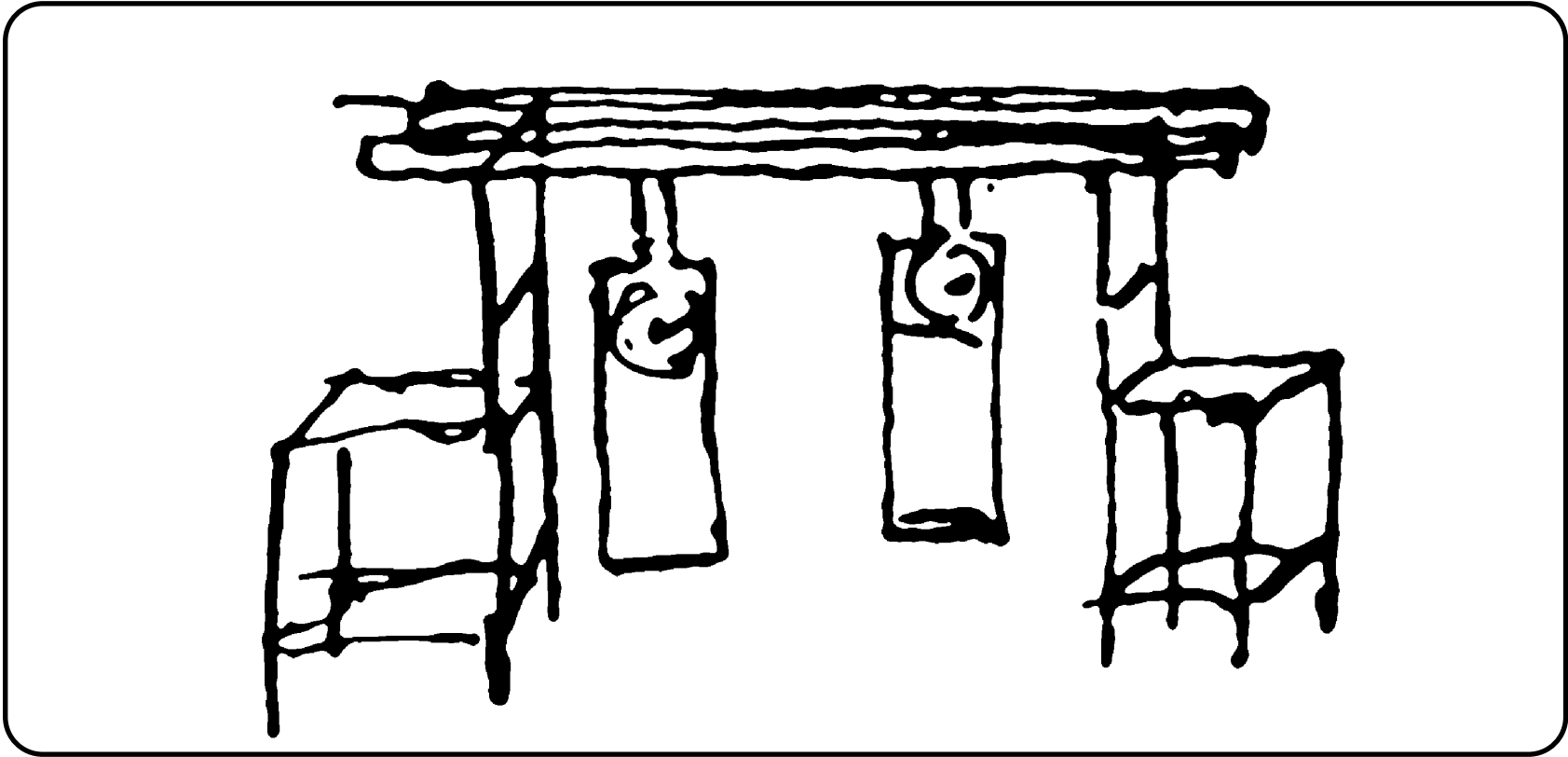
Huygens, 1665





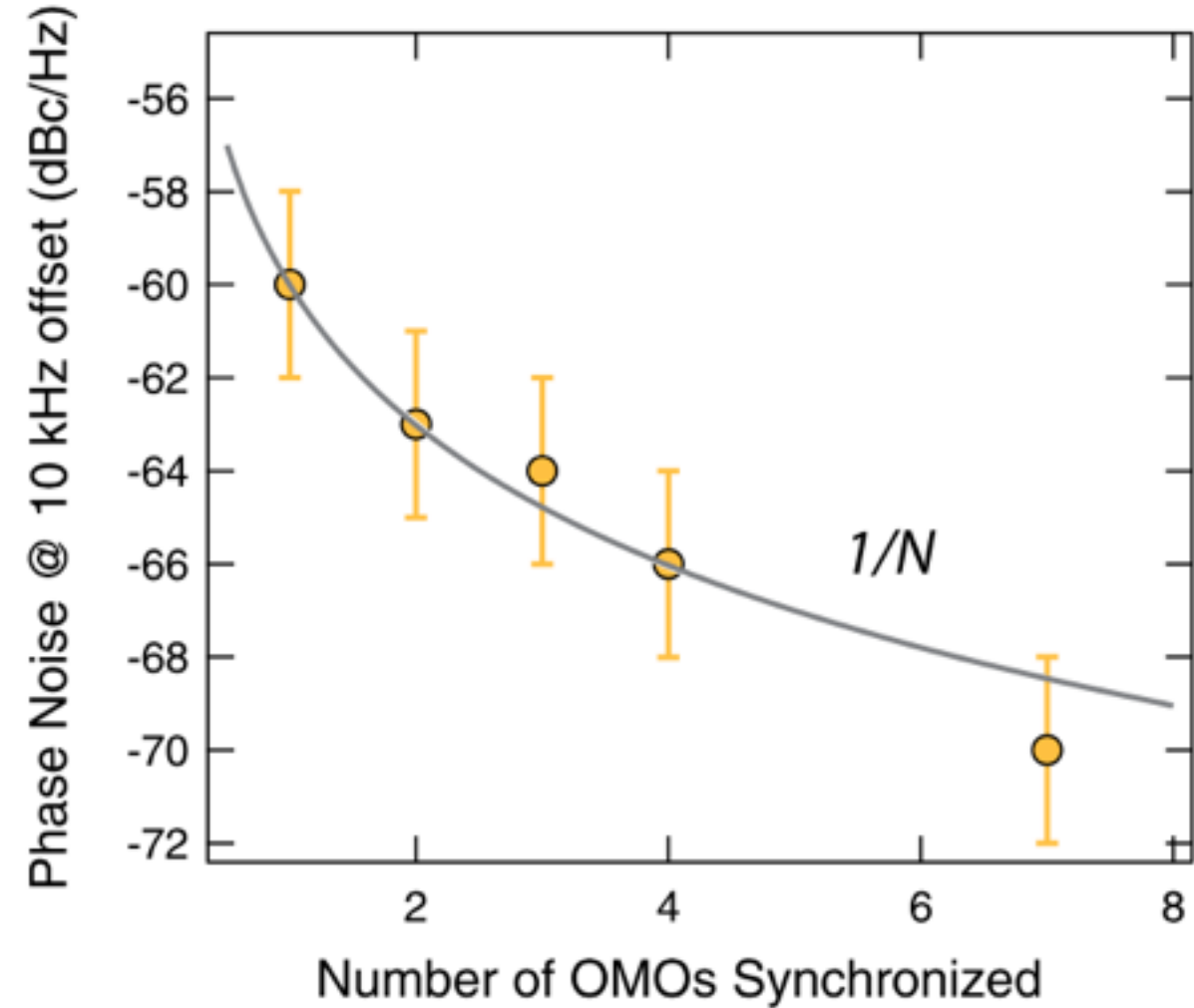
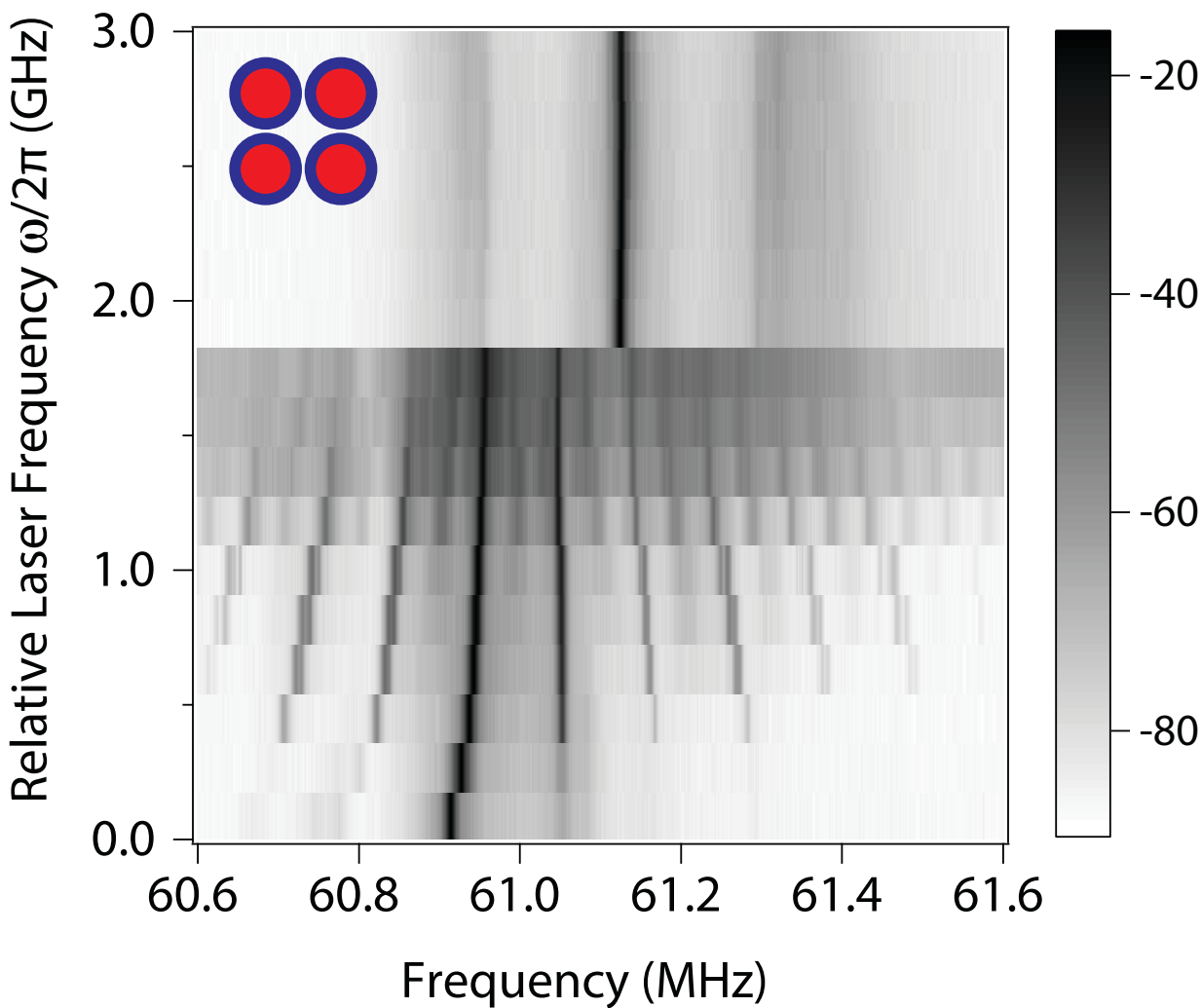
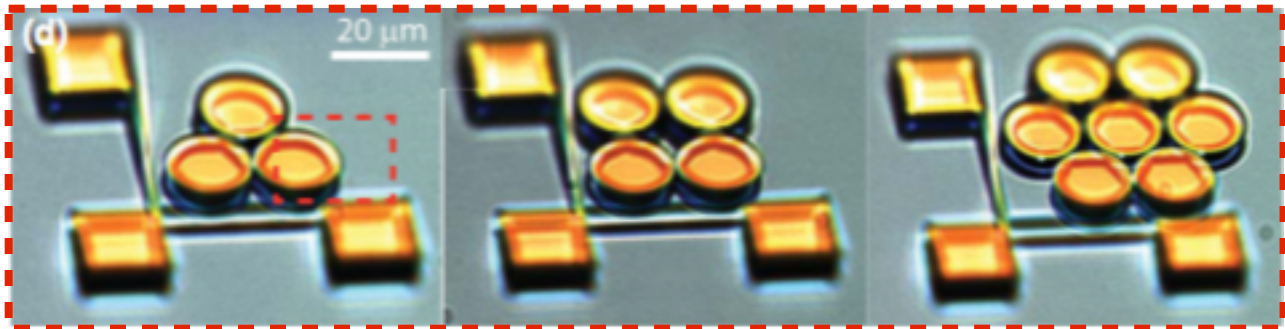
Synchronization of Oscillators

Huygens, 1665





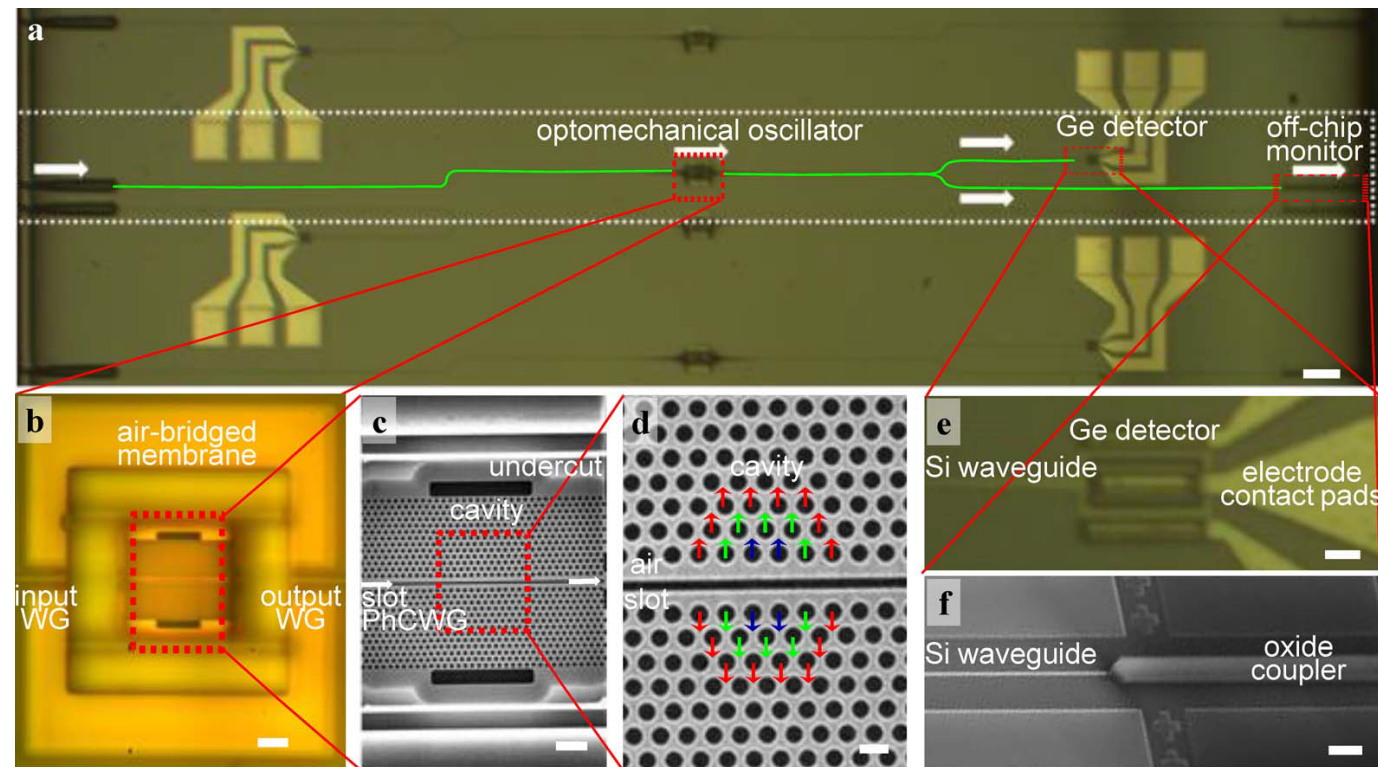
Array Synchronization



Zhang, M., et al (2015). PRL, 115(16), 163902.
Vahala, K. J. (2008). Physical Review A, 78(2), 023832.



Technological viewpoint: Si Nanophotonics Building Blocks



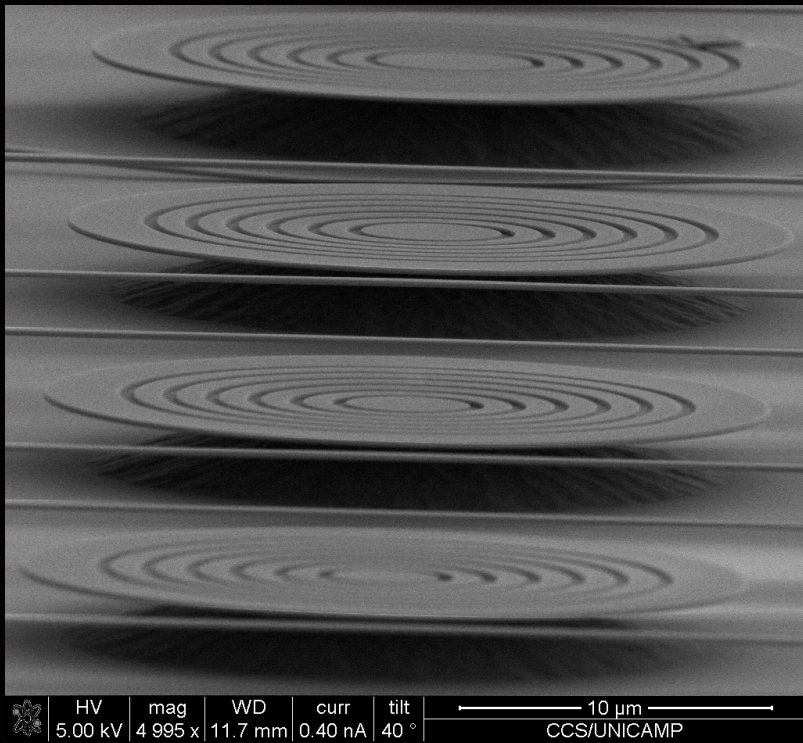
Integrated optomechanical oscillator chipset

Chee Wei Wong et. al SCIENTIFIC REPORTS 4 : 6842 (2014)

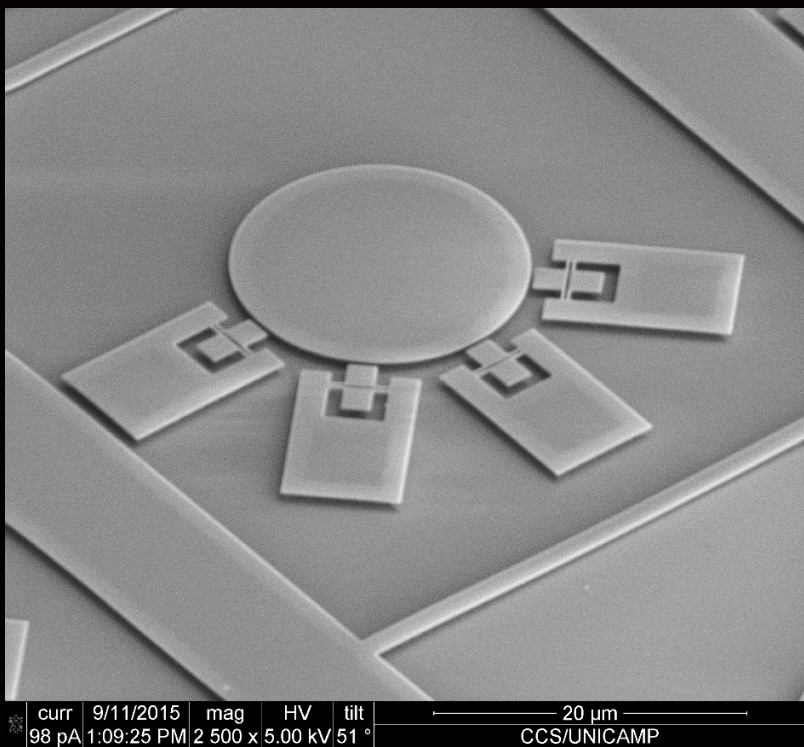


Candidates for OM arrays

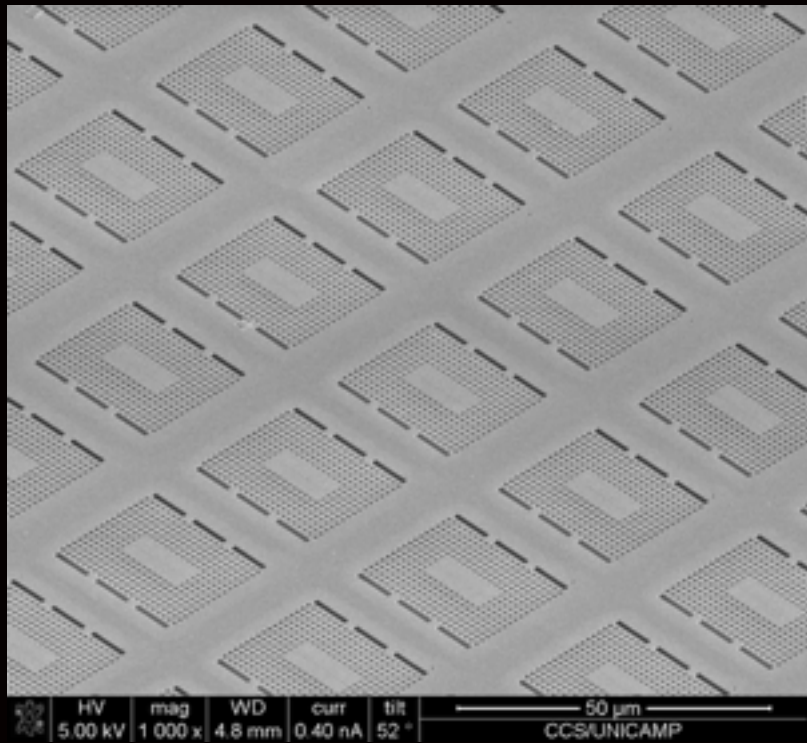
Bulls' eye cavity



Paddle resonators



Optomechanical crystals



Rodrigo Benevides



Gustavo Luiz



Felipe Santos

Benevides, et al. CLEO 2016
Luiz, et al CLEO 2106
Santos, F. G., et al. arXiv:1605.06318

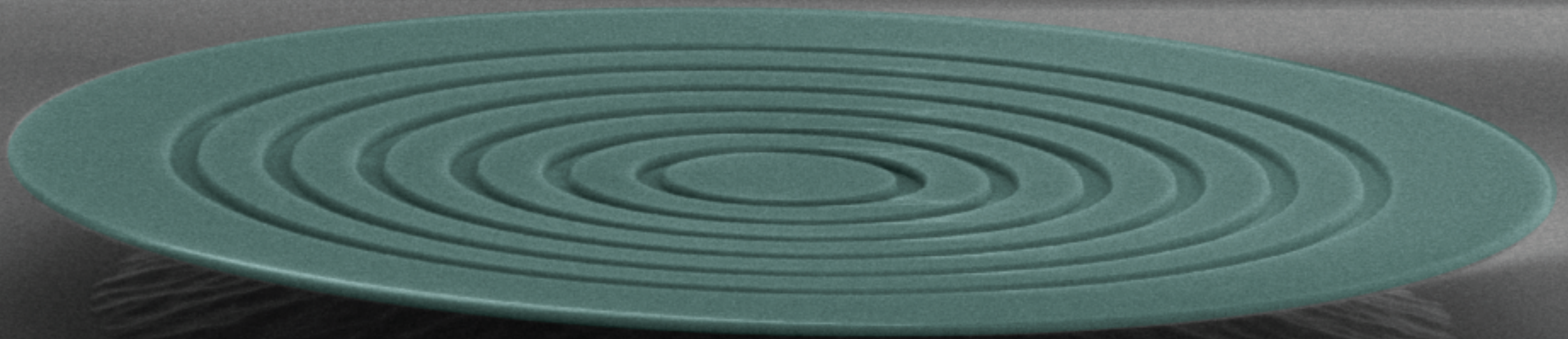


Outline

- ★ Optical and acoustic mode interaction
- ★ Optical force actuation
- ★ Dynamical back-action
- ★ Optomechanical clocks
- ★ **Bullseye - a case study**
- ★ Outlook



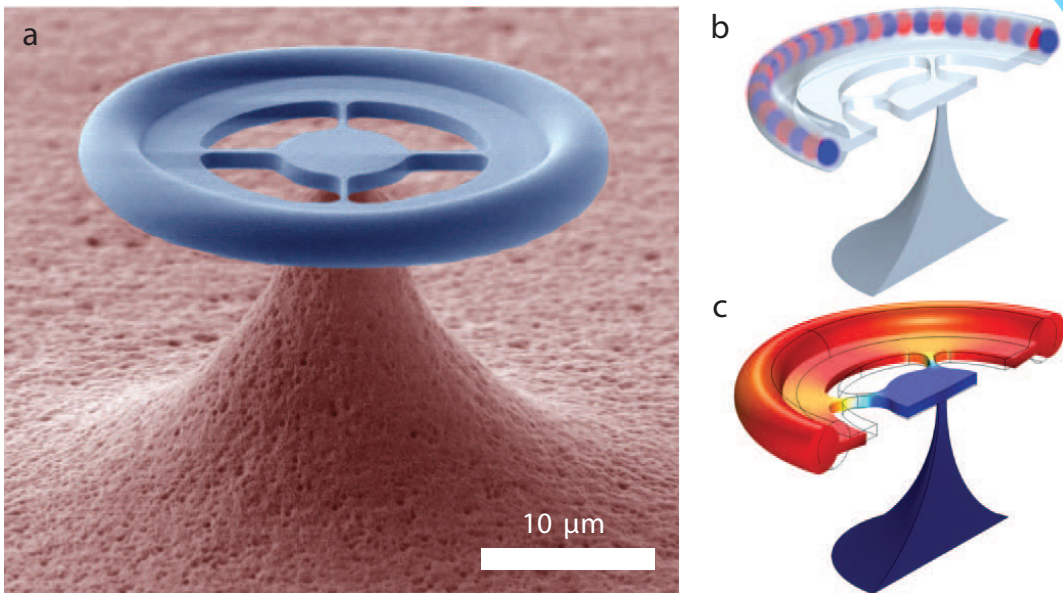
Bullseye





A new system for cavity optomechanics

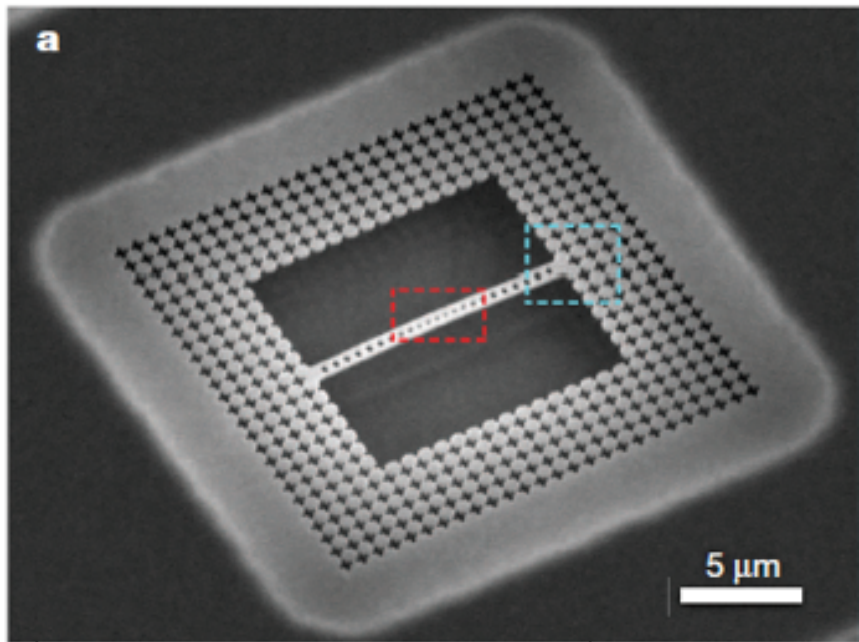
$\kappa \downarrow$



$\kappa_t/2\pi \approx 10 \text{ MHz}$ $\omega_m/2\pi \approx 80 \text{ MHz}$

E. Verhagen et. al, Nature **482**, 63 (2012)

$\Omega_m \uparrow$

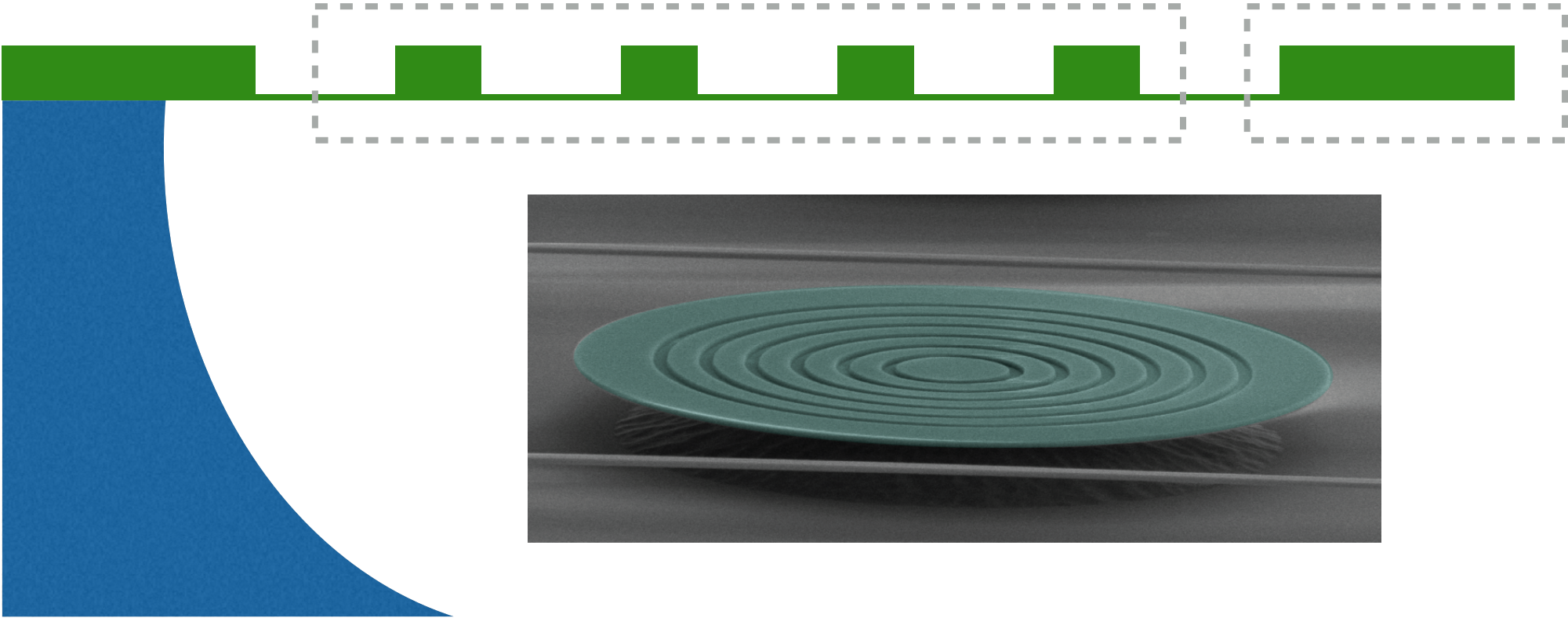


$\kappa_t/2\pi \approx 500 \text{ MHz}$ $\omega_m/2\pi \approx 4 \text{ GHz}$

J. Chan et. al, Nature **478**, 89 (2011)

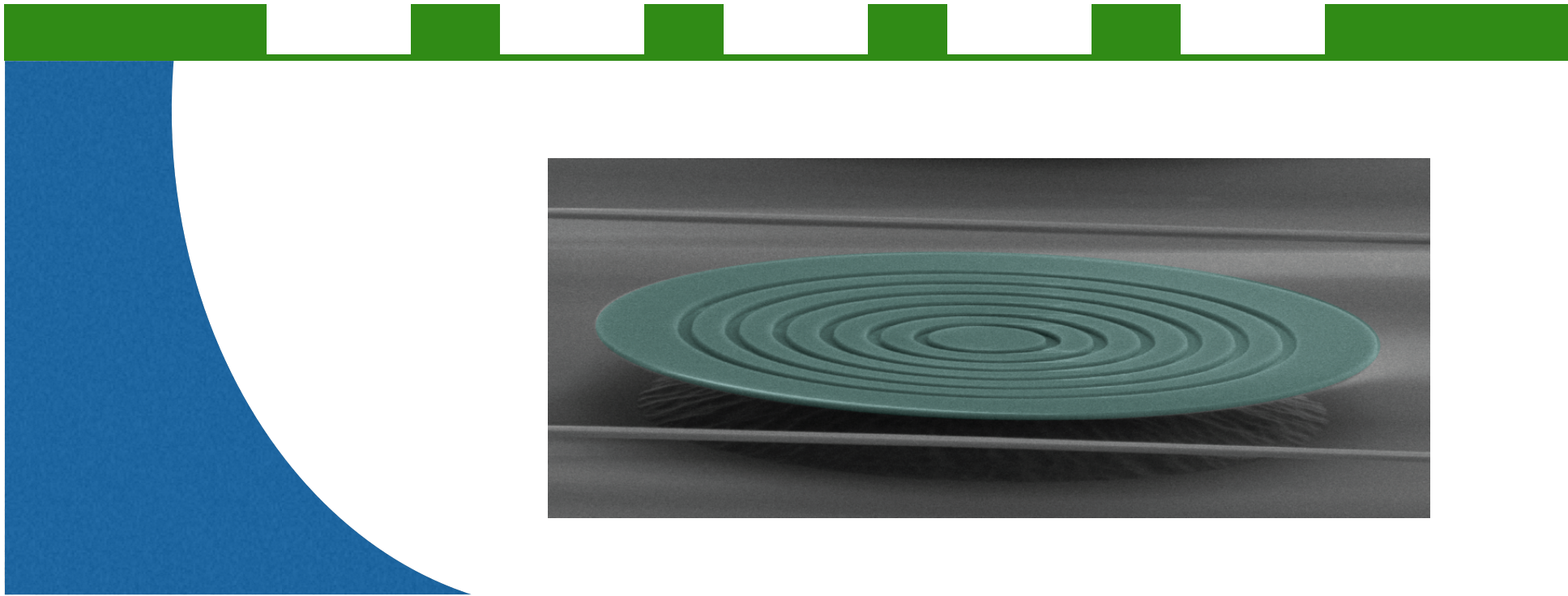


Optical Design

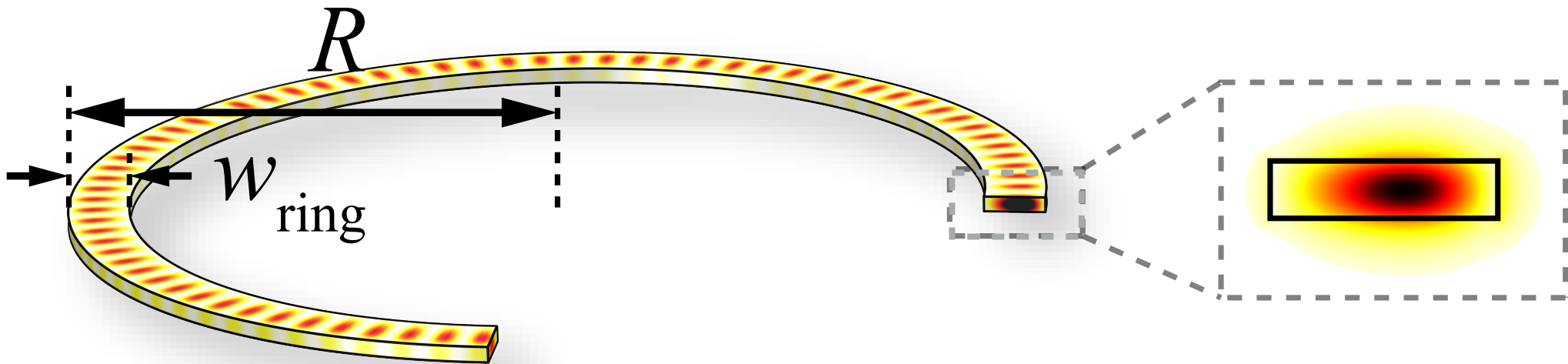




Optical Design

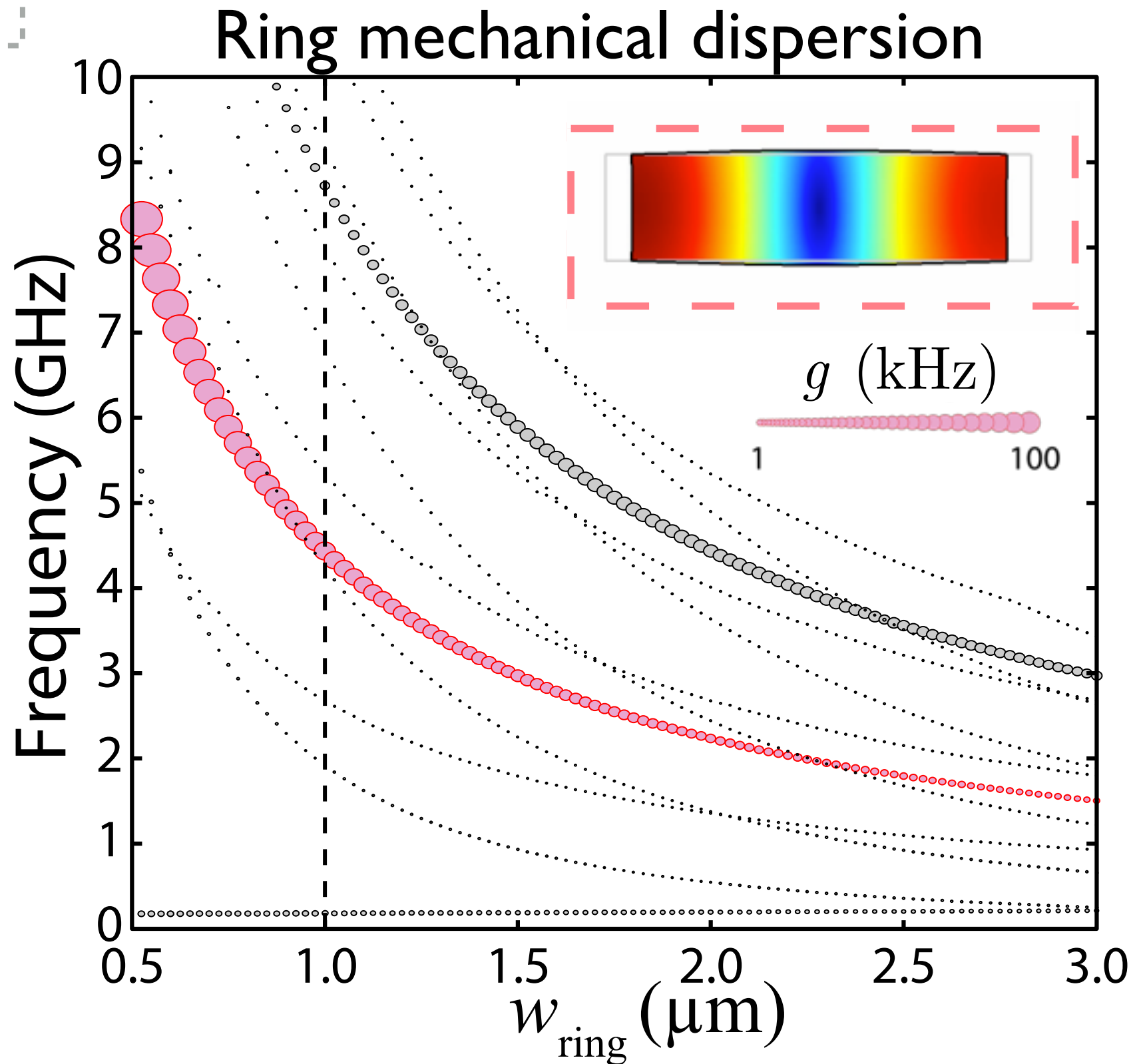
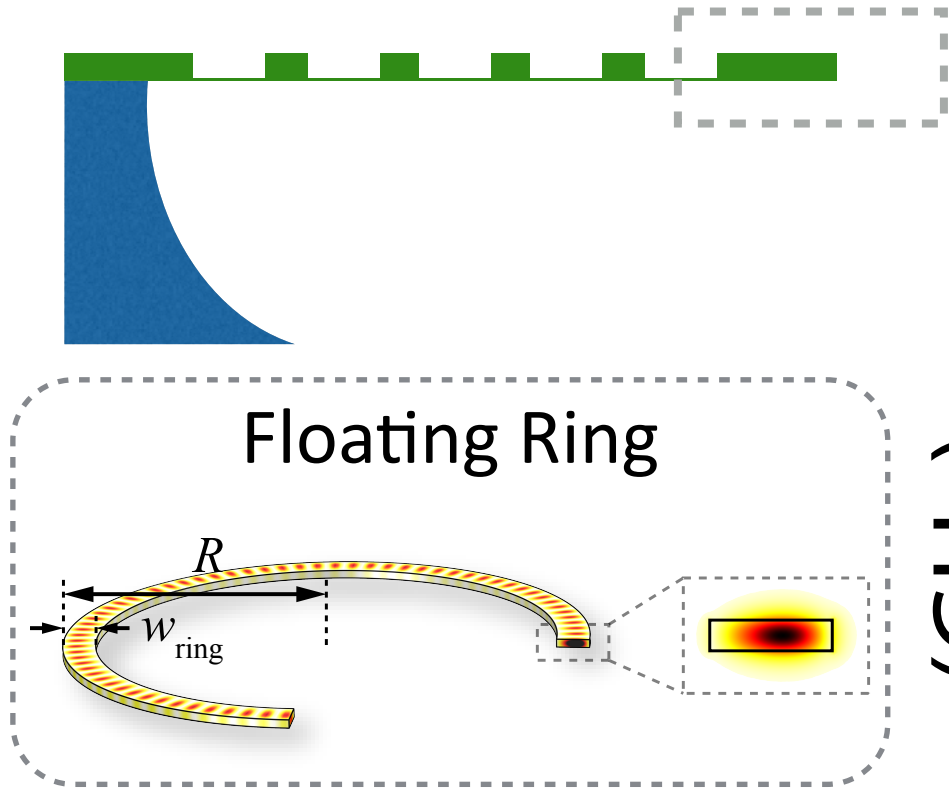


Floating ring resonator





Mechanical Design





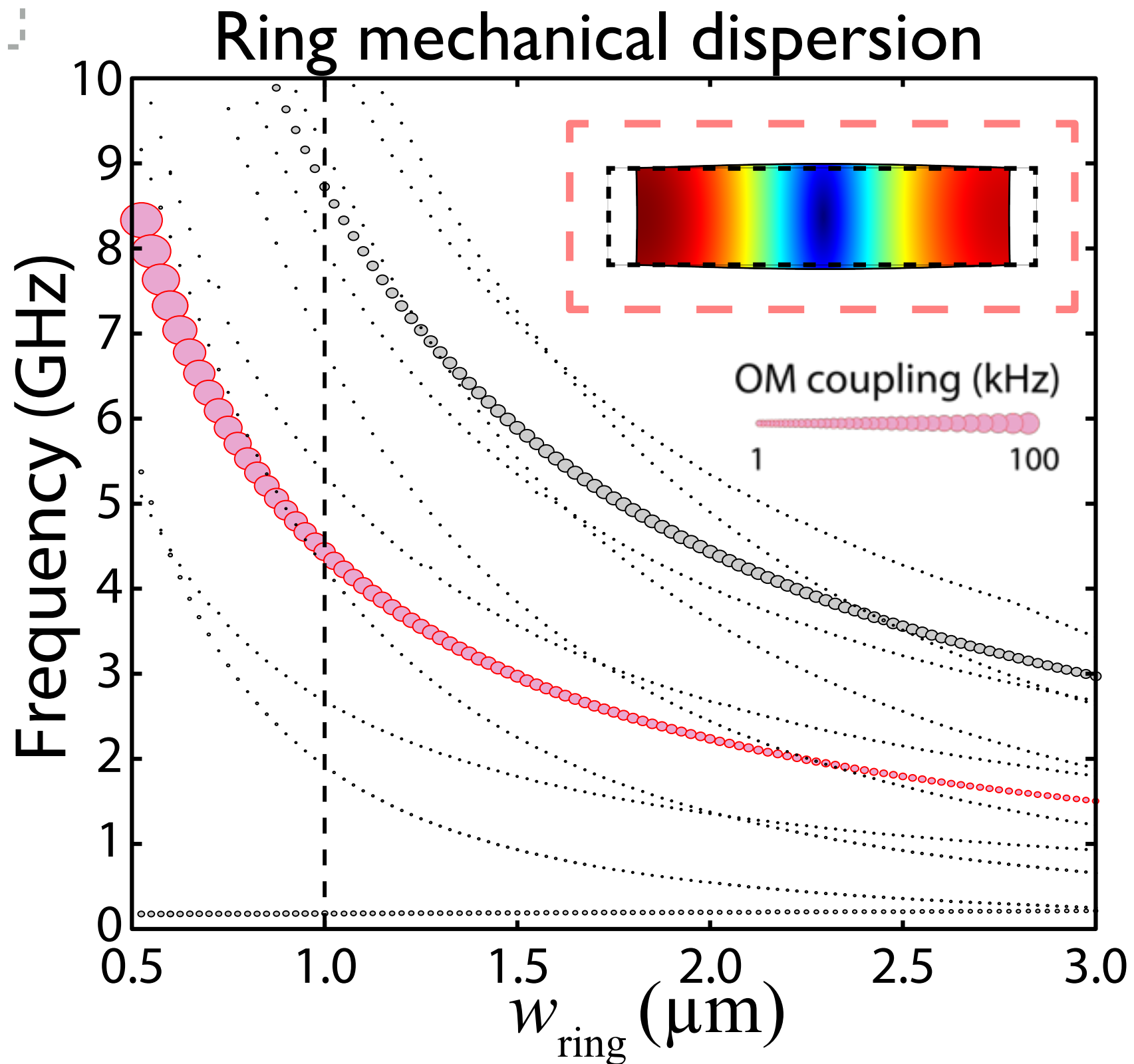
Mechanical Design

Floating Ring

High optomechanical coupling

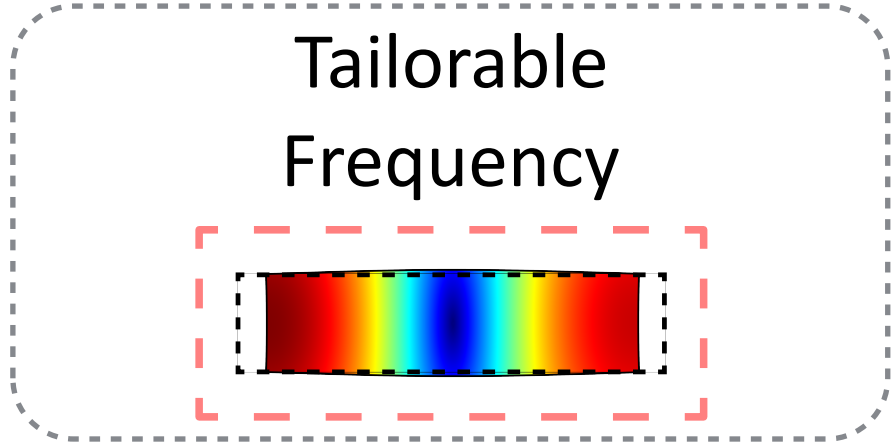
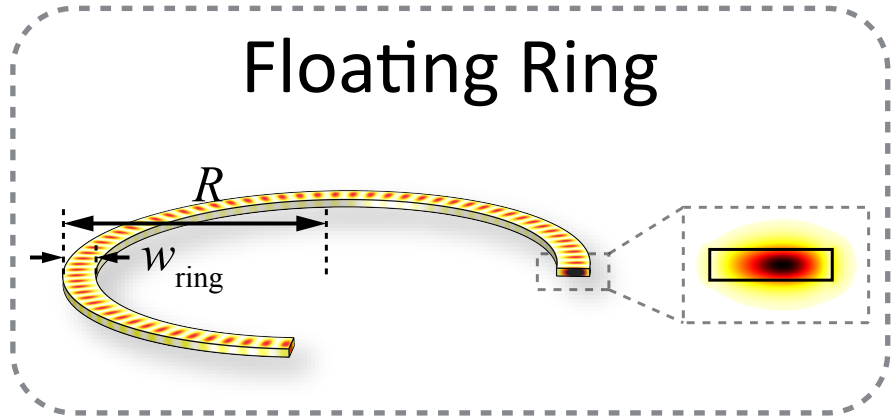
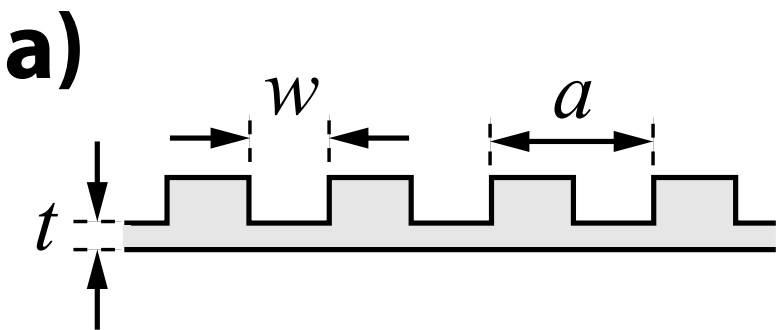
Photoelastic Effect

Tailorable mechanical frequency



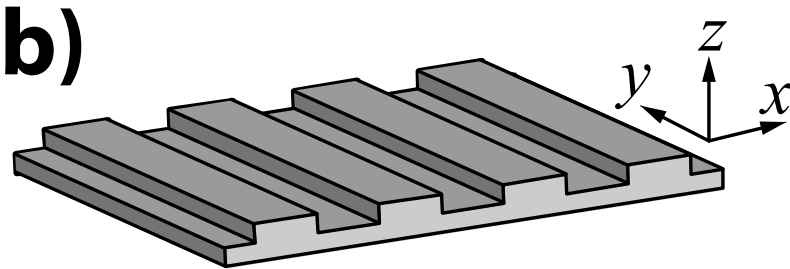
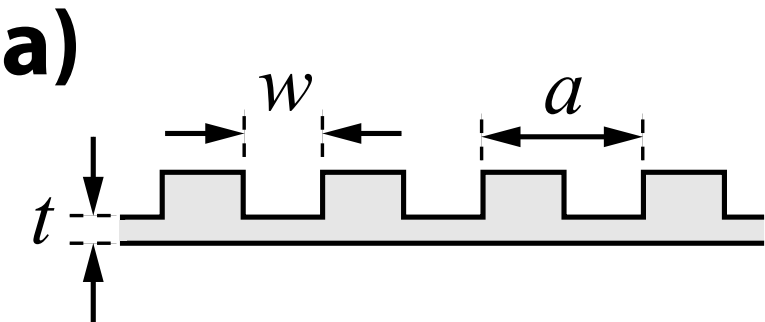
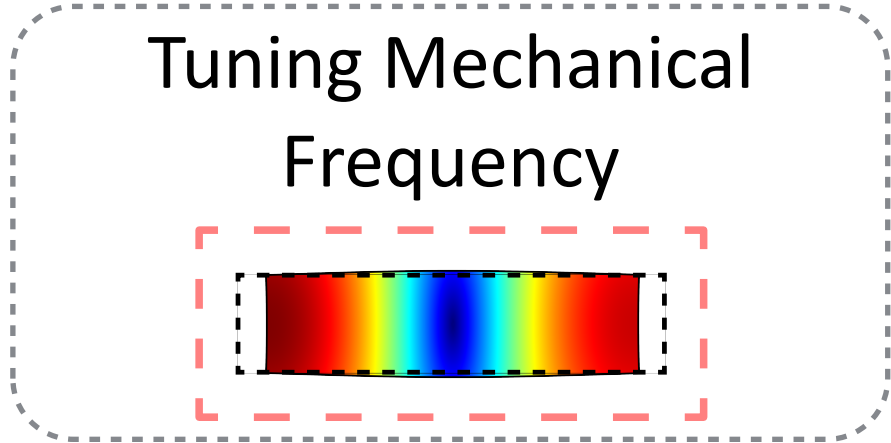
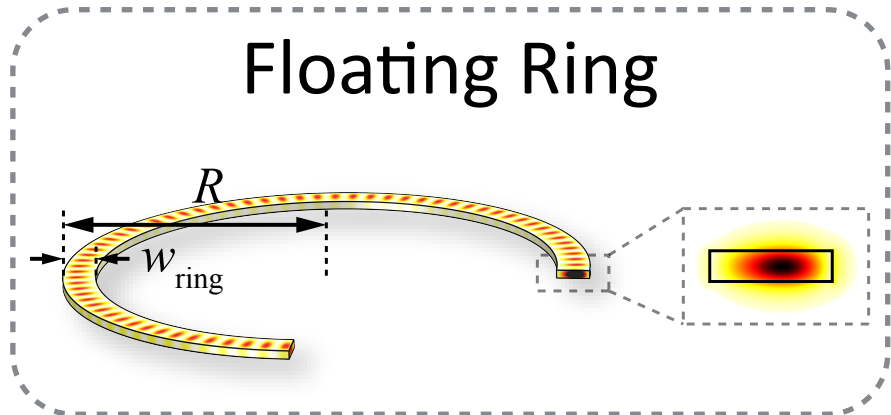


Mechanical Grating



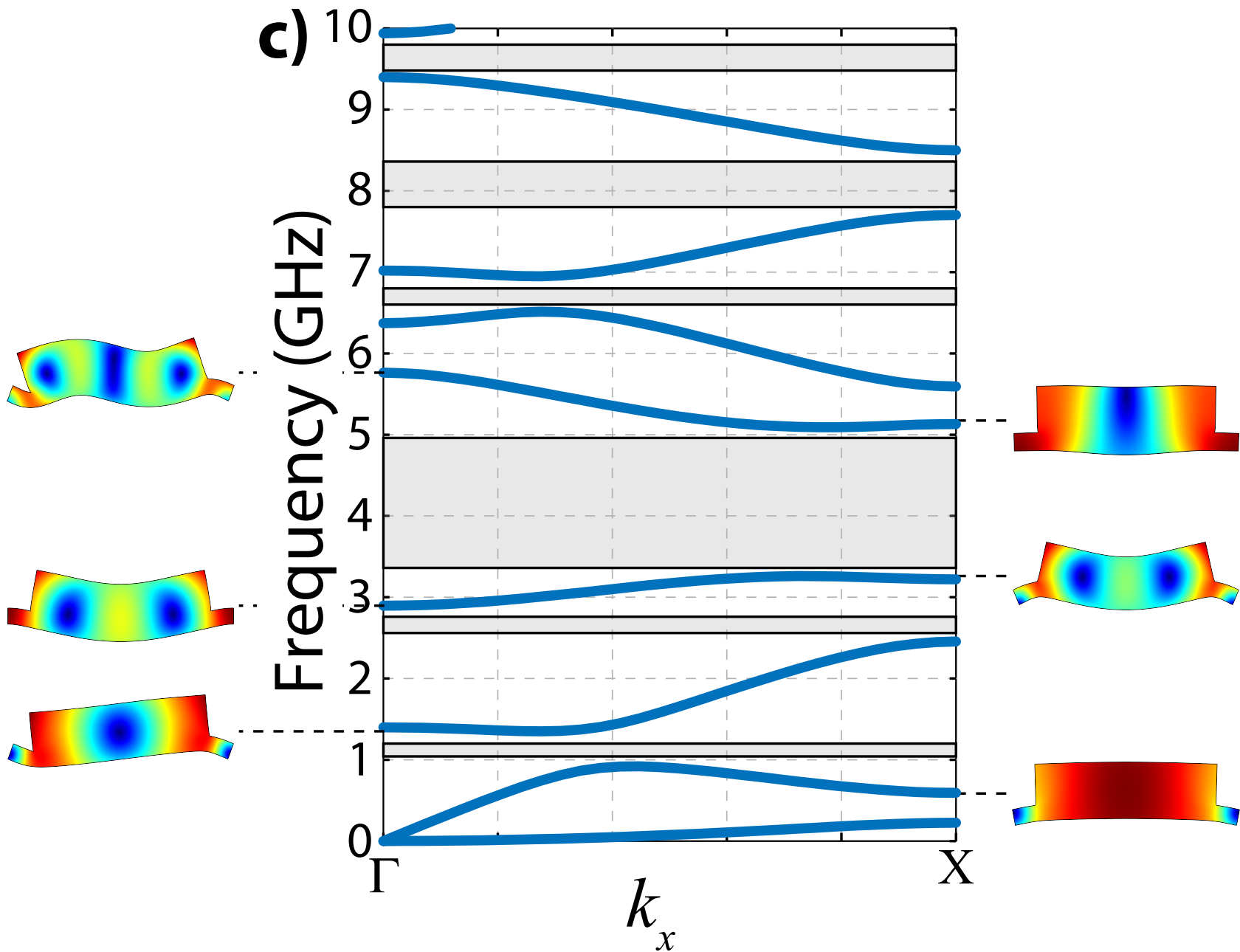
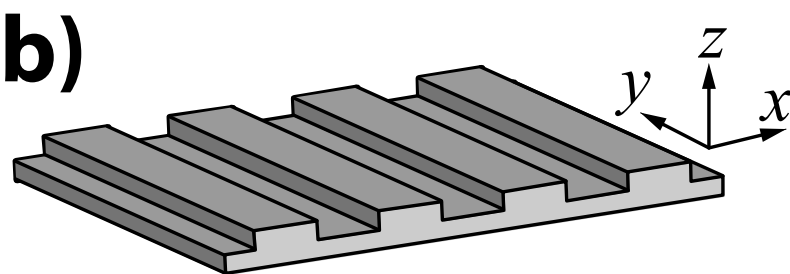
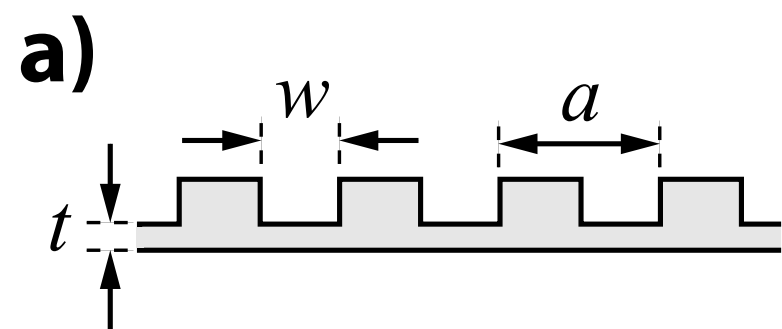
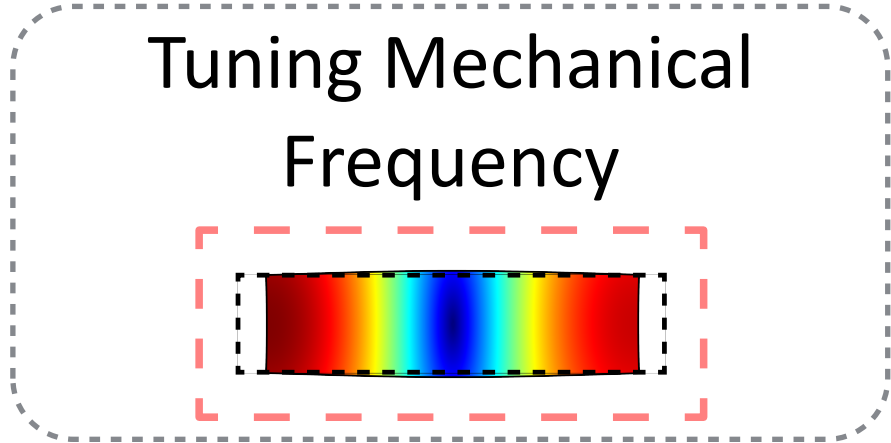
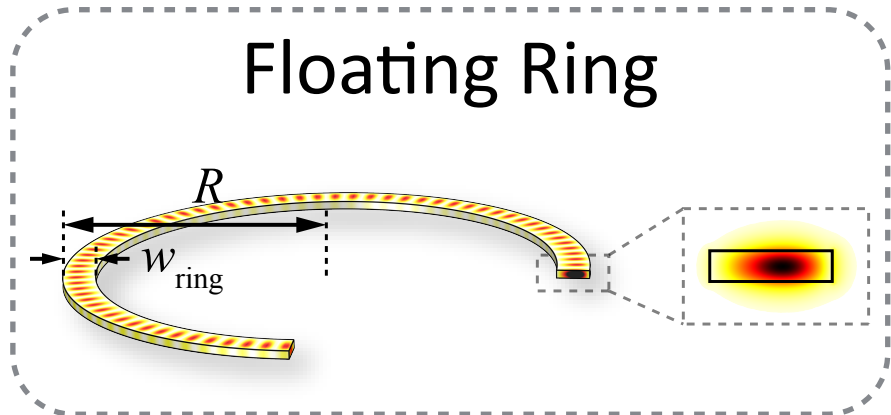


Mechanical Grating



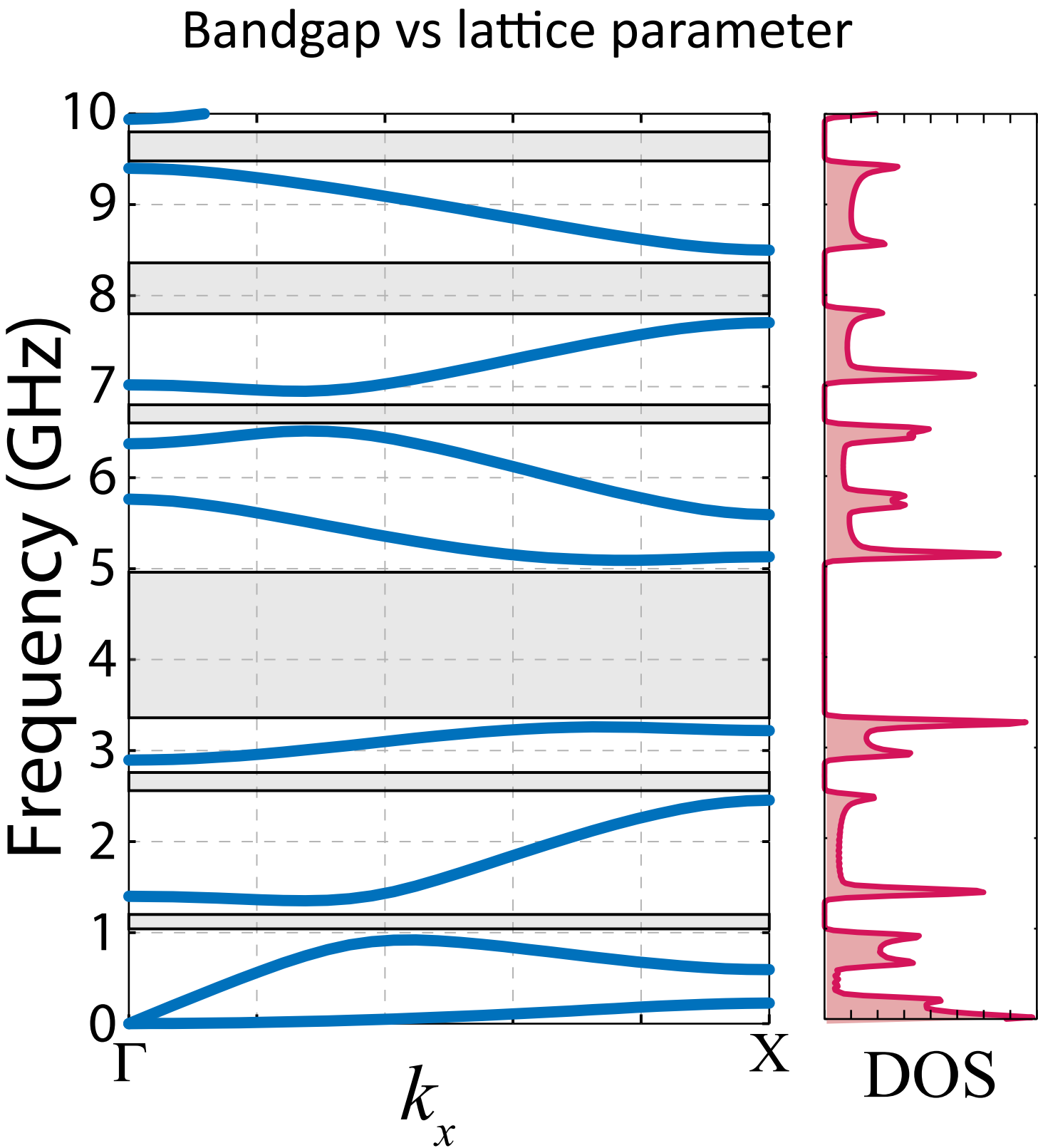
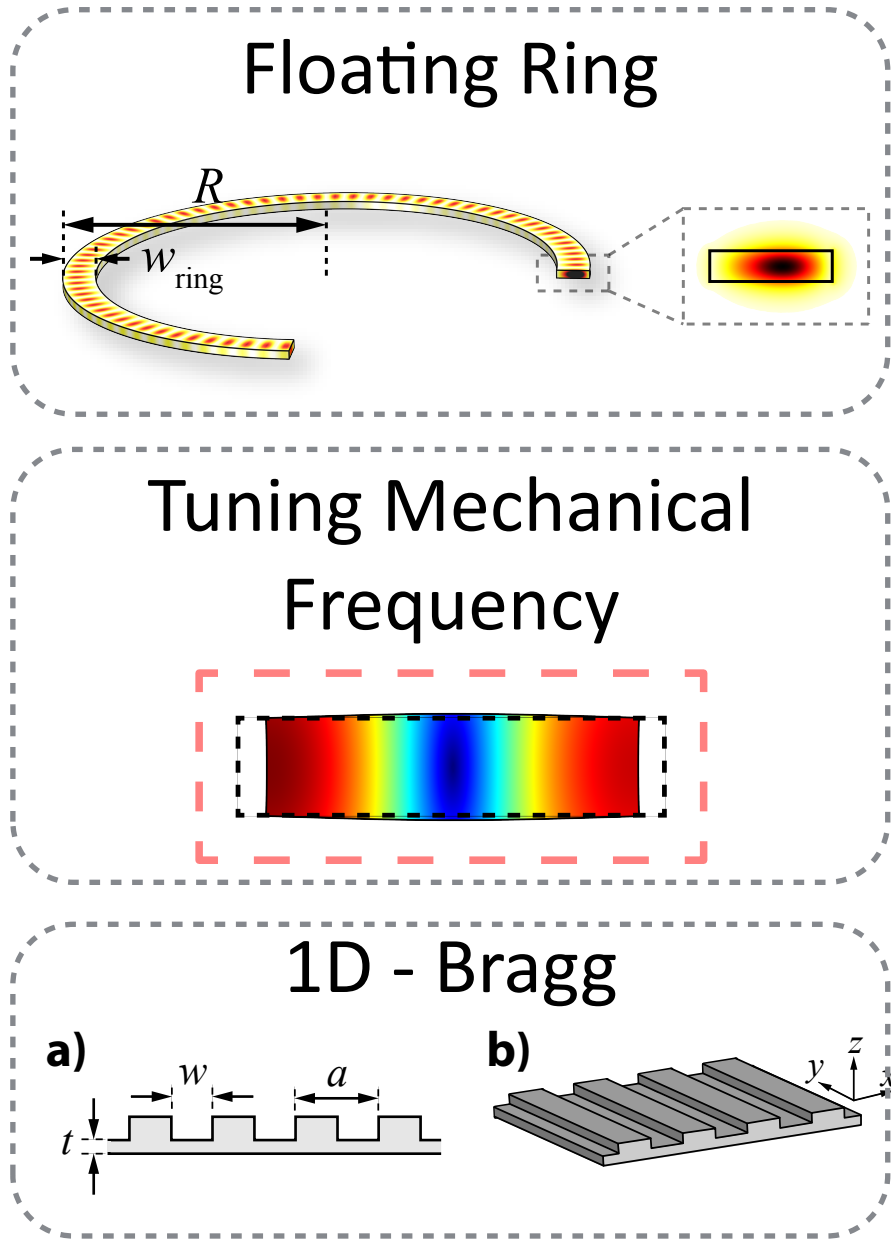


Mechanical Grating



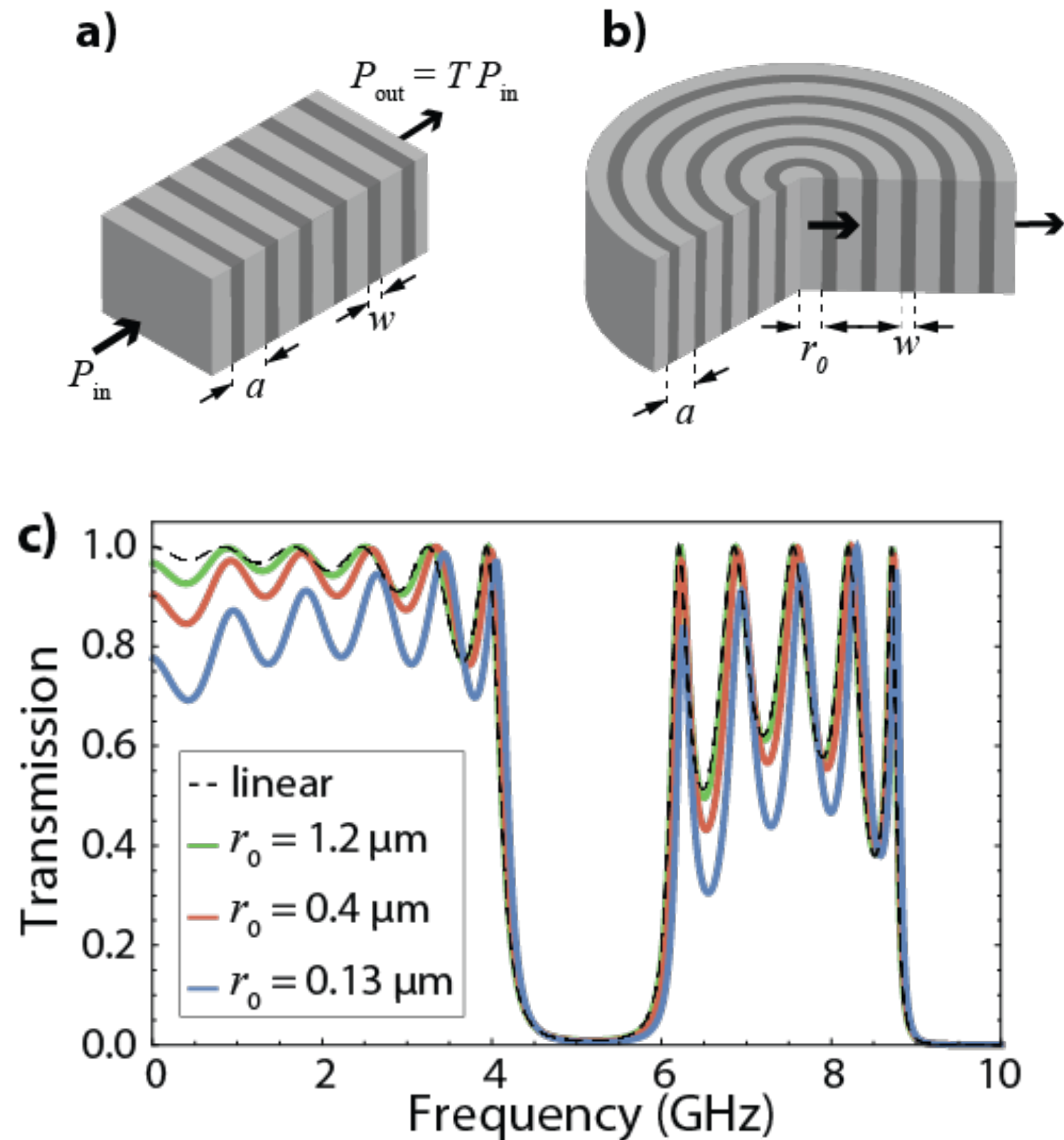


Mechanical Grating



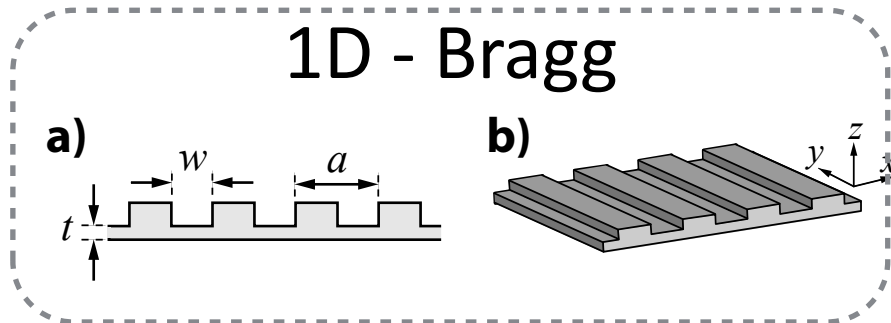
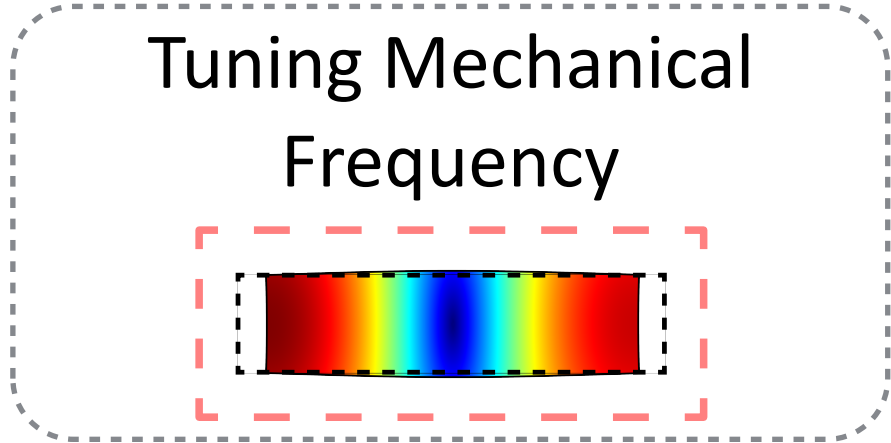
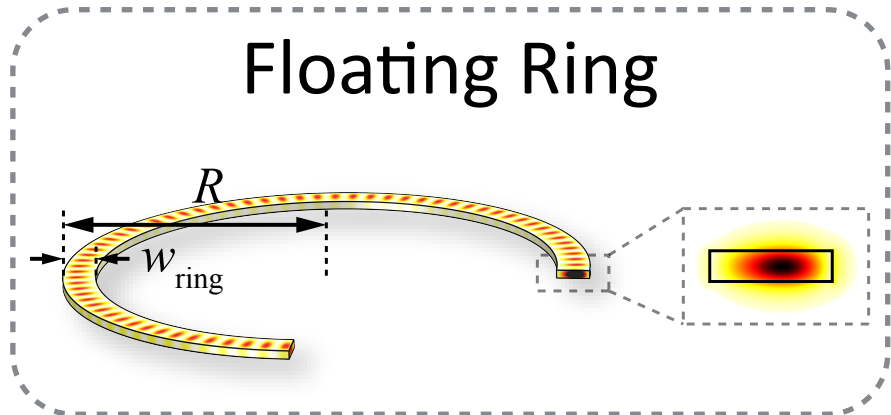


Acoustic linear grating?



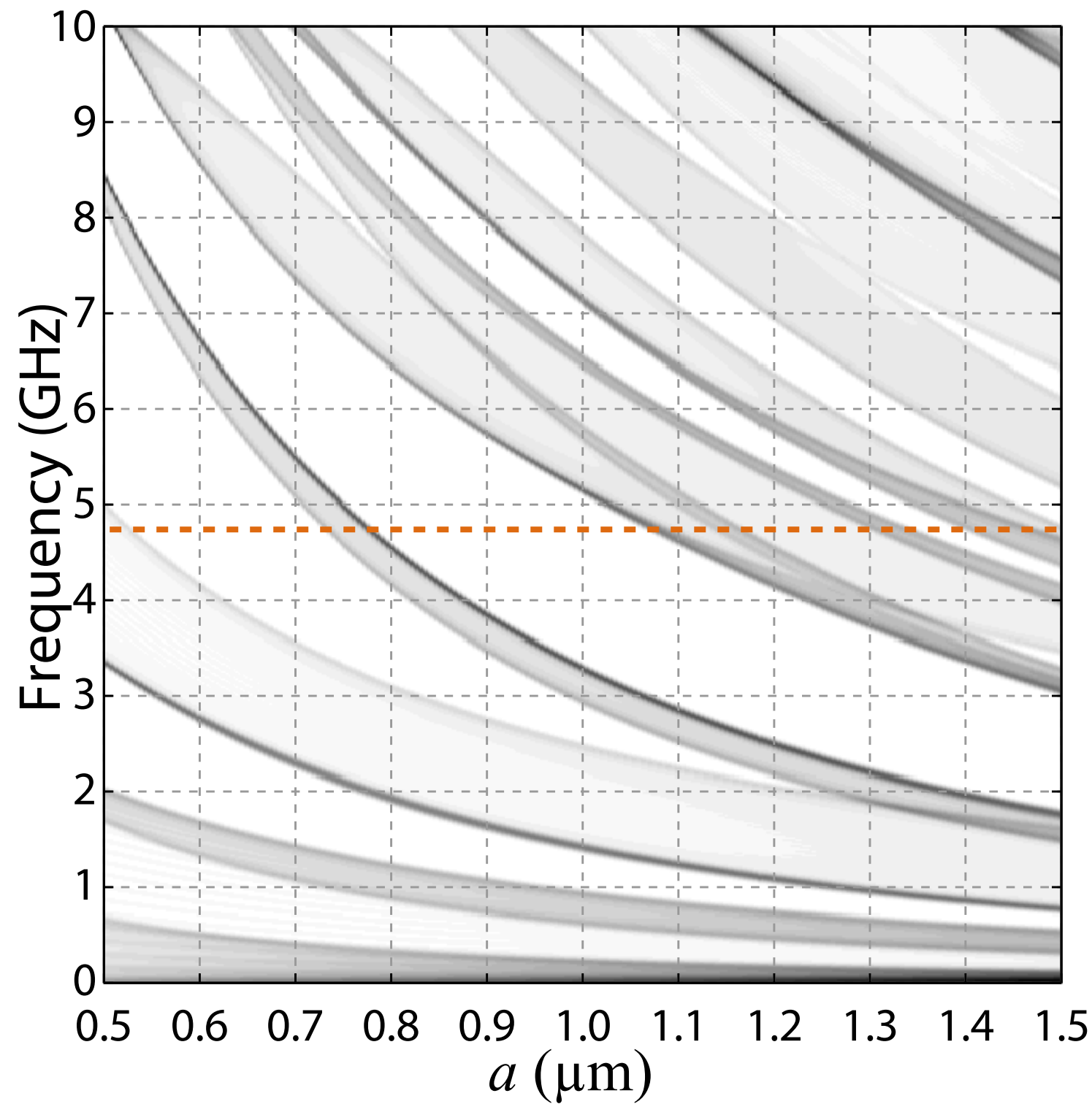


Mechanical Grating



$$v_{ac} = 9600 \text{ m/s}$$
$$w_{ring} = 1 \text{ }\mu\text{m}$$
$$f_{ac} = \frac{v_{ac}}{2w_{ring}} \approx 4.8 \text{ GHz}$$

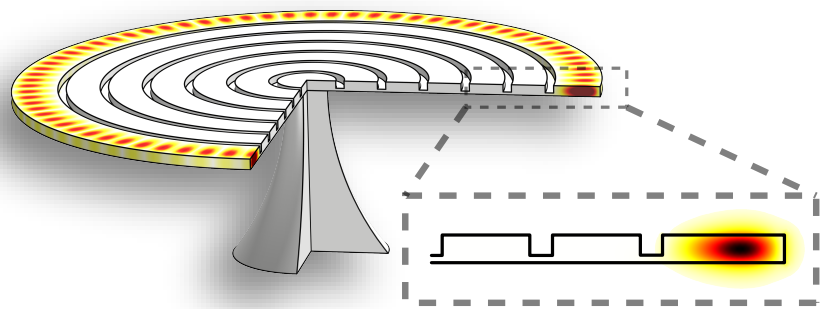
Bandgap DOS vs lattice parameter



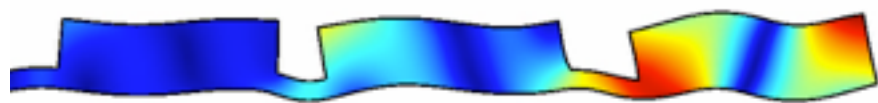


Mechanical Grating

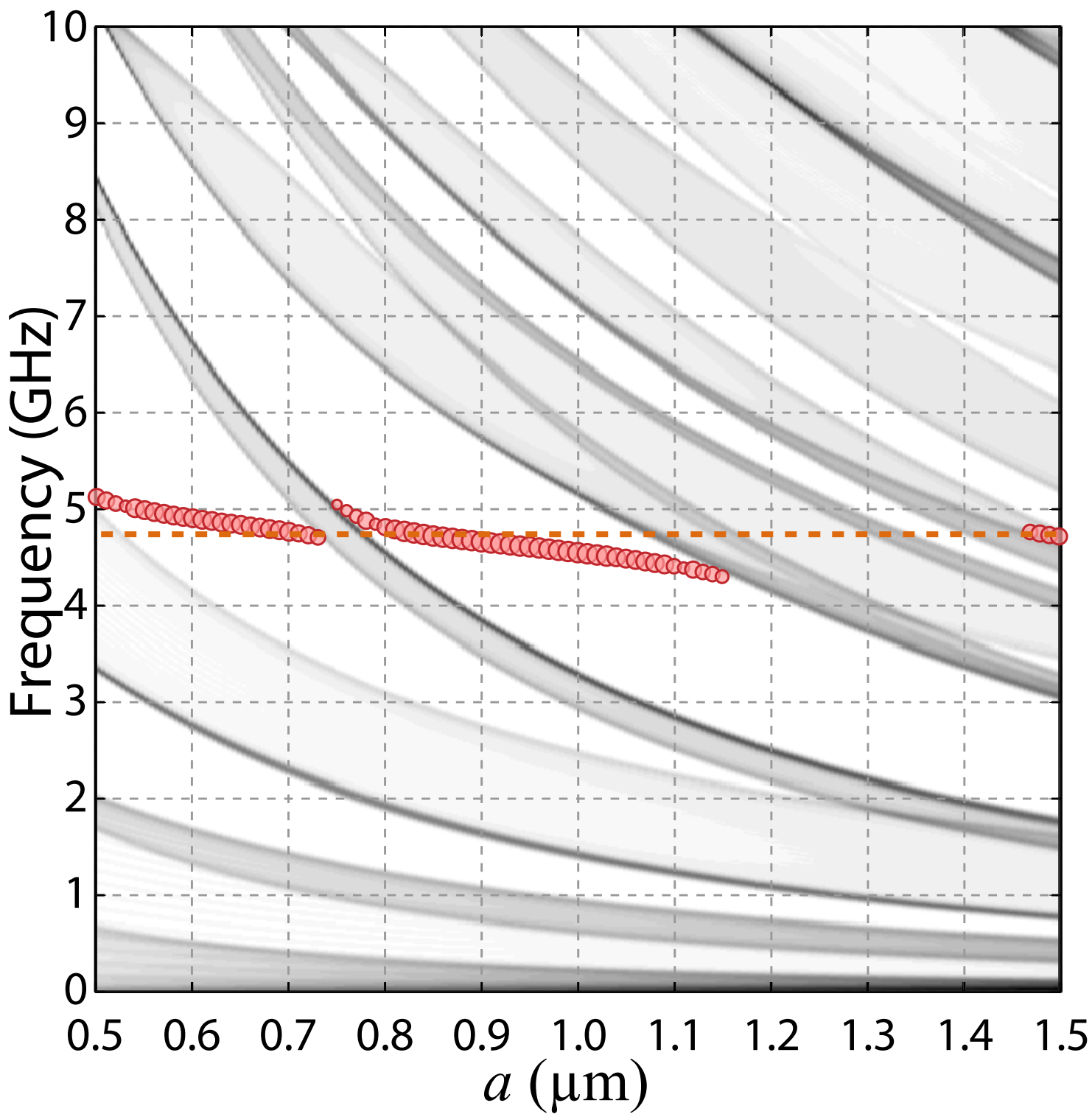
Bullseye Disk



Bullseye Disk



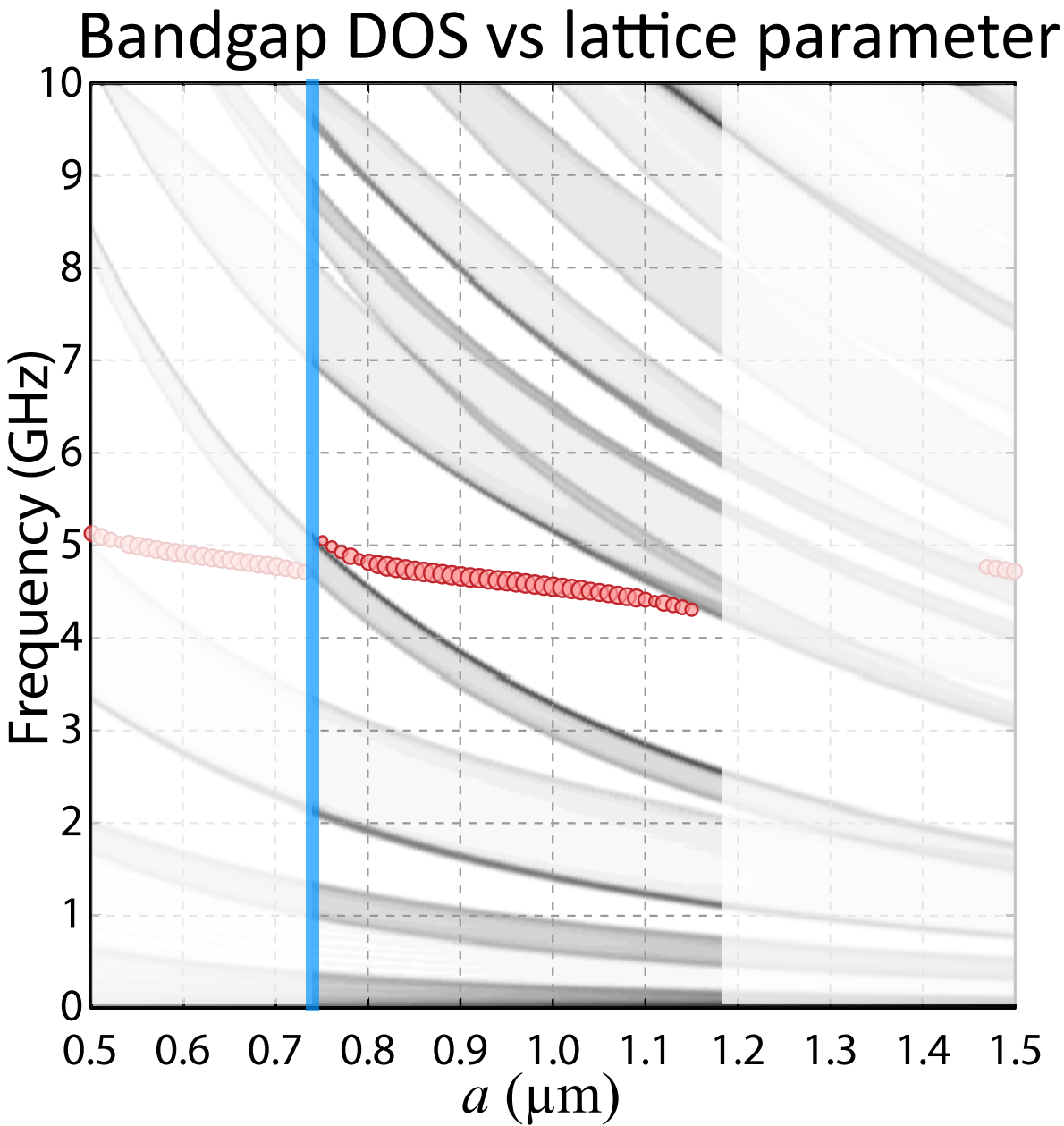
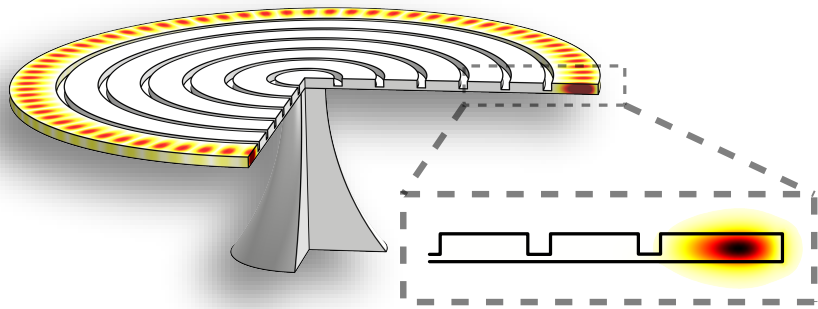
Bandgap DOS vs lattice parameter



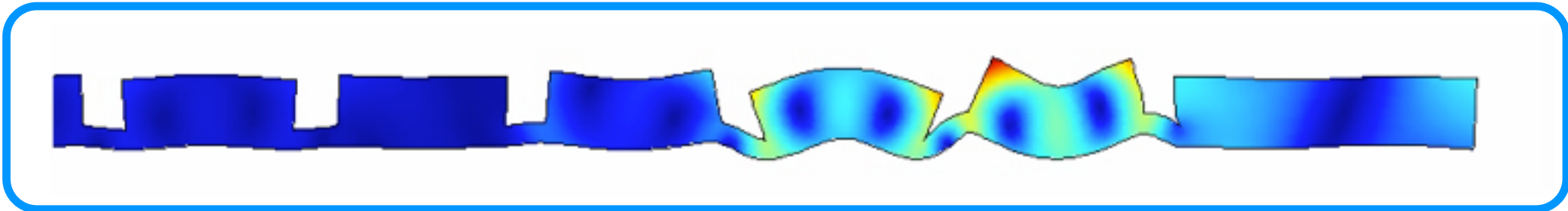


Mechanical Grating

Bullseye Disk



Simulated mode profile





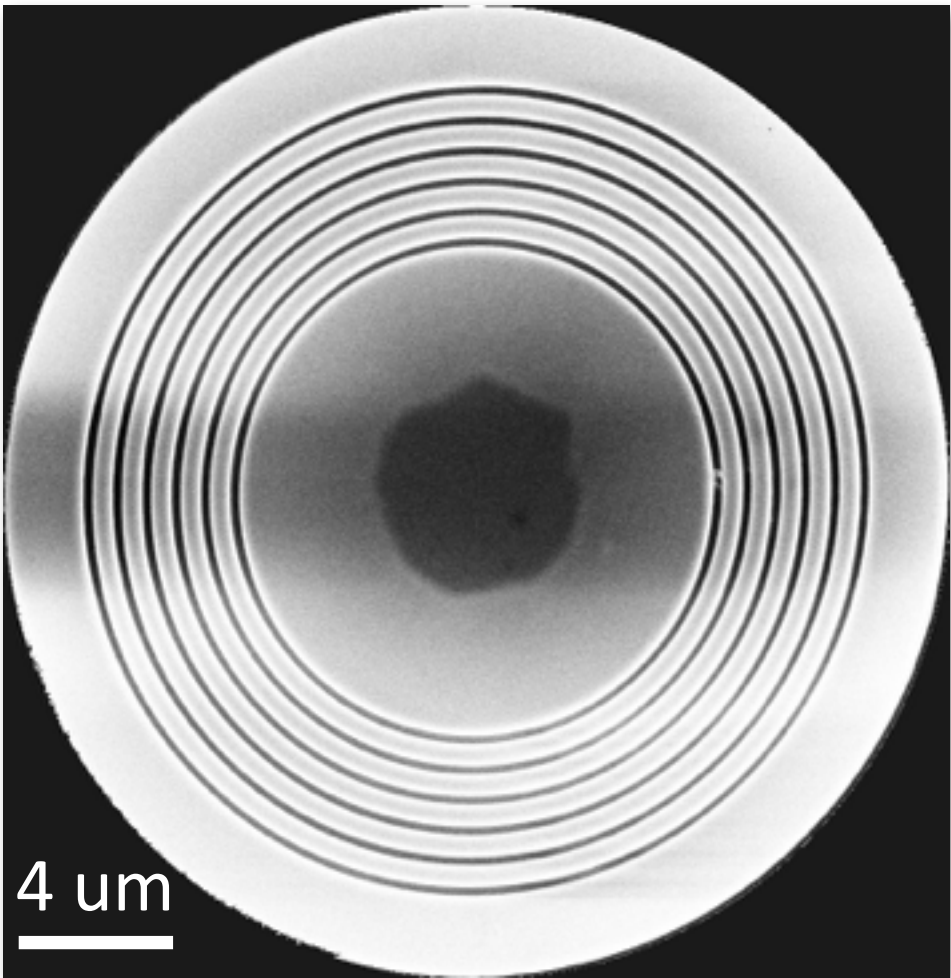
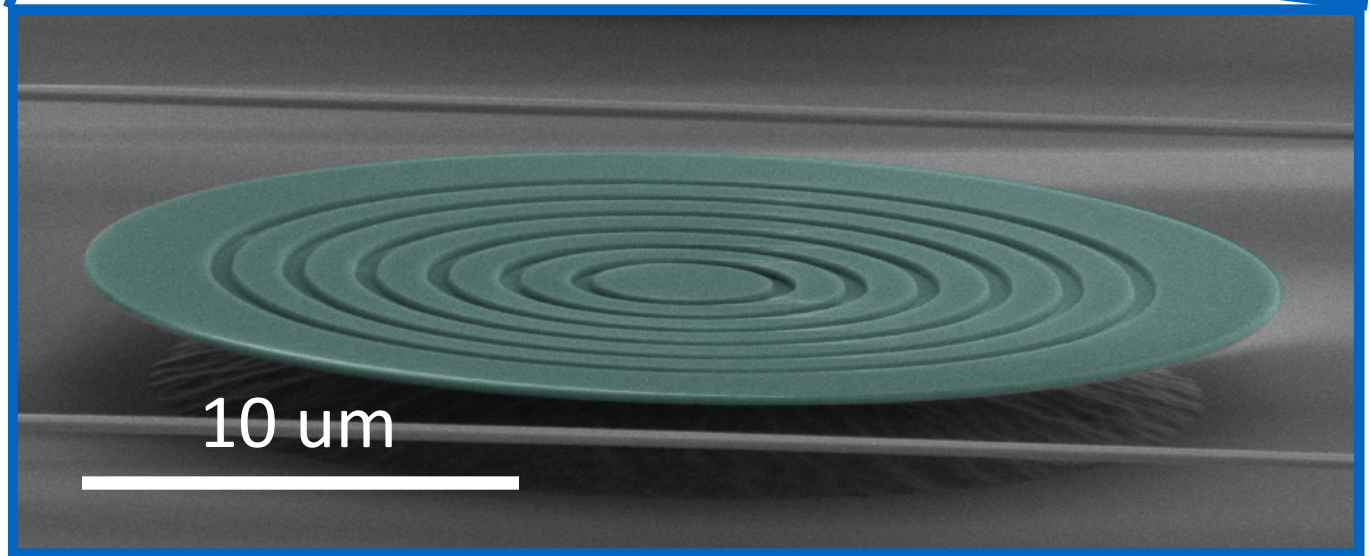
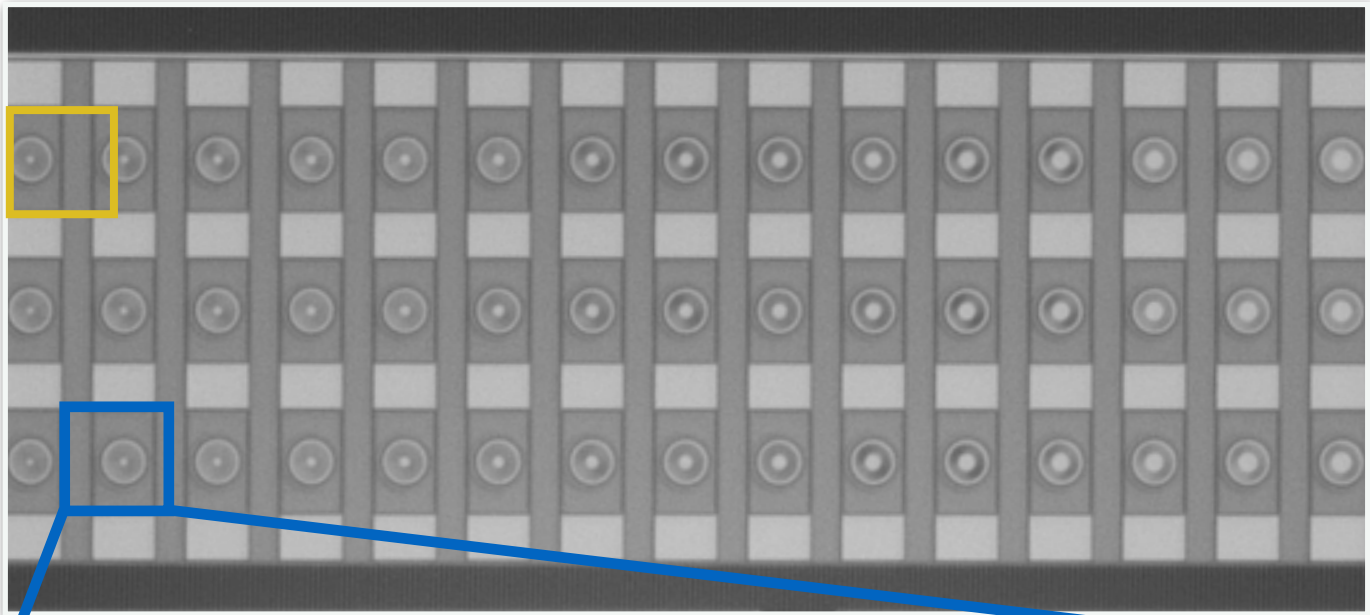
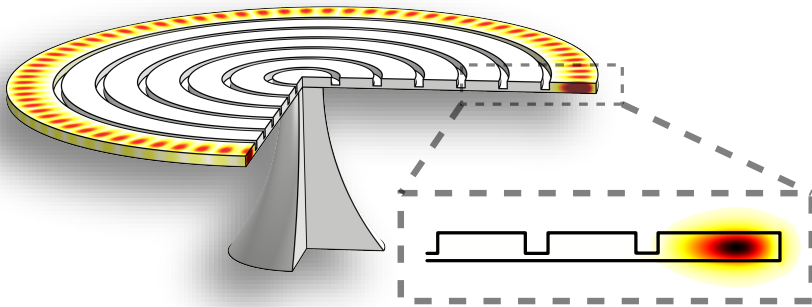
Bullseye fab: foundry based

Integrated and scalable production



Passive SiPhotonics imec-ePIXfab

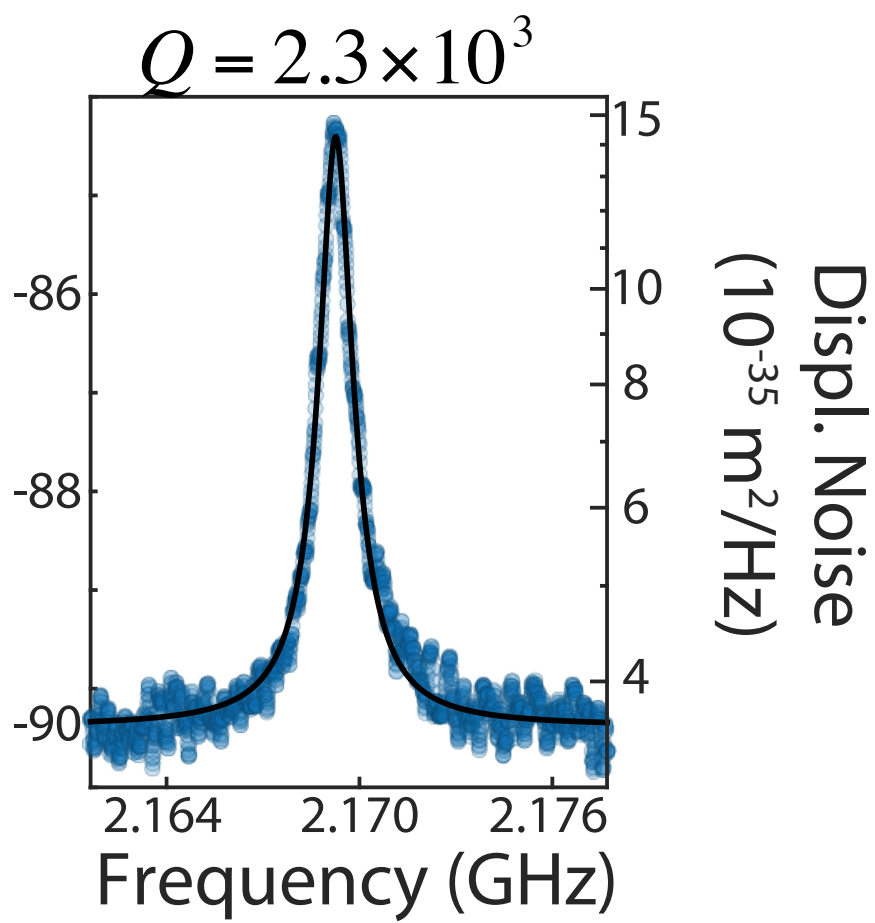
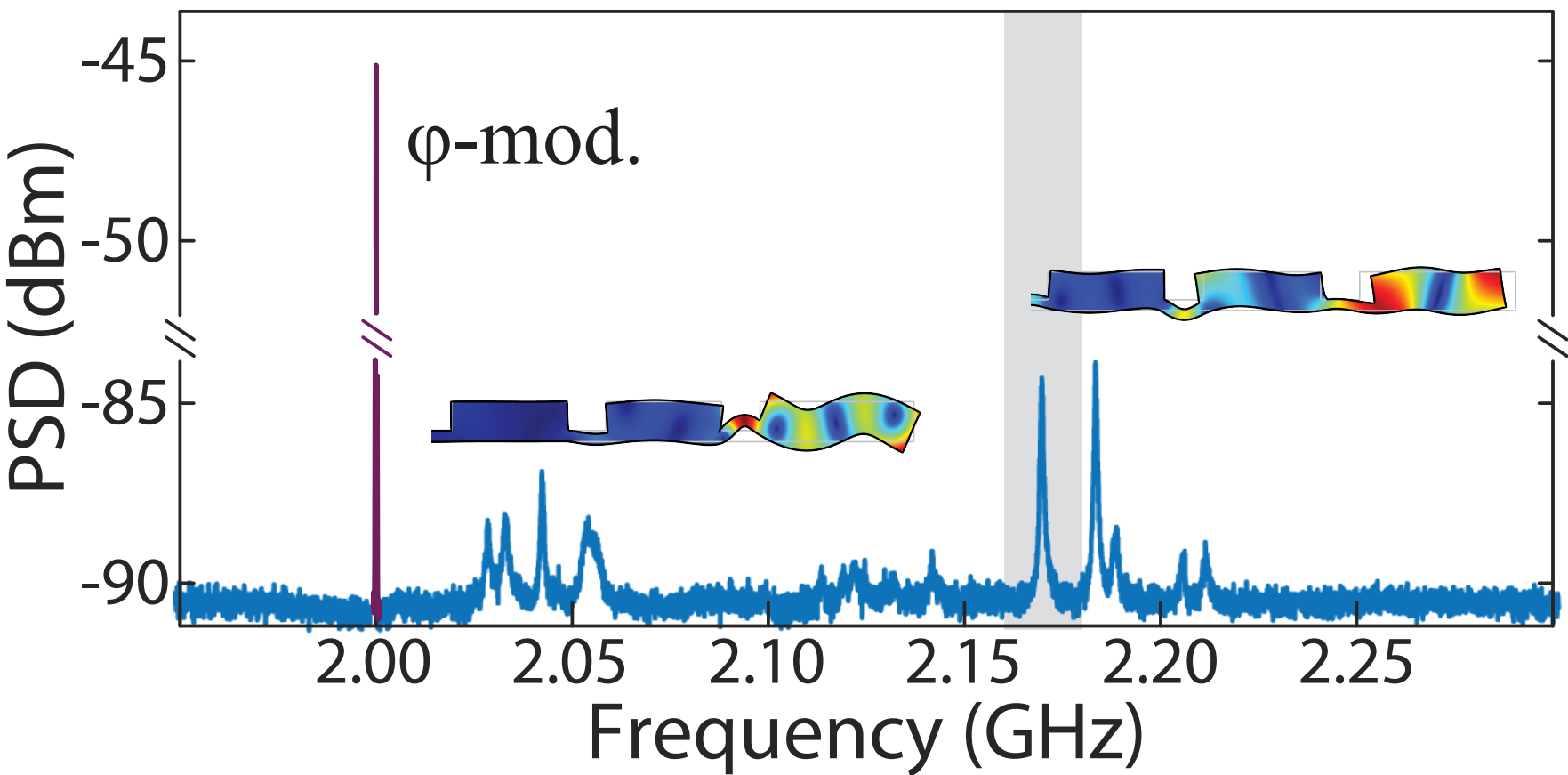
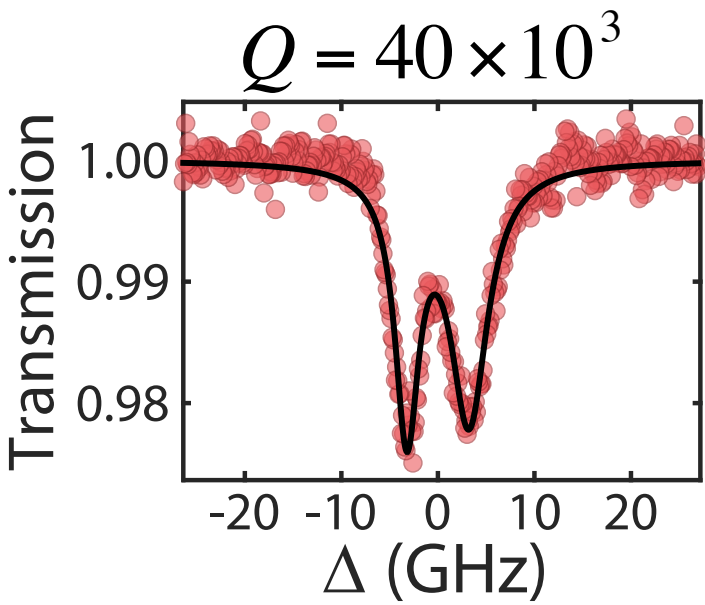
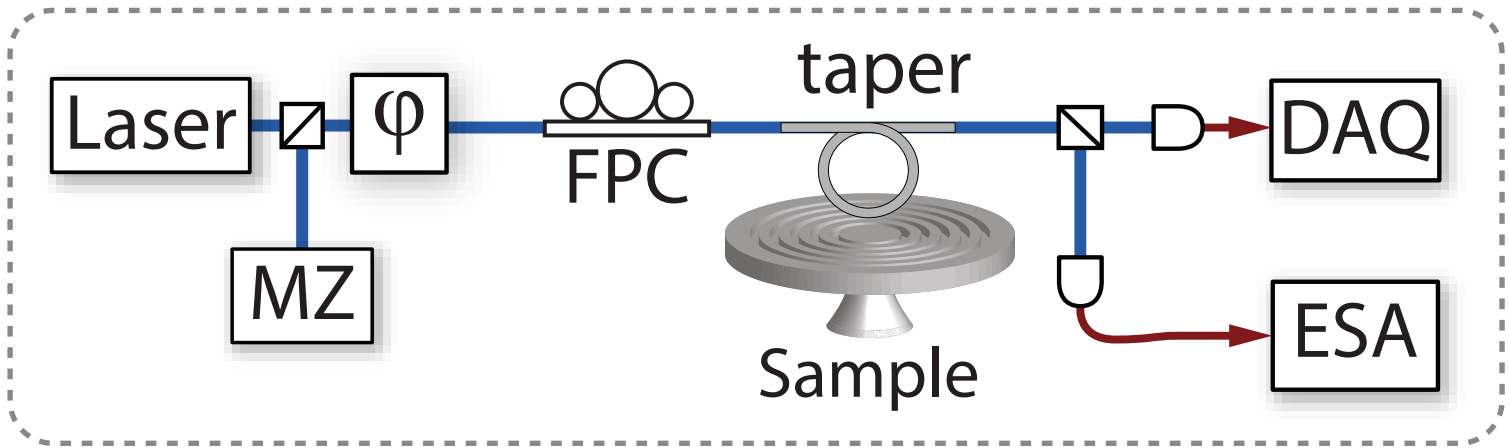
Bullseye Disk





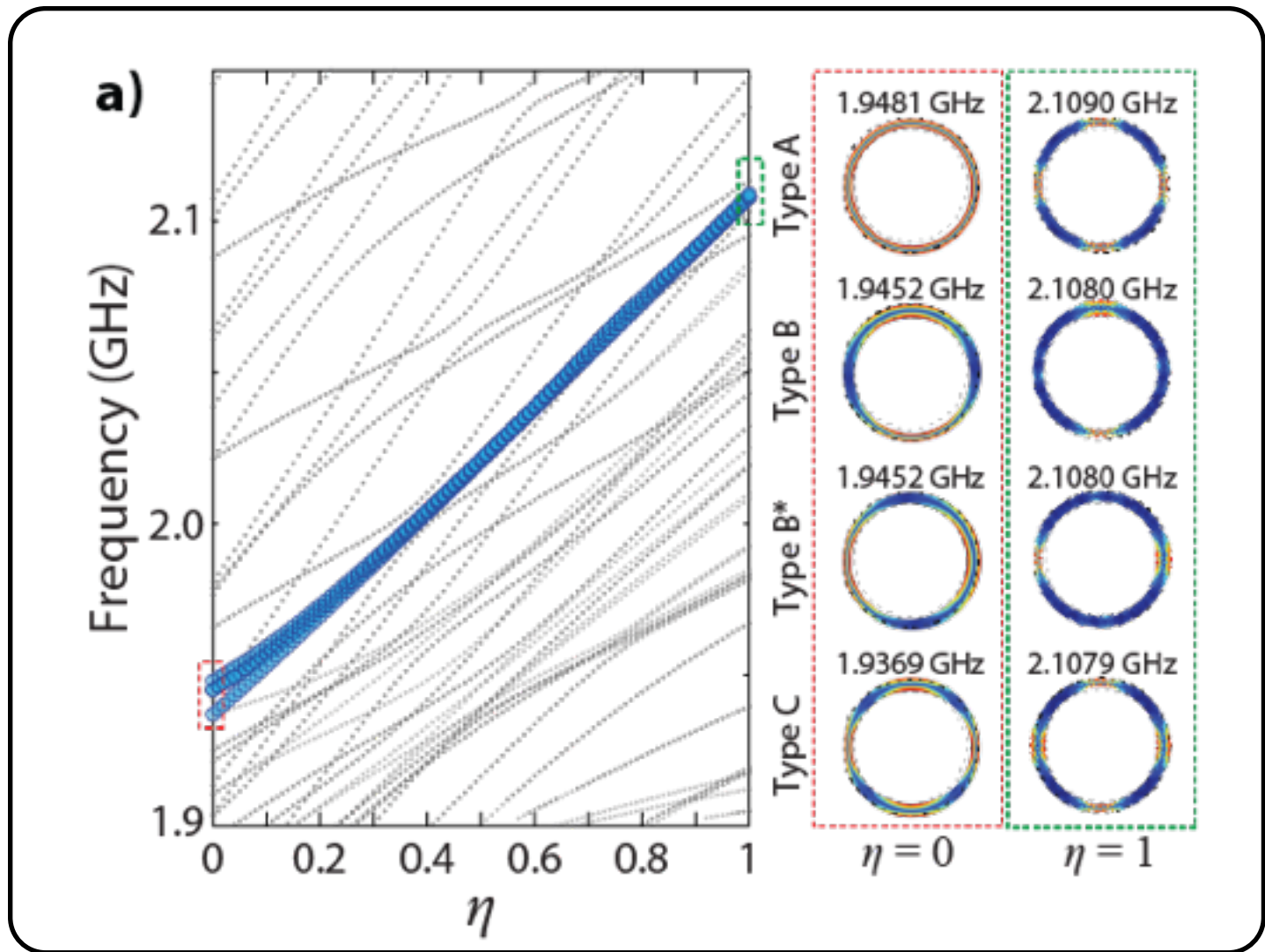
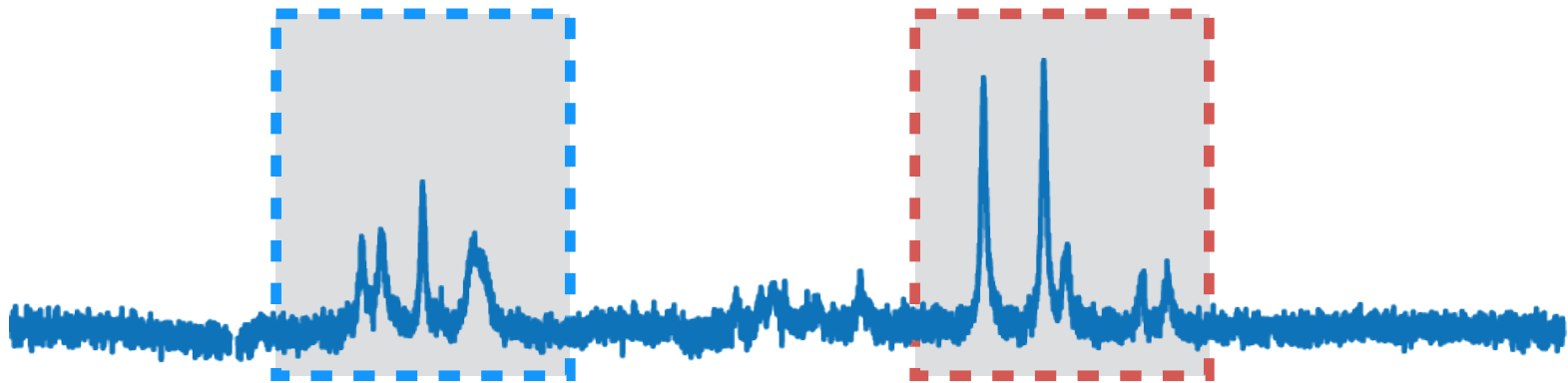
Mechanical mode testing

Experimental Setup





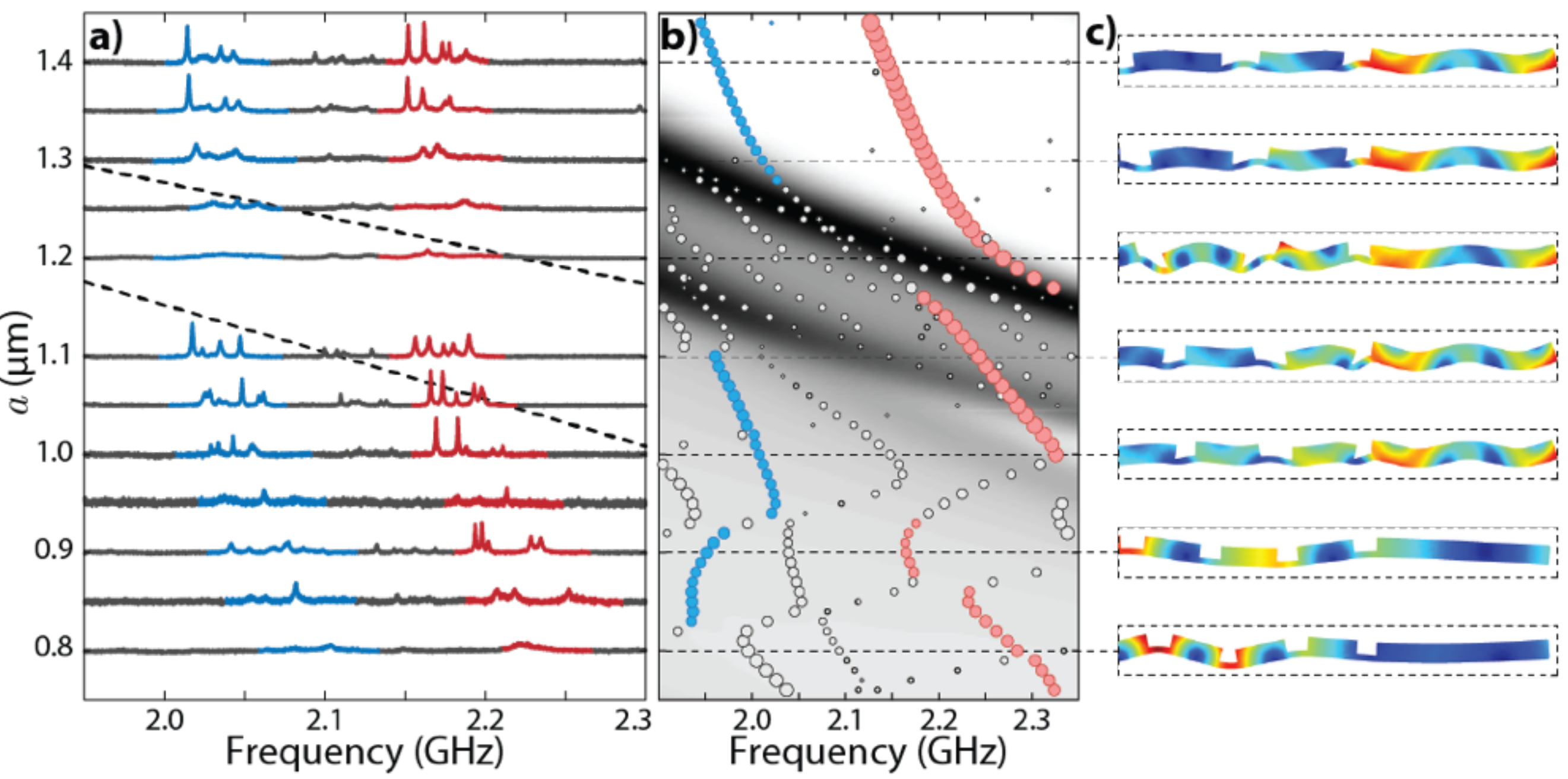
Multi-peak structure



Si Anisotropy

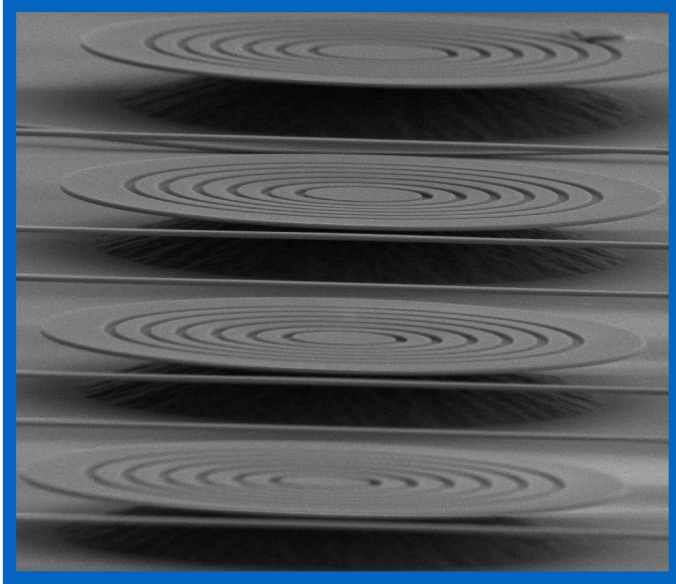


Mechanical mode testing





Perspective on Bull's eye arrays



Foundry based sample

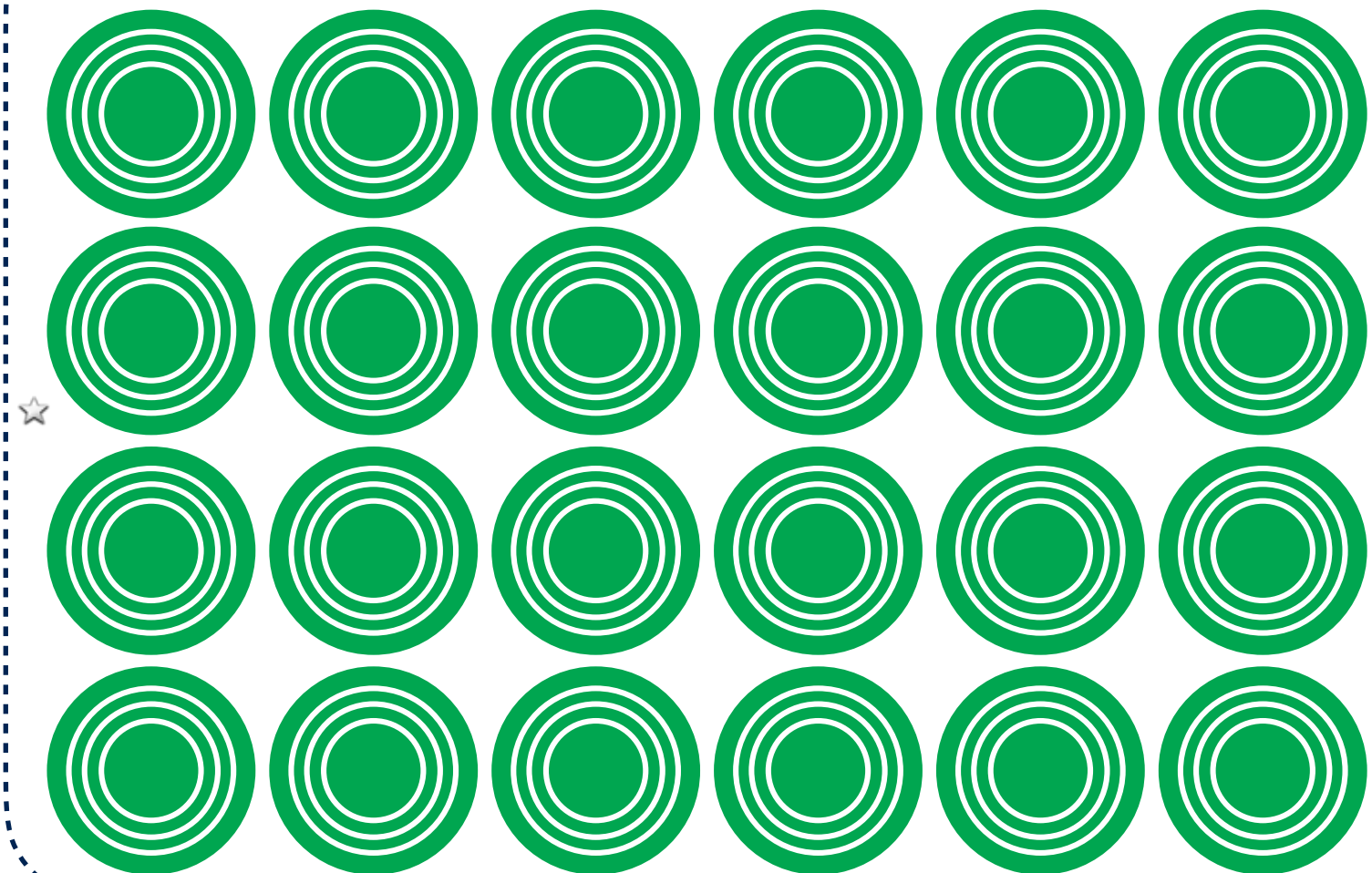
Independent
confinement
control

Mechanical:
radial grating

Optical: total
internal
reflection

- ★ Moderate optical damping rate ($Q \sim 100,000$)
- ★ High mechanical Q's (2300 @ 300K)
- ★ Array-scalable
- ★ Design flexible to different materials

Array-scalable





Outlook





Outlook

Linear Optomechanics

- Displacement detection
- Optical Spring
- Cooling & Amplification
- Two-tone drive: "Optomechanically induced transparency"
- Ground state cooling
- State transfer, pulsed operation
- Wavelength conversion
- Radiation Pressure Shot Noise
- Squeezing of Light
- Squeezing of Mechanics
- Light-Mechanics Entanglement
- Accelerometers
- Single-quadrature detection, Wigner density
- Optomechanics with an active medium
- Measure gravity or other small forces
- Mechanics-Mechanics entanglement
- Pulsed measurement
- Quantum Feedback
- Rotational Optomechanics

Multimode

- Mechanical information processing
- Bandstructure in arrays
- Synchronization/patterns in arrays
- Transport & pulses in arrays

Nonlinear Optomechanics

- Self-induced mechanical oscillations
- Attractor diagram?
- Synchronization of oscillations
- Chaos

○ White: yet unknown challenges/goals

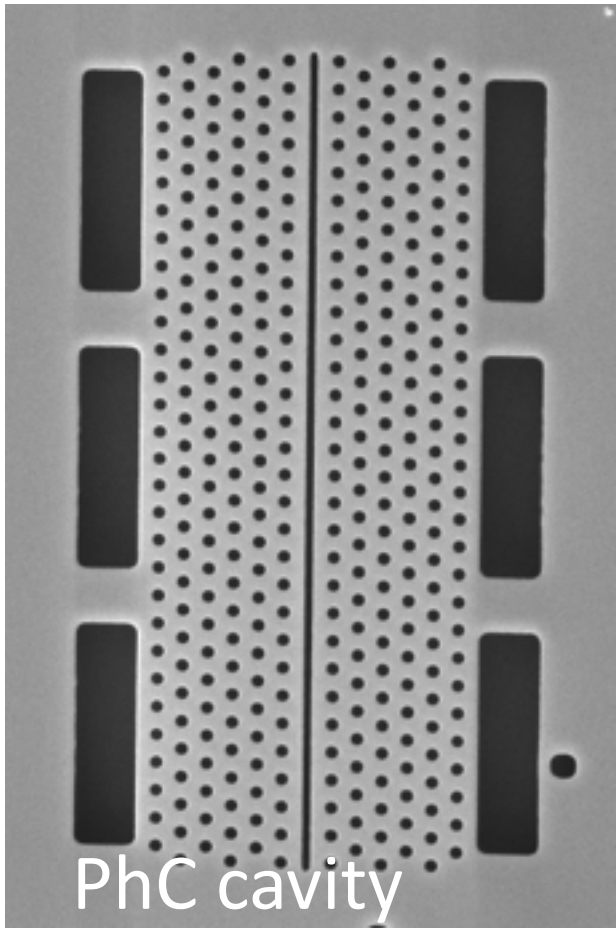
Nonlinear Quantum Optomechanics

- QND Phonon number detection
- Phonon shot noise
- Photon blockade
- Optomechanical "which-way" experiment
- Nonclassical mechanical q. states
- Nonlinear OMIT
- Noncl. via Conditional Detection
- Single-photon sources
- Coupling to other two-level systems
- Optomechanical Matter-Wave Interferometry

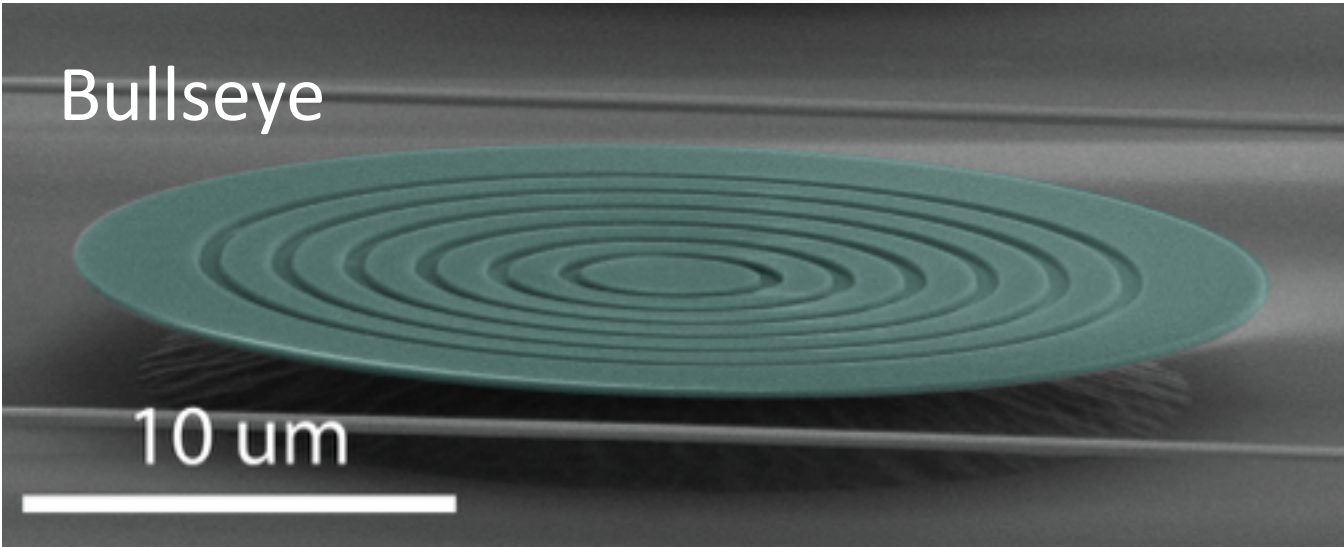


Outlook

High Q crystal
cavities



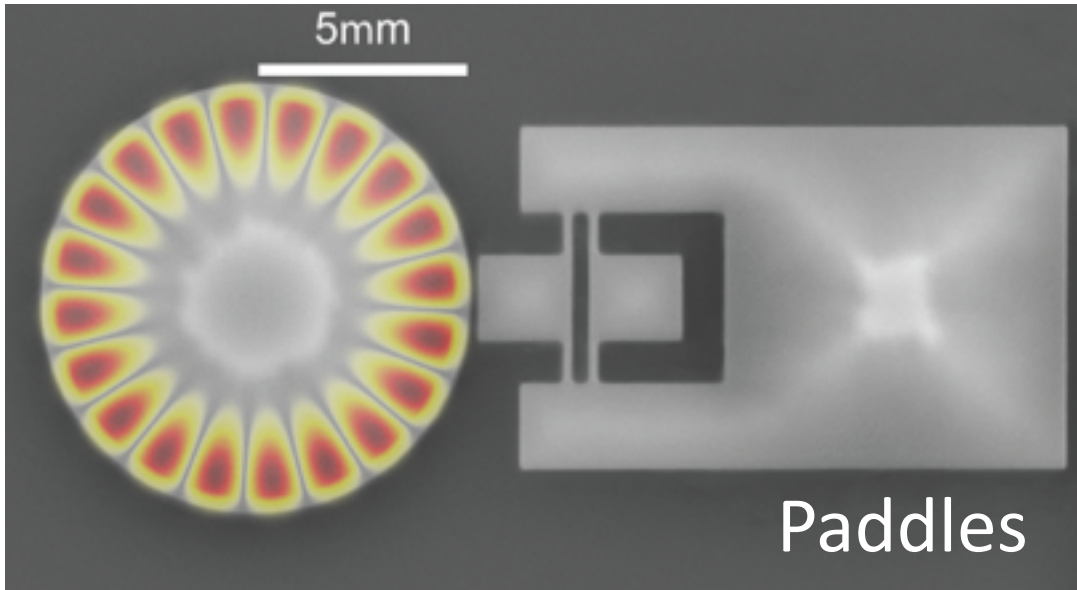
Large optomechanical coupling
on disk-like structure



Third harmonic
generation &
Frequency combs



Suppression of mechanical
radiation loss





Students



Felipe Santos



Rodrigo Benevides



Yovanny Espinel



Débora Princepe



Gustavo Luiz



Mário Machado



Lais Fujii



Guilherme Rezende



Jorge Henrique



Carlos Gois

Staff



Antônio Von Zuben



Celso Ramos

Prof.



Thiago Alegre



Newton Frateschi



Gustavo Wiederhecker

